Investigation of turbulence promoting inserts for augmenting heat transfer from gases in tubes

Velumylum Nirmalan
Iowa State University

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AUGMENTING HEAT TRANSFER FROM GASES IN TUBES

Iowa State University

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Velumylum Nirmalan

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This dissertation is dedicated in memory of my parents,

Drs. Vethanayaki and Saravanamuttu Velumylum.
NOMENCLATURE

A constant in correlation equation
A<sub>i</sub> surface area of insert in each segment, ft<sup>2</sup>
A<sub>s</sub> surface area of wall in each segment, ft<sup>2</sup>
A<sub>c</sub> cross-sectional area of teflon spacer, ft<sup>2</sup>
A<sub>x</sub> cross-sectional flow area of calorimeter section, ft<sup>2</sup>
B constant in correlation equations
C constants in correlation equations or model
C<sub>1</sub> proportionality constant in model
C<sub>2</sub> proportionality constant in model
C<sub>D</sub> orifice constant
C<sub>p</sub> specific heat of fluid packets, Btu/lbm-°F
C<sub>pa</sub> specific heat of air, Btu/lbm-°F
C<sub>pa,4</sub> specific heat of air in segment 4, Btu/lbm-°F
C<sub>pw</sub> specific heat of water, Btu/lbm-°F
C<sub>pw,4</sub> specific heat of water in segment 4, Btu/lbm-°F
D constant in correlation equations
D<sub>i</sub> inside diameter of calorimeter pipe, ft
D<sub>t</sub> diameter of location of thermocouples in calorimeter pipe wall, ft
F<sub>i-s</sub> shape factor from insert to tube wall
f Fanning friction factor
f<sub>a</sub> augmented Fanning friction factor
f<sub>o</sub> empty tube Fanning friction factor under R<sub>2</sub> or R<sub>3</sub> conditions
smooth tube Fanning friction factor
heat-transfer coefficient of air, Btu/hr-ft\(^2\)-°F
augmented heat-transfer coefficient, Btu/hr-ft\(^2\)-°F
average wall heat-transfer coefficient for air in the fourth segment, Btu/hr-ft\(^2\)-°F
average insert heat-transfer coefficient for air in the fourth segment, Btu/hr-ft\(^2\)-°F
empty tube heat-transfer coefficient, Btu/hr-ft\(^2\)-°F
thermal conductivity of fluid, Btu/hr-ft-°F
thermal conductivity of air in the fourth segment, Btu/hr-ft-°F
thermal conductivity of calorimeter wall, Btu/hr-ft-°F
thermal conductivity of teflon spacer, Btu/hr-ft-°F
pitch of insert, ft, as defined in Figure 2.9
insert entrance length for segment 1, ft
insert entrance length for segment 2, ft
axial length of a calorimeter segment, ft
insert entrance length, ft
overall length of calorimeter (distance between pressure taps), ft
mass flow rate of air, lbm/hr
mass flow rate of water, lbm/hr
number of tubes
augmented Nusselt number
Nusselt number based on inside pipe diameter
empty tube Nusselt number under \( R_2 \) or \( R_3 \) condition
smooth tube Nusselt number
P  pumping power, ft-lbf/sec
Pr  Prandtl number
P(t)  distribution function of contact time
P_{atm}  atmospheric pressure, in. of Hg
P_{st}  static pressure, in. of water
q  average heat transfer, Btu/hr
q_a  augmented heat transfer, Btu/hr
q_{a,0}  heat lost by air in the whole test section, Btu/hr
q_{a,4}  heat lost by air in segment 4, Btu/hr
q_{c,4}  convective heat transfer from the air to the tube wall in segment 4, Btu/hr
q_{cd}  heat transfer by conduction in the downstream calorimeter segment, Btu/hr
q_{ci,4}  convective heat transfer from the air to insert in segment 4, Btu/hr
q_{cu}  heat transfer by conduction in the upstream calorimeter segment, Btu/hr
q_{i(t)}  instantaneous heat transfer from fluid packet to the wall, Btu/hr
q_o  empty tube heat transfer, Btu/hr
q_{r,4}  radiant heat transfer from insert to wall in segment 4, Btu/hr
q_{w,4}  total heat gained to water in segment 4, Btu/hr
q_{wc,4}  heat gained by water after correction for axial conduction in segment 4, Btu/hr
q_{w,o}  heat gained by water in the whole test section, Btu/hr
R  universal gas constant
R^2  correlation coefficient
R_2  performance evaluation criterion under constant pressure
drop condition

R₃  performance evaluation criterion under constant pumping power condition

Reₐ  augmented Reynolds number

Reₚ  Reynolds number based on inside pipe diameter

Reₒ  empty tube Reynolds number

s  frequency of surface renewal, 1/sec

T  temperature of fluid packet, °F

T atm  ambient room temperature, °R

Tₐ₁,₄  inlet air temperature at segment 4, °F

Tₐ₂,₄  outlet air temperature at segment 4, °F

T b  bulk temperature of air, °R

T b  bulk temperature of fluid packets, °F

T b,₄  bulk temperature of hot gas in segment 4, °R

T i  temperature of eddies at first instant of renewal, °F

T i,₄  average insert temperature in segment 4, °R

T s  tube wall (surface) temperature, °F

T s,₄  average tube wall temperature in segment 4, °R

T w₄,₄  average wall temperature at the thermocouple location in segment 4, °F

T w₁,₄  inlet water temperature at segment 4, °F

T w₂,₄  outlet water temperature at segment 4, °F

t  time of contact of a fluid packet with the tube wall, sec

τₘ  mean residence time of fluid packet in tube, sec

U  overall heat-transfer coefficient, Btu/hr-ft²-°F

uₘ  mean velocity in pipe, ft/sec
W  width of insert, ft
\( W_{x_i} \) uncertainty in any variable \( x_i \)
y  radial distance from tube wall, ft
\( \alpha \) thermal diffusivity, ft\(^2\)/hr
\( \alpha \) bending angle of insert, °
\( \gamma \) twist angle of insert, °
\( \Delta p \) measured pressure drop, in. of water
\( \Delta p_c \) pressure rise due to temperature drop, lbf/ft\(^2\)
\( \Delta p_{or} \) pressure drop across orifice, in. of water
\( \Delta T \) temperature difference, °F
\( \Delta T_{LMTD} \) log mean temperature difference, °F
\( \Delta T_t \) temperature difference across teflon spacer, °F
\( \Delta X \) axial distance of teflon spacer, ft
\( \delta \) characteristic length of the fluid packet, ft
\( \varepsilon_i \) emissivity of insert
\( \varepsilon_s \) emissivity of wall
\( \theta \) bending line angle of insert, °
\( \mu \) viscosity of fluid packet, lbm/ft-hr
\( \mu_{a,4} \) viscosity of air in segment 4, lbm/ft-hr
\( \rho_a \) average density of air, lbm/ft\(^3\)
\( \rho_{a1} \) inlet density of air, lbm/ft\(^3\)
\( \rho_{a2} \) outlet density of air, lbm/ft\(^3\)
\( \sigma \) Stefan-Boltzmann constant, Btu/hr-ft\(^2\)-°R\(^4\)
\( \phi \) longitudinal angle of insert, °
\( \phi(L/D) \) mathematical function to correct for entry length effects
in smooth tube
I. INTRODUCTION

A. The Problem

The recent, though now nearly forgotten, energy crisis required thermal industries to give increased attention to methods of conserving energy. The thermal industry uses about 17% of the total energy consumed in the country to produce process steam and hot water (Yaverbaum, 1979). An increase in efficiency of water heating would significantly impact energy conservation nationwide.

One of the many methods used by the heating and ventilating industry for water heating involves combustion gases flowing inside water-jacketed tubes. Fire-tube boilers and domestic water heaters are familiar examples. Some representative applications of combustion gases flowing inside tubes other than the conventional usage mentioned above include (a) waste heat recovery in the chemical industry, (b) heat recovery in cogeneration systems, and (c) process gas cooling from steam/methane furnaces in the petroleum industry. If the heat transfer in any of these heat exchangers is improved by enhancing the convective heat-transfer coefficient on the gas side of the tube, the energy consumption or the required heating surface could be reduced, and thereby decreasing the cost of operation and the capital cost of the heating surface.

Enhancing or augmenting heat transfer is one of the many ways of reducing costs in the thermal industry. In recent years, the
exponential increase in the amount of published literature in the area of enhanced heat transfer (Bergles et al., 1983) indicates that enhanced heat transfer has developed into a major specialty area in heat-transfer research and development. Single-phase convective heat transfer can be enhanced by several techniques that can be broadly classified into passive and active techniques.

Some examples of passive enhancement techniques include surface roughness, extended surfaces, displaced enhanced devices, and swirl flow devices. These enhancement devices augment heat transfer by disturbing the existing flow. Surface roughness is produced by knurling, corrugating, or creating repeated ribs. Rough surfaces—ranging from random sand grain type roughness to discrete protuberances—enhance heat transfer by disturbing the fluid sublayer at the surface. Displaced enhanced devices, such as baffles and mixing elements, enhance heat transfer by altering the flow at the surface by disturbing the core flow. Swirl flow devices are inserts or other geometrical device which create a rotating and/or secondary flow. Examples include vortex generators, twisted-tape inserts, and axial-core inserts.

Active augmentation techniques use additional external sources of energy to enhance convective heat transfer. Some of the commonly used techniques are surface vibration, fluid vibration, surface scraping, use of electric fields, and jet impingement.

The gas convective coefficients are normally much lower than the liquid coefficients in fire-tube boilers or water heaters. One of the
most popular and easily accomplished methods to improve the gas convection and radiation is to introduce inserts inside the tubes. Inserts can be made in an infinite variety of configurations. Commonly used insert designs include twisted ribbons or tapes of near-tube-diameter width, wire coils similar to springs, and bent-strip inserts. A typical bent-strip insert configuration is a narrow, thin metal strip bent and twisted in a zig-zag fashion to allow periodic contact with the tube wall. However, outside of the twisted-tape insert, very little fundamental work has been done to establish the heat-transfer and flow-friction data required for the design of bent-strip inserts (Junkhan et al., 1982). Hence, the present investigation concentrates on establishing the performance of bent-strip inserts.

B. Scope of Investigation

To meet the objectives of predicting the performance of bent-strip inserts, three goals were established. They were

1. To acquire heat-transfer and flow-friction data for a variety of geometrical variations of bent-strip inserts in order to formulate design equations for these inserts

2. To better understand the mechanism of enhancement of convective heat transfer by visualizing flow patterns created by the inserts

3. To devise an analytical interpretation of the results so that the mechanism of enhancement is incorporated into correlation of the heat transfer.
To meet the above goals, the following research program was conducted in the Heat Transfer Laboratory of the Department of Mechanical Engineering and Engineering Research Institute at Iowa State University. A survey of the current techniques in cooling hot gases in tubes was conducted. An experimental facility was designed and built to obtain thermal-hydraulic data from hot gases flowing inside tubes. Different geometrical variations of the bent-strip inserts were tested by using the experimental facility. The heat-transfer and friction data were used to obtain an empirical correlation to predict the performance of these inserts. At the same time, two flow visualization techniques were used to study the flow patterns created by the bent-strip inserts. They provided insight into the mechanism of enhancement. Finally, the surface renewal/penetration theory was used as a basis for a model to provide an analytically based correlation for heat transfer in tubes in the presence of bent-strip inserts.
II. THERMAL HYDRAULIC EXPERIMENTS

A. Literature Survey

A literature survey using the computerized "Bibliography on Augmentation of Convective Heat and Mass Transfer - II" (Bergles et al., 1983) has revealed only limited information regarding inserts used in the cooling of gases in tubes. The three types of turbulence-producing inserts most commonly applicable to the present study are twisted tapes, wire coils, and bent strips. Figure 2.1 shows typical configurations of a wire coil, a twisted tape, and a bent-strip insert.

As early as 1896, Whitham used "loosely fitting and gently twisted" (a 360 degree twist in 10 ft) strips made of No. 10 sheet iron. The strips were inserted into forty-four 4 in. i.d. and 20 ft long horizontal tubes of a 100-hp boiler. Whitham observed a fuel savings as high as 18% when "the boiler plant is pushed and the draught is strong." Kirov (1949), 50 years later, fitted an Economic Boiler having 76 tubes of 2.74 in. i.d. with 2.5 in. wide twisted tapes having a twist ratio of 1 turn in 1.17 ft or 1 turn in 2.33 ft. With full-length tapes, the increase in boiler efficiency was about 5% to 7% which was equivalent to a fuel saving of 7% to 10%. The increase in fan horsepower was considered negligible for the forced draft system under test; however the use of twisted tapes in a natural draft system was not advised.

Heat-transfer and flow-friction references found for various inserts are described below; these are the only cases found where air or
FIGURE 2.1. Typical configurations of (a) a wire coil, (b) a twisted tape, and (c) a bent-strip insert.
combustion gases were used as the working fluid and cooling was present. To establish a uniform basis of comparison with the various references cited in this study, the pitch of an insert was defined as the length along the tube axis for one full twist of a twisted tape or one full cycle of a bent-strip insert as noted in Figure 2.1. The heat-transfer and pressure-drop data reported by the various investigators for these three types of inserts are compared in in Figures 2.2 and 2.3, respectively. The data are compared as Nusselt number and Fanning friction factor as a function of Reynolds number. The variation of Prandtl number was considered small enough to neglect Prandtl number in the heat transfer comparison. Pressure-drop data were corrected for momentum gain caused by the cooling.

An early but comprehensive study of twisted-tape turbulators with air cooling was reported by Royds (1921). Electrically heated air was allowed to flow through a tube of 2.31 in. i.d. and 6.58 ft length. A large number of tests were made with inlet air temperatures varying from 430°F to 1000°F. Nine different twisted tapes of width 1.94 in., sooted to simulate actual conditions, were tested with pitches varying from 10.9 to 88.9 in. The results obtained show that pitches larger than 20.0 in. had very little influence on heat transfer. Data were reported as curves of heat-transfer coefficients and of pressure drops as a function of mass velocities. The bulk temperature was assumed to be 550°F so that the data could be converted to standard dimensionless groups as presented in Figure 2.2. The pressure-drop data reported were
I WIRE COIL (COLBURN AND KING, 1931)
2 TT PITCH = 7 IN. (COLBURN AND KING, 1931)
3 TT PITCH = 7 IN. (EVANS AND SARGEANT, 1951)
4 TT PITCH = 14 IN. (EVANS AND SARGEANT, 1951)
5 TT PITCH = 10.9 IN. (ROYDS, 1921)
6 TT PITCH = 88.9 IN. (ROYDS, 1921)
7 TT PITCH = 3.15 IN. (THORSEN AND LANDIS, 1968)
8 BS PITCH = 10.7 IN. (JUNKHAN ET AL., 1985)
9 BS PITCH = 7.5 IN. (JUNKHAN ET AL., 1985)
10 TT PITCH = 20.8 IN. (JUNKHAN ET AL., 1985)
11 SMOOTH TUBE (GNIELINSKI, 1976)

TT - TWISTED-TAPE INSERT
BS - BENT-STRIP INSERT

FIGURE 2.2. Comparison of published heat-transfer data for wire-coil, twisted-tape, and bent-strip inserts
FIGURE 2.3. Comparison of published friction-factor data for wire-coil, twisted-tape, and bent-strip inserts
not corrected for momentum change caused by the temperature drop in air. Since this correction could not be calculated with the data reported, the pressure drop data are not compared in Figure 2.3.

Colburn and King (1931) investigated enhancement effects of twisted tapes and conically shaped wire coils by using hot air heated by electrical resistance heaters. The test section was water cooled by a helically wound copper tube soldered to the outside. End axial conduction gains to the test section were reduced by using guard coolers on sections of tube just upstream and downstream of the test section. Tube wall temperatures were measured on the outside wall of the tube at five equally spaced axial locations by thermocouples embedded in the solder between copper cooling coils. A traversing thermocouple was used to obtain an average inlet air temperature while the average outlet air temperature was measured by thermocouples placed in an insulated mixing chamber located on the outlet of the test section. Inlet air temperatures varied from $225^\circ\text{F}$ to $734^\circ\text{F}$, and inlet wall temperatures varied from $48.6^\circ\text{F}$ to $52.7^\circ\text{F}$.

Wire coils consisted of wire wound in a three dimensional conical helix, beginning with a coil essentially of the same diameter as the 2.625 in. i.d. test section and spiraling to the apex of the cone in ever-decreasing radii. The length of the cone thus formed was about 4 in. Nine of these coils were placed in the 3 ft long test section in an apex-to-base fashion. No information was given on the wire diameter and the direction of flow with respect to the cone sequence. Two full
length, snug fitting twisted tapes were also tested by Colburn and King (1931). However, only data for a 3 ft long steel tape of 7 in. pitch were tabulated. Data were reported as curves of $h_a/c_{pa}$ and pressure drop (corrected for momentum change) as functions of mass velocity through the test section. Their data, recalculated as $Nu_D$ and $f$ as functions of $Re_D$, are compared in Figures 2.2 and 2.3.

Studies on twisted-tape enhancement were also conducted by Evans and Sarjant (1951), who used electrically heated air in a setup somewhat similar to that used by Colburn and King (1931). Twisted-tape inserts of width 2.5 in. and pitches ranging from 7 to 14 in. were tested in an 8 ft long, 3 in. i.d. test section. The outside of the tube was water cooled by helically wound copper tubing soldered on the surface. Radial gas temperature and velocity measurements were made at several axial locations so that the bulk temperature and mean velocity could be calculated. Energy balances between cooling water gain and air energy loss were used to validate data for each run. The inlet air temperature was approximately 700°F while inlet Reynolds numbers varied from 3000 to 13,000. Results were presented as curves of the convective heat-transfer coefficients and of pressure drops as functions of Reynolds numbers for different twisted tapes and coils. Data for twisted-tape pitches of 7 and 14 in., converted to $Nu_D$ and $f$ as a function of $Re_D$, are compared in Figures 2.2 and 2.3. Bulk gas temperature was again assumed as 550°F to enable the calculation of Nusselt number and friction factor.
Evans and Sarjant attempted to separate the effects of radiation and convection in twisted-tape enhancement. While noting that their analysis was not conclusive in this respect, they estimated the maximum contribution by radiation as 25% of the measured heat-transfer coefficient when gas temperatures above about 1000°F were encountered. The contribution from radiation dropped rapidly as the gas temperature was reduced.

More recent tests of the effects of twisted tapes on cooled air flow were reported by Thorsen and Landis (1968). Inlet air temperature to a 1 in. i.d. water cooled test section ranged up to 475°F. Twisted tape pitches were 3.15, 5.17, and 8.00 in., and the tube gap was about 0.003 in. A major concern of this work was the assessment of the variable properties. A temperature-ratio correction factor was determined primarily from heating data because of the limited values of \(\frac{T_s}{T_b}\) that could be attained in cooling. The best correlation of the cooling data was achieved by ignoring the temperature-ratio correction. Heat-transfer and friction-factor data for the smallest pitch are given in Figures 2.2 and 2.3.

In a study conducted for DOE at Brookhaven National Laboratory, it was found that the addition of turbulators to residential oil-fired boilers resulted in fuel savings of 2% to 8% (McDonald et al., 1979). Details of the turbulator installation were not given in the report.

More recently, two commercial, arbitrarily designed bent-strip inserts and one twisted-tape insert were tested by Junkhan et al.
Tests were conducted using electrically heated air with inlet air temperatures around 300°F. One bent-strip insert of pitch 10.7 in. and width 0.875 in. displayed a 175% increase in the heat-transfer coefficient at a Reynolds number of 10,000, while the other insert of pitch 7.5 in. and width 0.75 in. displayed a 135% increase in the heat-transfer coefficient at the same Reynolds number. The two bent-strip inserts were made by different commercial manufacturers and bent in different fashions. The twisted tape, having a pitch of 28 in. and a width of 2.6 in, provided a 65% increase in the heat-transfer coefficient. The friction-factor increases accompanying the heat-transfer coefficients were about 1000% for the bent-strip inserts and 160% for the twisted-tape insert. The Nu and f data for the three inserts are also compared in Figures 2.2 and 2.3, respectively.

The literature was also reviewed to determine a method of obtaining an estimate of the net heat transfer due to radiation from the insert to the tube wall. Watanabe et al. (1984) indicate that a major portion of the heat transfer can be caused by radiation in the case where the fluid is being heated. In their case, the temperatures of the tube wall and the insert were of the order 1750°F. Beckermann (1984) concludes that even for small temperature differences of 40°F the radiation effects could be considerable. In his case, the temperature of the flue gas was in the range of 600°F to 1300°F. He also indicates that in the case of large tube-length-to-tube-diameter ratios, it is possible to neglect axial radiation and consider only radial radiation.
The literature reviewed above indicates that most of the research on turbulators is focused on twisted-tape inserts. Comparison of heat-transfer data in Figure 2.2 show that augmentation of heat transfer achieved by the use of bent-strip inserts is higher than twisted-tape inserts. The wire-coil inserts tested by Colburn and King (1931) attain an augmentation level marginally higher than bent-strip inserts but at a considerably higher pressure drop penalty, as seen in Figure 2.3. The review of literature shows that very little fundamental work has been done to obtain the heat transfer/flow friction basis for the design of wire-coil or bent-strip inserts. The present experimental study concentrates on establishing the heat transfer/flow friction basis for the design of bent-strip inserts.

B. Experimental Apparatus

An apparatus was developed to determine the thermal hydraulic performance of various inserts. The test apparatus consists of a blower suitable to produce flows in the Reynolds number range of 3000 to 20000, a heater box to heat the air to over 300°F, and a test section in which the actual performance of the inserts is to be evaluated. Prior to its use for insert tests, the performance of the apparatus as a smooth tube without inserts was obtained and compared with accepted correlations.

1. Flow system

The flow system that provided hot air to the test section is shown in Figure 2.4. Airflow was provided by a centrifugal blower that
FIGURE 2.4. Flow system (not to scale)
delivered 600 SCFM at a head of 6 in. of water. After leaving the fan, the air passed through a blast gate and into a 6 in. i.d. aluminum tube. This tube contained flow straighteners consisting of 1 in. diameter and 3 ft long plastic tubes packed inside the aluminum tube. After a flow development length of 10 ft, the air entered an orifice meter. Eight feet downstream from the orifice meter, the air passed through an elbow before entering the heater box.

The heater box was fabricated from 22-gage sheet steel in three parts: an inlet rectangular diffuser section varying in cross section from 6 in. x 6 in. at the entrance to 21 in. x 11.25 in. over a length of 72 in., an 8 in. long heater housing section, and a converging section with its cross section reducing from 21 in x 11.25 in. to 6 in. x 6 in. over a length of 50 in. The housing section for the heater contained a three-phase Chromalox Process Air Duct Heater, Model ADHT-010XX, rated at 10 kW, 240 V. A General Electric motor-driven three-phase autotransformer was used to control the power to the heater. An ammeter and a voltmeter were included in the circuitry to monitor the current and voltage input to the heater. An over-temperature protector for the heater was included in the electrical circuit to shut off the electric power to the heater if a predetermined limit temperature was reached. A motorized mixer was provided to minimize thermal stratification within the heater box.

The hot air coming out of the heater box passed through a bend and entered a 2.67 in. i.d. and 20 ft long drawn-on-mandrel steel tube,
which acted as the developing length required before the hot air entered the test section. After passing through the test section, the air was exhausted into a 5 ft long pipe section before being discharged outside the building. The heater box and the pipe leading into and out of the test section were covered with a layer of calcium silicate insulation for this high temperature service. A layer of glass fiber insulation was added to reduce the heat loss to a minimum.

During the latter part of the investigation, the flow loop had to be dismantled and moved to the Heat Transfer Laboratory in the new Mechanical Engineering/Engineering Science and Mechanics Building. The configuration of the flow loop was redesigned to reduce the floor space required. In its rebuilt form, as shown in Figure 2.5, the loop was about 40 ft long with the blower at one end and the heater box at the other end. After the air was heated, the hot air proceeded through two 90-degree bends separated by a short straight section, flowed above the heater box through the test section, and was vented to the outside through the roof. Care was taken to preserve the upstream conditions for the test section.

2. Test section

The test section shown in Figure 2.6 was based on a 3 in. o.d. carbon, drawn-on-mandrel steel tube with an i.d. of 2.67 in. The tube was 71.75 in. long and was flanged at the ends to enable easy installation and removal of inserts. To facilitate the attachment of thermocouples on the tube wall, 1 in. long and 1/8 in. wide axial slots
FIGURE 2.5. Elevation view of reassembled test loop
SPIRALLY WOUND 1/4" COPPER TUBE SOLDERED TO TUBE SURFACE
HARD ASBESTOS GASKET
TEFLON SPACER
WALL THERMOCOUPLES
INSULATION
SEGMENT 1
SEGMENT 3
SEGMENT 2
SEGMENT 4
PRESSURE TAP
WATER FLOW IN
PRESSURE TAP
WATER FLOW OUT
THERMOCOUPLE FITTING TO MEASURE OUTLET AIR TEMPERATURE DISTRIBUTION
THERMOCOUPLE FITTING TO MEASURE INLET AIR TEMPERATURE DISTRIBUTION

FIGURE 2.6. Test section schematic (not to scale)
were milled partway through the tube wall. These slots were provided at five axial locations corresponding to the ends of the four calorimeter segments described below. At each axial location, three slots were milled equidistant around the circumference. Also, each end of the test section was fitted with three pressure taps spaced equally around the circumference.

The water calorimeter for cooling the gas and conducting energy balances was made of 0.25 in. i.d. copper tubing wound around the tube and soldered to the outside with a low temperature (430°F) silver solder. Small 0.1 in. i.d. copper tubes were placed in the axial slots made for the thermocouples before the water calorimeter copper tubing was soldered to the outside of the test section. The calorimeter tubing was connected as four separate segments in series.

To minimize the axial conduction from the test section, teflon spacers were placed between the test section flanges and the inlet and outlet tube flanges. Finally, the test section was insulated with glass fiber.

3. Instrumentation

The basic measurements needed were the air and water mass flow rates and the various air and water bulk temperatures. Wall temperatures were also required for calculation of the heat-transfer coefficients. In addition, the pressure drop over the entire test section was required to calculate the friction factor.
Tube wall temperatures were measured with thermocouples fitted into the 15 axial slots milled partway into the test section wall. The welded thermocouple junctions were inserted into the slots through the small copper tubes passing under the calorimeter segments. The water temperature was measured by thermocouples attached to the calorimeter tubing at the inlet and outlet of each of the four calorimeter segments. The thermocouples installed as described above, permitted the determination of four sectional-average, gas-side heat-transfer coefficients. The bulk air temperatures at the inlet and outlet of the test section were measured by using four thermocouples placed at different radial locations just upstream and downstream of the test section. Also, to estimate the axial conduction at the ends of the test section, four thermocouples were placed on both sides of each teflon spacer to measure the temperature differences across the spacers. The axial conduction gains by the test section were calculated from the spacer dimensions, the thermal conductivity of the spacers, and the temperature differences across the spacers. The thermocouples were made of 28 ga chromel-alumel wire and were independently referenced to a Whittaker model BRJ14-50TP chromel-alumel 150°F constant temperature junction.

The static pressure at the test section inlet and the pressure drop across the test section were measured by using a Setra Systems model 239 differential pressure transducer. This transducer was connected via a six-channel Scanivalve to a piezometer ring attached to the pressure
taps at the inlet and outlet of the test section. To measure very small pressure differences, a Meriam model 34FB2 micromanometer, which is capable of reading 0.001 in. water, was used.

The air flow rate, controlled by the blast gate at the blower outlet, was measured by using the ASME orifice. The pressure drop across the orifice was measured by the differential pressure transducer, which was also connected via the Scanivalve to a piezometer ring attached to the pressure taps. The same system was used to measure static pressure at the orifice. The mass-flow rate was found by using standard ASME procedures and equations.

The water flow rate was measured directly using a container on a sensitive balance and a digital stopwatch. This was considered more accurate than using a rotameter because of the low water flow rate and variations in water temperature.

The outputs from the constant temperature junction and the pressure transducer were recorded digitally using a data acquisition system consisting of a Hewlett Packard (HP) 9845B Desktop computer, a HP 3495A 40-channel scanner, a HP 3456A digital voltmeter, two HP 9885M flexible disk drives, and a HP 2316G printer. Finally, the water-flow rate and atmospheric pressure were input into the computer by hand.

4. Test procedure

The test procedure consisted of starting the blower and switching on the electric heater. The air mixer in the heater box was also started at this time to reduce stratification of the low velocity flow
in the heater box. The system was allowed to heat up, a process that normally took about three hours. The air-flow rate was adjusted for a specific run requirement by setting the blast gate at the blower outlet. The data acquisition system was programmed to scan the air inlet and outlet water temperatures to determine when a steady state was obtained. The water flow was adjusted to obtain at least a 40°F temperature rise in water from test section inlet to outlet. Once steady state was obtained, temperature and pressure readings were measured in about 90 seconds by the data acquisition system.

5. Data reduction

The data acquisition program (listed in Appendix A) was used to promptly reduce the data, and the energy balance for the overall test section was checked by using the equation

\[ q_{w,0} - q_{cu} - q_{cd} = q_{a,0} \quad (2.1) \]

Agreement of both sides of Equation 2.1 to within 5% was required to validate a run; however, energy balances typically agreed to within 2%.

Heat transfer to each calorimeter segment was determined by writing the energy equation for the segment. The calculation procedure was started at the fourth segment. The process of calculation described below was repeated using the inlet air temperature at one segment as the outlet air temperature for the next segment. The second and the third segments had no conduction correction; however, heat gain by the coolant in the first segment required a conduction correction similar to that used for the fourth segment.
The calculation procedure for inserts was started from the fourth segment instead of the first segment because the mixing effect of an insert reduces possible temperature stratification in the flow and thereby yields a better mixed-mean outlet air temperature. In the case of the smooth tube, the calculation procedure was started from the first segment, because this method was considered better in the case in which there was no insert in the test section.

For the fourth segment, the net energy added to the water from the air flowing in the segment was obtained from

\[ q_{w,4} - q_{cd} = m_w c_w (T_{w2,4} - T_{w1,4}) - q_{cd} \]  \hspace{1cm} (2.2)

where \( q_{cd} \) is the estimated heat transfer to the water by conduction from the downstream ducting. Then, the air-side energy balance,

\[ q_{w,4} - q_{cd} = q_{a,4} = m_a c_{pa,4} (T_{a1,4} - T_{a2,4}) \]  \hspace{1cm} (2.3)

was used to find the air inlet temperature to the segment, \( T_{a1,4} \). The convective heat transfer from the air to tube wall is given by

\[ q_{c,4} = q_{a,4} - q_r,4 \]  \hspace{1cm} (2.4)

where \( q_r,4 \) is the radiant heat transfer from the insert to the tube wall. The average wall temperature near the surface for this segment, \( T_{w4} \), was taken as the linear average of the six wall temperature readings at the ends of each segment. The log-mean temperature difference between the bulk air temperature and the tube wall in the fourth segment was obtained by using the equation given below.
\[ \Delta T_{\text{LMTD}} = \frac{(T_{a1,4} - T_{t,4}) - (T_{a2,4} - T_{t,4})}{\ln \left( \frac{T_{a1,4} - T_{t,4}}{T_{a2,4} - T_{t,4}} \right)} \]  

(2.5)

The overall heat-transfer coefficient times surface area from the hot air to the wall thermocouples is given by

\[ (U_s A_s)_4 = \frac{q_{c,4}}{\Delta T_{\text{LMTD}}} = \left[ \frac{1}{h_{a,4} \pi D_{i} L_{c}} + \frac{1}{2 \pi k_{s} L_{c}} \ln \left( \frac{D_{t}}{D_{i}} \right) \right]^{-1} \]  

(2.6)

The convective heat-transfer coefficient, \( h_{a,4} \), was obtained by combining Equations 2.5 and 2.6 and is given by

\[ q_{c,4} = \left[ \frac{1}{h_{a,4} \pi D_{i} L_{c}} + \frac{1}{2 \pi k_{s} L_{c}} \ln \left( \frac{D_{t}}{D_{i}} \right) \right]^{-1} \frac{(T_{a1,4} - T_{t,4}) - (T_{a2,4} - T_{t,4})}{\ln \left( \frac{T_{a1,4} - T_{t,4}}{T_{a2,4} - T_{t,4}} \right)} \]  

(2.7)

To calculate \( q_{r,4} \), the radiation transfer model used to estimate the radiation heat transfer from the insert to the wall is

\[ q_{r,4} = \frac{\sigma (T_{i,4}^4 - T_{s,4}^4) A_{i}}{1/F_{i-s} + (1/\varepsilon_{i} - 1) + (1/\varepsilon_{s} - 1)(A_{i}/A_{s})} \]  

(2.8)

where suffixes \( i \) and \( s \) are for the insert and tube wall, respectively. The shape factor, \( F_{i-s} \), cannot be assumed to be unity as the strip insert is bent in such a fashion that it can radiate to itself. To simplify the calculation, one can assume constant wall and insert.
temperatures in each segment of the test section. This is reasonable as far as the insert is concerned, since the conduction from the insert to the tube wall can be neglected because of the very small contact area between the tube wall and the insert. Data from Kreith (1962) and Feingold (1966) were used to estimate the shape factor from the insert to the wall. The insert was assumed to be two finite plates with a common edge joined at different angles depending on the pitch-to-diameter ratio.

The convective heat transfer from the hot gas to the insert is given by the following equation:

\[ q_{ci,4} = h_{i,4} (T_{b,4} - T_{i,4}) \]  

(2.9)

where the bulk air temperature, \( T_{b,4} \), is the average of the inlet and the outlet air temperatures in the fourth segment. By considering an energy balance for the insert in one segment and neglecting any conduction in the insert,

\[ h_{i,4}(T_{b,4} - T_{i,4}) = \frac{\sigma(T_{i,4}^4 - T_{s,4}^4)}{1/F_{i-s} + (1/\varepsilon_i - 1) + (1/\varepsilon_s - 1)(A_i/A_s)} \]  

(2.10)

The system of Equations 2.2 through 2.10 required an iterative solution. Initially, assuming \( q_{r,4} = 0 \), \( h_{a,4} \) was obtained in Equations 2.2-2.7. Then assuming \( h_{a,4} = h_{i,4} \) in Equation 2.10, the value of \( T_{i,4} \) was obtained. Equation 2.8 was used to calculate \( q_{r,4} \) when the value of \( T_{i,4} \) was known. The value of \( q_{r,4} \) was used to correct further the
convective heat transfer from the hot gas to the tube in Equation 2.4. A new value of $h_{a,4}$ was calculated by using Equation 2.7 and was compared with the previous value. This iteration was repeated until the value of $h_{a,4}$ converged to within 1%.

The corrections caused by radiation were about 2% to 5% of the total heat transfer for the temperature and Reynolds number ranges of this study. In order to verify the assumptions regarding the calculation procedure of the radiant component of heat transfer, the insert temperature was measured for one insert, and the radiant heat transfer from the insert to the wall was estimated. These results, given in Appendix B, indicated that by using the above assumptions the radiant component of heat transfer is underpredicted by about 30%. The most likely cause for this discrepancy is the assumption of $h_{i,4} = h_{a,4}$. However, since the total radiant component is only about 2 to 5% of the convective heat transfer, the uncertainty in the radiation correction based on these assumptions does not affect the convective heat-transfer coefficient by more than 2%.

The average Fanning friction factor is given by

$$f = \left[ \frac{\Delta p}{\left( \frac{m_a}{A_x} \right)^2 \left( \frac{1}{\rho_{a2}} - \frac{1}{\rho_{a1}} \right)} \right] \left/ \frac{4L_t \rho_a u_m^2}{D_i^2} \right. \right]$$  

(2.11)

$$f = \left[ \frac{\Delta p}{\left( \frac{m_a}{A_x} \right)^2 \left( \frac{1}{\rho_{a2}} - \frac{1}{\rho_{a1}} \right) D_i \rho_a A_x^2} \right] \left/ \frac{2 m_a^2 L_t}{m_a^2} \right. \right]$$  

(2.12)
where the pressure drop is corrected for the momentum change. The cross-sectional area is that of the empty tube without considering inserts. Also, in calculating the Reynolds number, no allowance was made for the blockage of the inserts. A sample calculation and the uncertainty analysis are given in Appendix C.

C. Experimental Program

1. Empty tube experiments

Before starting to evaluate the performance of the inserts, the experimental apparatus was tested for its performance as a plain tube without inserts. Average heat-transfer coefficients were determined for each calorimeter segment of the test section. The average friction factor was calculated by using the pressure drop across the test section.

The conditions just upstream of the test section approximate an almost uniform temperature gas with a fully developed velocity profile. At the cooled test section, a temperature profile develops and the heat-transfer coefficient in the first segment should be that for a developing flow. Further downstream, the temperature profile should become fully developed. This behavior is shown in the top half of Figure 2.7 where the data for segment 1 is above that for the other three segments. The average Nusselt numbers, corrected for thermal entry length effects according to curves given by Mills (1962) and property variations, are shown in the bottom half of Figure 2.7 together
FIGURE 2.7. Heat-transfer data for smooth tube compared with the Gnielinski correlation.
with a line representing the Gnielinski (1976) correlation for cooling expressed as

$$\frac{\text{Nu}_D \left( \frac{T_S}{T_b} \right)^{0.45}}{\text{Nu}_D} = \frac{\left( f/2 \right) \left( \text{Re}_D - 1000 \right) \text{Pr}}{1 + 12.7 \sqrt{\left( f/2 \right) \left( \text{Pr}^{2/3} - 1 \right)}}$$

where $T_S$ is the average surface temperature of the inside of the tube, obtained from the average of the six wall thermocouple readings at the ends of each segment. Shah and Johnson (1981) recommend the Gnielinski correlation for constant property fully developed turbulent flow through a smooth circular tube for Reynolds numbers ranging from 2300 to 5,000,000. To correct for property variations in cooling of air, Gnielinski recommends an exponent of -0.45 for use with his correlation.

The average Fanning friction-factor data, compared with the Filonenko (1954) correlation

$$f = \left[ 1.58 \ln(\text{Re}_D) - 3.28 \right]^{-2}$$

are shown in Figure 2.8 where the agreement is within ±2%. Property variation effects for the friction data have not been considered because of the uncertainty in the correction factor (Shah and Johnson, 1981). Figures 2.7 and 2.8 show that the heat-transfer data are within ±5% of the Gnielinski correlation and the friction data are within ±2% of the Filonenko correlation. This confirms the basic smooth tube performance of the test apparatus.
FIGURE 2.8. Friction-factor data for smooth tube compared with the Filonenko correlation
2. Basic thermal hydraulic experiments

a. Inserts tested  The type of insert tested is shown in Figure 2.9. Heat-transfer and pressure-drop studies were conducted with eleven different geometrical variations of the same type of insert. Included were pitches of 6.0, 9.5, and 16.0 in., and insert widths of 0.5, 0.75, and 0.875 in., as shown in Table 2.1. Three of the 9.5 in. pitch inserts were made by a commercial manufacturer. Included was one insert fabricated from perforated steel. The other eight inserts were made at Iowa State University (ISU). One of the ISU inserts had a twist angle of 0°, while all the other inserts had a twist angle of 25°. Equations relating tube diameter, insert pitch, width, longitudinal angle and twist angle are given in Appendix D.

b. Results and discussion  Nusselt numbers obtained from the tube side heat-transfer coefficients of the four segments are plotted against Reynolds numbers in Figures 2.10-2.20. The Nusselt numbers were corrected for variable properties by using the exponent -0.45 in order to compare the data with the empty tube correlation of Gnielinski (1976). The pressure-drop data, converted to the Fanning friction factor, are also plotted in these figures. Heat-transfer and pressure-drop results are given in a tabular form in Appendix E.

The heat-transfer data in Figures 2.10-2.20 show that segment 1 heat transfer is lower than the other three segments, which indicates possible thermal hydraulic entrance effects. In the following discussion, only results from the second, third, and fourth segments
FIGURE 2.9. Insert geometry (not to scale)
TABLE 2.1. Insert characteristics

<table>
<thead>
<tr>
<th>Insert Label</th>
<th>Pitch P (in.)</th>
<th>Width W (in.)</th>
<th>Long. Angle $\phi$ (°)</th>
<th>Twist Angle $\gamma$ (°)</th>
<th>Manufacturer</th>
</tr>
</thead>
<tbody>
<tr>
<td>A1</td>
<td>9.5</td>
<td>0.875</td>
<td>122</td>
<td>25</td>
<td>Commercial</td>
</tr>
<tr>
<td>AN</td>
<td>9.5</td>
<td>0.75</td>
<td>122</td>
<td>25</td>
<td>Commercial</td>
</tr>
<tr>
<td>A2</td>
<td>16.0</td>
<td>0.875</td>
<td>143</td>
<td>25</td>
<td>ISU</td>
</tr>
<tr>
<td>A5</td>
<td>16.0</td>
<td>0.75</td>
<td>143</td>
<td>25</td>
<td>ISU</td>
</tr>
<tr>
<td>A4</td>
<td>6.0</td>
<td>0.875</td>
<td>90</td>
<td>25</td>
<td>ISU</td>
</tr>
<tr>
<td>A3</td>
<td>6.0</td>
<td>0.75</td>
<td>90</td>
<td>25</td>
<td>ISU</td>
</tr>
<tr>
<td>A6</td>
<td>9.5</td>
<td>0.875</td>
<td>122</td>
<td>25</td>
<td>ISU</td>
</tr>
<tr>
<td>A7</td>
<td>9.5</td>
<td>0.75</td>
<td>122</td>
<td>25</td>
<td>ISU</td>
</tr>
<tr>
<td>A8</td>
<td>9.5</td>
<td>0.5</td>
<td>122</td>
<td>25</td>
<td>ISU</td>
</tr>
<tr>
<td>A9</td>
<td>9.5</td>
<td>0.75</td>
<td>122</td>
<td>0</td>
<td>ISU</td>
</tr>
<tr>
<td>PA</td>
<td>9.5</td>
<td>0.75</td>
<td>122</td>
<td>25</td>
<td>Commercial</td>
</tr>
<tr>
<td></td>
<td></td>
<td></td>
<td></td>
<td></td>
<td>(Perforated)</td>
</tr>
</tbody>
</table>

have been considered; the entrance effects are considered in a later section. Figures 2.21 and 2.22 show the effect of width and fabrication technique on heat-transfer and pressure-drop data, respectively, by using the data for the inserts A6, A7, A8, A1, and AN. These inserts all have the same pitch of 9.5 in. but the widths vary as shown in the figures. From Figures 2.21 and 2.22, it is seen that the heat-transfer coefficient and friction factor decrease with decreasing width. It is also seen that the commercially manufactured inserts A1 and AN have higher heat transfer and pressure drop. Inserts A1 and A7 (also AN and A7) are identical in most dimensions, except that inserts A6 and A7 have sharp bends (radius of bend about 0.1 in.) and inserts A1 and AN have rounder bends (radius of bend about 1 in.). A least-squares line is
FIGURE 2.10. Heat-transfer and pressure-drop data for insert A1
Figure 2.11. Heat-transfer and pressure-drop data for insert AN.
FIGURE 2.12. Heat-transfer and pressure-drop data for insert A2
FIGURE 2.13. Heat-transfer and pressure-drop data for insert A5
FIGURE 2.14. Heat-transfer and pressure-drop data for insert A4
FIGURE 2.15. Heat-transfer and pressure-drop data for insert A3
FIGURE 2.16. Heat-transfer and pressure-drop data for insert A6
FIGURE 2.17. Heat-transfer and pressure-drop data for insert A7
FIGURE 2.18. Heat-transfer and pressure-drop data for insert A8
FIGURE 2.19. Heat-transfer and pressure-drop data for insert A9
FIGURE 2.20. Heat-transfer and pressure-drop data for insert PA
FIGURE 2.21. Effect of insert width and fabrication technique on heat-transfer coefficient
FIGURE 2.22. Effect of insert width and fabrication technique on friction factor
drawn through the data for each insert to give a better idea of the enhancement above the smooth tube base line.

Figures 2.23 and 2.24 show the heat-transfer and pressure-drop data for inserts A3, A4, A2, and A5. In addition to the effect of width, Figures 2.23 and 2.24 also show that heat transfer and pressure drop increases with decreasing pitch. This trend is shown more clearly in Figures 2.25 and 2.26, in which heat-transfer data of inserts A6, A7, and A2 (all of width 0.875 in.) and of A7, A3, and A5 (all of width 0.75 in.) are plotted, respectively, as a function of Reynolds number. For example, the data in Figures 2.21 and 2.23 show a 6% increase in Nusselt number for an increase of 17% in width. Similarly, Figure 2.25 indicates a 15% increase in Nusselt number for a 37% decrease in pitch. The same trend of an increase in friction factor with a decrease in pitch is also shown in Figures 2.27 and 2.28.

The performance of the perforated insert PA (identical to insert AN except for the perforations) is shown in Figures 2.29 and 2.30. Insert PA has a much lower heat transfer and pressure drop than insert AN. To estimate the effect of twist angle, the performances of inserts A7 and A9 are compared in Figures 2.31 and 2.32. With lower twist angles, the flow encounters a larger projected cross section of the insert, and therefore, an increase in heat transfer and pressure drop is expected.

3. Insert entrance length experiments

The flow conditions entering the test section approximate a fully developed turbulent flow velocity profile with a temperature
FIGURE 2.23. Effect of insert width and pitch on heat-transfer coefficient
FIGURE 2.24. Effect of insert width and pitch on friction factor
FIGURE 2.25. Effect of insert pitch on heat-transfer coefficient. Insert width = 0.875 in.
FIGURE 2.26. Effect of insert pitch on heat-transfer coefficient. Insert width = 0.75 in.
FIGURE 2.27. Effect of insert pitch on friction factor. Insert width = 0.875 in.
FIGURE 2.28. Effect of insert pitch on friction factor. Insert width = 0.75 in.
FIGURE 2.29. Effect of insert perforations on heat-transfer coefficient.
FIGURE 2.30. Effect of insert perforations on friction factor
FIGURE 2.31. Effect of insert twist angle on heat-transfer coefficient
FIGURE 2.32. Effect of insert twist angle on friction factor
distribution characteristic of a long, insulated pipe. As mentioned before, the heat-transfer data from the first segment of the test section are generally lower on the coordinates of Nusselt number versus Reynolds number than the data from the other segments because fully developed enhanced conditions have not been attained in the first segment. Also, it is found that the insert length-to-diameter ratios in most insert applications are rather short. Hence, a study of the conditions in this entrance region is appropriate in order to assess the insert length required to attain fully developed enhanced conditions.

a. Method of testing for insert entrance effects The study of the entrance region heat transfer for an insert was conducted using the same experimental test rig. Figure 2.33 shows how the insert was moved axially in the tube so its upstream end could be placed at any desired location with respect to the first and second calorimeter segments. By moving the insert to a given position, the location where the enhancement begins is known in relation to the instrumented segments. The "insert entrance length" for the any particular segment, $L_i$, is defined as the distance from the upstream end of the insert to the upstream end of segment $i$. For example, the insert entrance length for the second segment, $L_2$, is defined as the distance from the upstream end of the insert to the upstream end of segment 2.

Experiments were conducted for Reynolds numbers of 3000, 5000, 10,000, 15,000, and 20,000, for eight insert positions and for two inserts. Inserts A4 and A6, which were of different pitch but same
FIGURE 2.33. Location of insert to obtain entry region data
width, were tested because a pitch effect of the insert entrance length was anticipated.

b. Results and discussion The entrance region heat-transfer results are shown in Figures 2.34 and 2.35, where Nusselt numbers for segments 1, 2, and 3 are plotted against the ratio of insert entrance length to insert pitch. The figures indicate that as the ratio of insert entrance length to pitch is reduced, the heat-transfer coefficient from the hot gas to the tube decreases almost linearly. When the insert is absent from a particular segment, the Nusselt numbers correspond to the smooth tube condition for the particular Reynolds number. Both inserts indicate that a fully developed enhanced heat-transfer condition is attained at an insert entrance length to pitch ratio of about 1.5. After the fully developed enhanced condition is reached, the Nusselt number should not vary with the insert entrance length. However, the figures show that there are variations in the data at insert entrance-length-to-pitch ratios greater than 1.5 and that the fluctuations are more pronounced at higher Reynolds numbers. This could be caused by a variation in the number of contact points between the insert and the tube wall in the particular segment, which in turn depends on the location of the insert with respect to the segment.

4. Experiments to differentiate the effects of core and wall regions

The flow visualization studies on inserts reported in Chapter III indicate that the flow at the core of the tube and the flow at the wall do not mix as well as intuition suggests. Hence, it was deemed
FIGURE 2.34. Insert entrance length heat-transfer data for insert A6
FIGURE 2.35. Insert entrance length heat-transfer data for insert A4
necessary to study the effects of the core region and the wall region separately.

a. Inserts tested  To facilitate this study, the bent-strip insert A1 was cut apart to provide separate wall and core inserts AW and AC as shown in Figure 2.36. First, 1/8-in. diameter holes were drilled at three different axial locations in each of the insert sides, and three rigid 1/8-in. wires were inserted through the holes and welded to the insert. Then the insert was cut axially into three sections by using a band saw. Spacers were welded to keep the two wall sections in their correct locations, thereby forming the insert AW. Similarly, spacers were welded to keep the core section in the center of the tube to make the core region insert AC. The core section insert AC had a width of 2/3 times the diameter. These two inserts were separately tested for heat transfer and pressure drop.

b. Results and discussion  Figures 2.37 and 2.38 give the heat-transfer and pressure-drop data for the core region insert AC and the wall region insert AW. These data are compared with insert A1, the original insert from which the two sectional inserts were made. Figure 2.37 shows that the core region insert enhances the heat transfer more than the wall region insert. Although the flow visualization studies in Chapter III of this study indicate that the flow along the wall is not much disturbed by the entire insert, the results of this chapter indicate that the core region plays a significant, if not dominant, role in enhancing the heat transfer. However, it is likely that the core
FIGURE 2.36. Inserts AW and AC. Fabricated to test core and wall region separately.
FIGURE 2.37. Heat-transfer data for the core (AC) and wall (AW) region inserts compared with data for insert A1
FIGURE 2.38. Pressure-drop data for the core (AC) and wall (AW) region inserts compared with data for insert A1
insert performs better on its own than the core region of the full insert because of the eddies shed from the noncontinuous sharp-edged strips. Figure 2.38 also shows the same sort of results for the case of the friction factor, with insert AC exhibiting nearly the same friction factors as the full insert.

D. General Correlations and Performance Evaluations

1. Empirical correlation of data

A statistical approach (Ostle and Mensing, 1975) was used to find a correlation by assuming a model and minimizing the least squares differences in the dependent variables to fit all the data that had been obtained for the unperforated inserts. The data of insert A9 (with twist angle 0°) was also not considered in the correlation. The following equation models were used for the correlations:

\[
\text{Nu}_D = A \left( \text{Re}_D \right)^B \left( \frac{L}{D_1} \right)^C \left( \frac{W}{D_1} \right)^D \left( \frac{T_s}{T_b} \right)^{-0.45} \tag{2.15}
\]

\[
f = A \left( \text{Re}_D \right)^B \left( \frac{L}{D_1} \right)^C \left( \frac{W}{D_1} \right)^D \tag{2.16}
\]

where A, B, C, and D are constants. In the above equations, \(\text{Nu}_D\) is based on sectional-average heat transfer, while \(f\) is based on pressure drop across the whole test section. In Equation 2.15, the Prandtl number is not considered in the correlation because there is no appreciable variation in Pr in the experiments conducted. Therefore, the correlation developed is valid only at a Pr of 0.7 (i.e., for most
gases). A property variation correction was not considered in the correlations because the variation of temperature ratio, $T_s/T_b$, was not large enough. However, the heat-transfer correlation was corrected for bulk properties by using the exponent suggested by Gnielinski (1976).

To obtain the values of the constants, a matrix approach to multiple linear regression was used. The methodology of this approach is given in Appendix F.

Due to the apparent entry length behavior, two equations of the same model were necessary to predict the heat transfer. One equation is used to predict the heat transfer in the "entry region," which is defined as the region from the upstream end of the insert to 1.5 times the insert pitch. The second equation predicts the heat transfer in the "fully developed region," which is defined as the region after an insert length of 1.5 times the insert pitch. The estimates of the constants, exponents and limits at 95% confidence levels are given in Table 2.2. $R^2$ is the correlation coefficient for each of the regression equations. The values of A, B, C, and D are valid in the following range of variables:

\[
\begin{align*}
Re_D &= 3000 - 20,000 \\
Pr &= 0.7 \\
L/D_i &= 2 - 6 \\
W/D_i &= 0.15 - 0.40
\end{align*}
\]

Figures 2.39 and 2.40 show the correspondence between predicted and
TABLE 2.2. Estimates and uncertainties of constants in correlation equations

<table>
<thead>
<tr>
<th>Equation</th>
<th>$\text{Nu}_D$ (Region)</th>
<th>$\text{Nu}_D$ (Entry)</th>
<th>$\text{Nu}_D$ (Developed)</th>
<th>$f$ (Full)</th>
</tr>
</thead>
<tbody>
<tr>
<td>A</td>
<td>0.222 ±0.0045</td>
<td>0.540 ±0.0029</td>
<td>4.350 ±0.0996</td>
<td></td>
</tr>
<tr>
<td></td>
<td>-0.0044</td>
<td>-0.0029</td>
<td>-0.0073</td>
<td></td>
</tr>
<tr>
<td>B</td>
<td>0.695 ±0.0004</td>
<td>0.649 ±0.0001</td>
<td>-0.100 ±0.0004</td>
<td></td>
</tr>
<tr>
<td>C</td>
<td>-0.347 ±0.0009</td>
<td>-0.391 ±0.0003</td>
<td>-1.286 ±0.0011</td>
<td></td>
</tr>
<tr>
<td>D</td>
<td>0.319 ±0.0031</td>
<td>0.503 ±0.0007</td>
<td>1.320 ±0.0030</td>
<td></td>
</tr>
<tr>
<td>$R^2$</td>
<td>0.9551</td>
<td>0.9714</td>
<td>0.9984</td>
<td></td>
</tr>
</tbody>
</table>

Experimental heat-transfer data for the entry region and developed region, respectively. Figure 2.41 shows the correspondence between the predicted and experimental pressure-drop data. The correlations predict most data within ±10%. The correlation obtained above was tested by applying it to data for insert B, shown in Figure 2.42, which was tested by Junkhan et al. (1985). Insert B was made by another commercial manufacturer and is bent in a different fashion from the other inserts in the present study. Figures 2.43, 2.44, and 2.45 compare the experimental heat-transfer and friction-factor data of insert B with the correlations. The friction-factor prediction is excellent, but the correlation overpredicts the heat transfer by about 10% in the developed region and underpredicts the heat transfer in the entry region by about 10%. This could be due to the fashion the insert B is bent at the tube wall as shown in Figure 2.42.
FIGURE 2.39. Predicted data for the entry region Nusselt number (Equation 2.15) compared with the experimental data for 9 inserts.
FIGURE 2.40. Predicted data for the developed region Nusselt number (Equation 2.15) compared with the experimental data for 9 inserts
FIGURE 2.41. Predicted friction-factor data (Equation 2.16) compared with the experimental data for 9 inserts.
FIGURE 2.42. Geometry of insert B tested by Junkhan et al. (1985)
FIGURE 2.43. Insert B entry region heat-transfer data compared with prediction.
FIGURE 2.44. Insert B developed region heat-transfer data compared with prediction
FIGURE 2.45. Insert B friction data compared with prediction
2. **Performance evaluation**

The results presented earlier show that bent-strip inserts augment heat transfer as much as 200% or more. However, the pressure-drop increase accompanying the enhancement in heat transfer is considerably higher. Figures 2.46 and 2.47 show the trends of heat-transfer enhancement and the increase in pressure drop with the variation in width and pitch of the insert. In these figures, the reference Nusselt number, $N_u_s$, is obtained by using the Gnielinski correlation at the same Reynolds number and a Prandtl number of 0.7. Figure 2.46, for a constant $Re_D = 10000$ and $L/D_i = 4.0$, shows a very steep increase in the friction factor but only a marginal increase in heat transfer as the width-to-tube-diameter ratio increases. Figure 2.47, again drawn at a constant $Re_D = 10000$ and $W/D_i = 0.4$, shows a drastic increase in the friction factor and a marginal increase in heat transfer as the pitch-to-diameter ratio is reduced. At first glance, these results indicate that using bent-strip inserts may not be considered practical. However, to actually evaluate and compare the performance of the bent-strip inserts, one must use one or more performance evaluation criteria (PEC) as described by Bergles et al. (1974).

Two performance evaluation criteria are applicable in enhancing convective heat transfer from hot gases in tubes using bent-strip inserts. They are:

1. Increasing heat duty with the basic geometry and pressure drop fixed ($R_2$) and
FIGURE 2.46. Predicted heat-transfer and friction-factor enhancement as a function of width-to-tube-diameter ratio
FIGURE 2.47. Predicted heat-transfer and friction-factor enhancement as a function of pitch-to-tube-diameter ratio
2. Increasing heat duty with the basic geometry and pumping power constant ($R_3$).

The notations for the criteria are the same used by Bergles et al. (1974). In the case of fire-tube boiler applications, improved efficiency is attained by reducing stack temperature and increasing heat duty without having a very drastic increase in pressure drop (Yaverbaum, 1979). Therefore, $R_2$ would be more appropriate in fire-tube boiler applications. In forced draft systems, the criterion $R_3$ would be more appropriate.

The criterion $R_2$, defined as the increase in heat duty for constant pressure drop and basic geometry, can be written as

$$R_2 = \left[ \frac{h_a}{h_o} \right] \left( \Delta p, \Delta T, l, D_i, N, T_{ai} \right) \frac{q_a}{q_o} = \frac{\text{Nu}_a}{\text{Nu}_o}$$  \hspace{1cm} (2.17)

The pressure drop is

$$\Delta p = 4f \left( \frac{L}{D_i} \right) \rho_a \frac{u_m^2}{2} = \text{constant}$$  \hspace{1cm} (2.18)

which becomes

$$f_a \frac{\text{Re}_a^2}{\text{Re}_o^2} = f_o \frac{\text{Re}_o^2}{\text{Re}_o^2}$$  \hspace{1cm} (2.19)

as the constraint for PEC $R_2$.

For a given $Re_a$ and insert geometry, $f_a$ was calculated by using the empirical correlation (Equation 2.14) developed in the previous section. The Filonenko (1954) correlation
\[ f_o = \left[ 1.58 \ln(Re_D) - 3.28 \right]^{-2} \quad (2.14) \]

was used in Equation 2.19 to obtain an equivalent \( Re_o \). A trial-and-error method was necessary to solve for the \( Re_o \) that satisfies the constraint Equation 2.19.

The equivalent \( Nu_o \) was obtained using the Gnielinski (1976) correlation

\[ \left( \frac{T_s}{T_b} \right)^{0.45} = \frac{(f/2)(Re_D - 1000)Pr}{1 + 12.7\sqrt{(f/2)(Pr^{2/3} - 1)}} \quad (2.13) \]

with \( Pr = 0.7 \). Finally, PEC \( R_2 \) was evaluated by obtaining the ratio \( Nu_a/Nu_o \). \( Nu_a \) was calculated by using the empirical correlation (Equation 2.15) developed earlier for the given \( Re_a \) and insert geometry.

The \( R_3 \) performance evaluation criterion, which is defined as the ratio of heat duty under constant pumping power and geometry conditions, is

\[ R_3 = \frac{h_a}{h_o} \left( \frac{q_a}{q_o} \right) = \frac{Nu_a}{Nu_o} \quad (2.20) \]

Pumping power is given by

\[ P = Nu_m A_x f D_i \rho_m u_m^2/2 = \text{constant} \quad (2.21) \]

which becomes

\[ A_x f_a Re_a^3 = A_x f_o Re_o^3 \quad (2.22) \]
This is the constraint equation for the $R_3$ criterion. The same method described earlier was used to calculate an equivalent $Re_o$, and $Nu_a/Nu_o$ was evaluated for a given $Re_a$ and insert geometry.

The performance evaluation results were obtained by using the correlations given for the developed region of the flow. Note that the PEC data are evaluated under idealized conditions where temperature changes along the tube and thermal resistance other than that of the gas-side have been disregarded; however, they demonstrate the qualitative advantages of the use of bent-strip inserts.

The optimum performance condition (flow and/or insert geometry) for the $R_2$ criterion was estimated by plotting PEC $R_2$ against Reynolds number $Re_a$ for three different $L/D_i$ and $W/D_i$ ratios. For a constant $W/D_i = 0.3$, three values of $L/D_i$, 2.0, 4.0, and 6.0 were considered in Figure 2.48. In Figure 2.49, three values of $W/D_i$, 0.2, 0.3, and 0.4, were considered for a fixed value of $L/D_i = 4.0$. From both of these figures it is apparent that $R_2$ is strongly dependent on the Reynolds number. Also, for the constant pressure-drop criterion, the advantages of the bent-strip inserts are found only at lower Reynolds numbers. Decreasing pitch or increasing width reduces the enhancement effectiveness. Moreover, the variation in pitch seems to have a larger impact on the effectiveness than the variation in width.

To estimate the optimum performance conditions according to the constant pumping power ($R_3$) criterion, the variation of $R_3$ with Reynolds number, $W/D_i$, and $L/D_i$ is plotted in Figures 2.50 and 2.51. Again, $R_3$
FIGURE 2.48. Effect of Reynolds number and $W/D_1$ on $R_2$. 

**DEVELOPED REGION**

$\frac{L}{D_1} = 4.0$

$\frac{W}{D_1}$

- 0.2
- 0.3
- 0.4

INCREASING $\frac{W}{D_1}$

$R_2 = \frac{N_{UA}}{N_{U_0}}$

**REYNOLDS NUMBER, $Re_a$**

3000 5000 10000 30000

0.5 0.6 0.7 0.8 0.9 1.0 1.1 1.2 1.3 1.4 1.5 1.6 1.7
FIGURE 2.49. Effect of Reynolds number and $L/D_1$ on $R_2$. 
shows a strong dependence on the Reynolds number, but with the use of this criterion, a much larger Reynolds number range (3000-30,000) is available for favorable applications of the bent-strip insert. Also, the enhancement effectiveness increases with increases in width; however, the variation in pitch has hardly any effect on the enhancement effectiveness.

Although the $R_2$ and $R_3$ PEC criteria yield encouraging results, the actual improvement is likely to be much lower. In the case of a computer simulation of a commercial heating boiler by Didion and Chern (1984) that used the turbulator results obtained by Junkhan et al. (1982), the results indicate that the bent-strip inserts increase the boiler seasonal efficiency by about 6%. These results are rather disappointing considering the high heat-transfer enhancement obtained in this study. One must realize, however, that in the computer simulation the inserts were placed only in 60% of the last pass due to practical limitations, and also that the majority of the heat transfer occurs in the fire box of the fire-tube boiler. Yaverbaum (1979) indicates that when inserts are used, allowances have to be made to ensure sufficient air flow for complete combustion. Yaverbaum also indicates that after inserts are installed in boilers, the burner control needs to be readjusted to obtain optimum efficiency.
FIGURE 2.50. Effect of Reynolds number and $W/D_1$ on $R_3$
FIGURE 2.51. Effect of Reynolds number and L/D₁ on R₃
III. FLOW VISUALIZATION

A. General

Flow visualization is considered one of the most effective methods to study and comprehend the basic physical nature of any flow condition. To better understand the enhancement mechanism and to study the flow patterns created by the bent-strip insert, two techniques of flow visualization were used. One technique consisted of injecting a colored dye in water flowing in a test apparatus and of photographing the dye motion (Merzkirch, 1974; Yang, 1977). The other technique was based on the observation of flow patterns on the wall of a plastic tube. The patterns of fluid flow adjacent to the tube wall were made visible by coating the wall with an oil/lamp-black mixture and using air as the working fluid (Tien, 1978; Tien and Sparrow, 1979; Motwani, 1984).

B. Flow Visualization Experiments

1. Water/dye-injection technique

The apparatus to view the dye motion in water flow consisted of a constant-head water tank which fed water through a smooth nozzle into a 5 ft long, 0.75-in. diameter clear plastic tube, as shown in Figure 3.1. The tank was divided into two sections by a wall of ceramic spheres, which acted as a baffle. The baffle assured an undisturbed water supply to the plastic tube. It was possible to obtain a Re_D range of 100 to 3000 by varying the flow rate using the needle valve at the exit. The flow rate was determined using a weigh tank.
FIGURE 3.1. Water/dye flow visualization setup
A bent-strip insert similar to Figure 2.9 was made of plastic to fit the tube. The dye, a mixture of concentrated food coloring powder and water, was fed into the water by gravity with the flow rate adjusted by a restrictor to match the dye and water velocities. The dye was inserted by means of hypodermic needle tubing near the center of the tube upstream of the insert and at the wall. When fed at the wall, the dye was inserted either at the side facing the viewer or at the top. In some experiments, two different color dyes were used in water to study the flow in the center of the tube and at the wall simultaneously.

Still photographs of the flow pattern were taken with a camera focused on the area of interest in the tube and with a fluorescent tube as a light source. Photographs were taken with 400 ASA color print film at shutter speeds ranging from $1/125$ to $1/1000$ second. Photographs could be obtained in either a vertical or horizontal plane. In order to study variations in the flow patterns as a function of time, the dye motions were videotaped with a portable Panasonic video cassette recorder. The videotape could be played back at a slower speed or stopped to look at single frame snapshots.

2. Oil/lamp-black technique

Visualization of the flow patterns adjacent to the wall was also accomplished by using the oil/lamp-black technique. The heat transfer experimental setup described in Chapter II was used with some modifications. Initially, the calorimeter test section was removed and replaced with a 2.70-in. diameter hard glass tube, and unheated air was
used as the working fluid. The flow-visualization technique consisted of evenly coating half the inside of the glass tube with a thin film of a mixture of oil and lamp black. Although, several types of oil with various weight ratios of oil to lamp black were tried, the best results were obtained with a rust removing oil in a weight ratio of oil to lamp black of about 4:1. The excess oil/lamp black was removed once the surface had been evenly coated. The assembly was immediately installed in the flow system, the blower was started, and the air was allowed to flow through the tube for about 20 minutes. $Re_D$ was adjusted to about 25,000. Under the action of the shear stress exerted by the air on the tube wall, the mixture is made to flow along the wall in the flow direction. The resulting black and white pattern of streaks represent the average flow field near the surface. In a stagnation region, the oil and lamp black mixture remains as it was initially applied, while a region clear of the mixture indicates a region of high shear stress. The flow patterns thus formed on the glass were photographed.

To enhance the photographic results, the author used a different scheme. In this case, a 2.75 in. i.d. plastic tube was used instead of the glass tube. A second, slightly smaller diameter plastic tube, 2.5 in. i.d. by 2.75 in. o.d. and 5 ft in length, was split longitudinally into two halves. White, plasticized adhesive paper was secured to the inner surfaces of the two halves. First, the insert was placed between the two half tubes which were then reassembled and inserted in the 2.75 in. i.d. plastic tube. Next, the inner white
surface of the reassembled two halves was coated with a mixture of oil and lamp-black. Finally, the entire assembly was installed in the flow system. The air was allowed to flow through the test section for about 20 minutes. The assembly was then taken out of the flow system, and the split halves were removed from the outer tube so that photographs could be taken of the flow patterns formed on the white contact paper. Black and white 400 ASA film was used to photograph these patterns.

C. Observations of Flow Phenomena

Figure 3.2(a) shows one of the photographs obtained from the dye-injection apparatus with the camera lens axis in a horizontal plane and with dye injection in the center of the tube. The flow pattern in Figure 3.2(a) shows that the core flow at a Reynolds number of 1667 begins to be substantially mixed at the core of the tube within two insert pitch lengths.

Figure 3.2(b) shows the flow patterns when the dye is inserted at the wall with the same camera orientation and at a Reynolds number of 1667. Note that the wall flow does not appear to be affected to a great extent by the bent-strip insert at this Reynolds number. This observation is somewhat contradictory to the thermal-hydraulic results obtained by using the separate core and wall region inserts. The core-region insert tests indicate that the core region of the insert disturbs the wall flow even more than the wall region insert. Figure 3.2(c) shows a view obtained with the camera lens axis in a vertical plane with
FIGURE 3.2. Flow patterns obtained by using the water/dye-injection flow visualization technique. $Re_D = 1667$
the dye injection from the top, and with the strip in the vertical orientation. Here, the wall layer at the top gets disturbed by the bluff-body-like flow downstream of the strip where it is in contact with the tube wall.

Simultaneous insertion of two colors of dye in Figure 3.2(d) shows these flow patterns at the same time. In this case, the camera lens axis was in the horizontal plane. The blue dye was inserted at the center of the tube ahead of the insert while the red dye was inserted at the wall just ahead of the insert. The blue dye at the center was fully mixed almost immediately because of the insert while the red dye required about two pitch diameters from the beginning of the insert to mix with the core. This suggests that an entrance effect exists and that the earlier assumption of the flow being fully developed by the end of the first segment (1-2 pitch lengths) in the heat-transfer/pressure-drop experiment is valid.

As mentioned before, the flow patterns obtained by the dye injection technique were videotaped. Some observations that could not be photographed but that are seen by re-playing the videotape include:

- The flow at the wall spirals along the surface, indicating the presence of secondary swirl flow.
- The flow at the wall and the flow at the core do not mix as rapidly as anticipated in 1-2 pitch lengths.
- By injecting the dye at the wall upstream of the insert, it is seen that the bent-strip insert affects the wall layer after a
length of 1-2 insert pitches. This indicates the presence of a developing length.

• At high turbulent core velocities, the wall layer remains comparatively unmixed and very laminar at locations upstream of the insert. However, at the same flow rate and in the presence of the insert, the sublayer at the wall appears to very turbulent and chaotic. This behavior is also observed away from the insert contact with the tube wall.

• It is possible to see streaks of dye from the wall layer being lifted away into the core region at random intervals. The frequency of lifting appears to increase at higher velocities.

Figures 3.3, 3.4, 3.5, and 3.6 show the flow patterns obtained by using the oil/lamp-black technique with the glass tube. These flow patterns are at the wall around the point of contact between the tube wall and the bent-strip insert. In all four figures, the plane passing through the center of the insert is in the vertical orientation. Figures 3.3 and 3.5 were photographed in the horizontal direction through the uncoated side of the glass tube. Figures 3.4 and 3.6 shows the same view as Figures 3.3 and 3.5 from below the glass tube.

Figures 3.7(a), 3.7(b), and 3.7(c) are flow patterns obtained at a Reynolds number of 25,000 by using the oil/lamp-black technique. They show, respectively, the first contact point between the wall and the insert, the space between the first and second contact points, and the
FIGURE 3.3. Flow patterns obtained by using the oil/lamp-black technique with the glass tube. Photographed from the front.
FIGURE 3.4. Flow patterns obtained by using the oil/lamp-black technique with the glass tube. Photographed from under the tube.
FIGURE 3.5. Flow patterns obtained by using the oil/lamp-black technique with the glass tube. Photographed from the front.
FIGURE 3.6. Flow patterns obtained by using the oil/lamp-black technique with the glass tube. Photographed from under the tube.
second contact point of the insert and the tube wall. The tube center line and the approximate location of the contact points of the tube wall and the insert—noted by the white rectangle—have been marked on the figures. Note that the insert touches the tube wall only at the two extremities and that there is thin gap between the insert and the tube wall at the contact area.

The photographs in Figures 3.3 through 3.7 show that most flow disturbance patterns occur near the contact points of the bent-strip turbulator and the tube wall. Since the flow patterns are caused by shear stresses at the wall, the patterns indicate that the shear stresses near the contact points of the inserts and the tube are much higher than at other regions of the wall. The flow patterns also indicate that the flow at the wall gets turned in the direction of the twist of the insert and that the effect of the insert appears at a distance of about 30% of pitch upstream of the contact point. After the contact point, the flow returns back to its original axial position within a downstream distance of about 15% of pitch. In the radial direction, the insert is seen to affect the flow pattern to about 30° from the plane of the insert. Further, it can be observed that between the two contact points, the flow at the wall is not greatly affected by the tape. That is, the flow on one side of the wall does not "feel" the effects of the insert contacting the wall on the diametrically opposite side.
FIGURE 3.7. Flow patterns obtained by using the oil/lamp-black technique with the split plastic tube
In conclusion, both flow visualization techniques show that the flow at the wall is affected most where the bent strip insert comes in contact with the tube wall. Nevertheless from the water/dye-injection technique, it appears that in the regions where the bent strip does not touch the tube wall, the wall sublayer tends to be turbulent and very choppy. Also, mixing of the core and wall regions seems to occur when the sublayer gets ejected into the core region. The observation of the streaks of dye being ejected from the wall into the core region underscores the possibility of using a surface renewal/penetration theory to model the heat-transfer process when it is enhanced by bent-strip inserts.
IV. SURFACE RENEWAL/PENETRATION MODEL OF HEAT TRANSFER

A. Background

The surface renewal/penetration model has been applied to a wide variety of transport problems. Applications include mass transfer across a liquid-gas interface (Higbie, 1935; Danckwerts, 1951), mass transfer across a solid-fluid interface (Hanratty, 1956; Toor and Marchello, 1958; Harriot, 1962; Pinczewski and Sideman, 1974), heat transfer between a surface and a turbulent fluid (Toor and Marchello, 1958; Thomas, 1970a; Thomas and Fan, 1971; Thomas, 1982; Ibrahim and Thomas, 1983), heat transfer in a fluidized bed (Chung et al., 1972), boiling heat transfer (Han and Griffith, 1965; Mikic and Rohsenow, 1969), scraped surface heat transfer (Harriot, 1958; Azoory and Bott, 1970; Hagge and Junkhan, 1975) and condensation heat transfer in a tube with a static mixer (Fan et al., 1978). Most of the above authors have had reasonable success in using the surface renewal/penetration theory to model the heat transfer. In at least one case, pressure drop has also been modelled (Meek and Baer, 1970).

B. Turbulence Burst Phenomena

The turbulent structure near a surface has been an area of active research for a long time. Early experimental studies of Kline et al. (1967), Popovich and Hummel (1967), and Corino and Brodkey (1969) have indicated that flow very near the wall (i.e., in the wall sublayer) is not laminar but is disturbed by small velocity fluctuations.
Kline et al. (1967), using the dye-injection and hydrogen-bubble flow-visualization techniques, have shown that the turbulent structure in the near wall region exhibits recognizable and reproducible patterns—lately known as "coherent structures"—of events. These events consist of a gradual lift up of a layer at the wall (also called "streaks") and a rapid ejection where the low-speed streaks move rapidly away from the wall. In later experiments, Kim et al. (1971) termed this ejection process as "bursting".

Corino and Brodkey (1969) reporting from visual studies of motion pictures of the trajectories of small suspended particles in the flow, substantiated the same idea as Kline et al. (1967) in a slightly different fashion. The sequence of events involved in the turbulent transfer process consists of (i) formation of a low-speed packet of fluid near the wall region; (ii) interaction of a much larger high-speed fluid with the low-speed packet of fluid (also known as "sweep") so that the low-speed packet is accelerated; (iii) a rapid ejection where the low-speed packet moves away from the wall with intense, abrupt, and chaotic movements.

Offen and Kline (1974, 1975) reiterated the work of Kline et al. (1967) and also suggested that the lifting up of the wall layer may be a result of vortex pairing that results from the interaction of the vortices of the previous burst and the lifting up streak. Visual and hot-wire anemometry studies by Head and Bandyopadhyay (1981) indicate that the boundary layer action consists of vortex loops, horseshoes, or
hairpins arrayed at a characteristic angle of approximately 45° to the wall. They also conclude that the shape of the vortices is dependent on the Reynolds number. Some later experiments conducted by Falco (1982) show that the turbulent production process is closely associated with the formation and evolution of a localized flow module called a pocket flow module. Pockets form in the wall region of the turbulent boundary layer as a result of the wall interaction of vortex-ring-like eddies that originate in the outer flow.

Parallel to the flow-visualization experiments, many investigators (Rao et al., 1971; Blackwelder and Kaplan, 1976; Willmarth and Bogar, 1977; Willmarth and Sharma, 1984) have studied the turbulent structure at the wall by using hot-wire anemometry. Results of Rao et al. indicate that the mean observed frequency of bursting is related to the outer flow variables and not with the wall flow variables. Willmarth and Sharma indicate similar results where the burst frequency they measured does not depend on the Reynolds number when scaled with wall parameters. They add, however, that the burst frequency increases with the Reynolds number when scaled with outer variables.

The realities, myths, and present state of the coherent structure of turbulence are surveyed in a recent paper by Hussain (1983). Despite many differences of opinion among experimentalists regarding observed details and methods used to study the coherent structure of wall turbulence, there is considerable agreement on the basic occurrence of events such as low-speed layers and their lift up and on the high
contribution to transport properties by the ejection and sweep phases of bursting. An idealized schematic picture of the turbulent burst process is shown in Figure 4.1. Grass (1971) describes how relatively large packets of low velocity fluid adjacent to the wall surface are periodically ejected from the wall region into the turbulent core region. Satisfaction of continuity requires that the fluid ejected must be replaced by a fresh inrush of fluid from the core having the properties associated with the turbulent core. During the period of time the fluid packet is in the vicinity of the surface, transient molecular transport (or conduction in the case of heat transfer) is said to occur. This basic principle is used to develop a model of heat transfer from gases in tubes with bent-strip inserts.

C. The Surface Renewal Model

The following development of a surface renewal/penetration model for heat transfer in a tube in the presence of bent-strip inserts assumes that

1. A packet of fluid at bulk temperature, driven towards the wall by the insert action and by unsteady turbulent motion, replaces another packet of fluid and comes into contact with the wall.

2. During the contact time, heat is transferred predominantly by transient conduction from the fluid to the wall.
FIGURE 4.1. Turbulent burst process showing action of a fluid packet.
The fluid packet remains in contact with the wall for a certain length of time before it is replaced by another packet of fluid from the core region. During this time, the following equation describes the transient conduction between the wall and the packet of fluid:

$$\frac{\partial T}{\partial t} = \alpha \frac{\partial^2 T}{\partial y^2} \quad (4.1)$$

It is realistic to assume that the temperature of the packet remains at a constant value of $T_i$ beyond a certain characteristic length $\delta$ from the tube wall, which is the thickness of the fluid packet as shown in Figure 4.1. $T_i$ represents the temperature of eddies at first instant of renewal. The initial and boundary conditions describing this transfer process are

$$T = T_i \quad \text{at } t = 0 \text{ and } y > 0$$

$$T = T_s \quad \text{at } t > 0 \text{ and } y = 0 \quad (4.2)$$

$$T = T_x \quad \text{at } t > 0 \text{ and } y = \delta$$

The solution of Equation 4.1 subject to the conditions given by Equation 4.2 is presented in Appendix G. The final result is

$$\frac{T - T_s}{T_i - T_s} = \frac{y}{\delta} - \frac{2}{\pi} \sum_{n=1}^{\infty} \left[ \frac{1}{n} \sin \left( \frac{n\pi y}{\delta} \right) \exp \left( \frac{-an^2\pi^2}{\delta} \right) \right] \quad (4.3)$$

If $q_i(t)$ is the instantaneous heat-transfer rate at the wall at contact time $t$, then
and the mean heat-transfer rate is given by the summation of the heat-transfer contribution by all the packets of fluids of different ages

\[ q = \int_0^\infty q_i(t) P(t) \, dt \]  

(4.5)

where \( P(t) \) is the contact-time distribution function of the fluid packets. Because of the highly mixed flow in the tube caused by the presence of the insert (observed in the flow-visualization experiments), it is reasonable to assume that the fluid packets on the surface are replaced in a completely random manner. Under this assumption, the distribution function according to Danckwerts (1951) is

\[ P(t) = s \exp(-st) \]  

(4.6)

where \( s \) is the frequency of renewal of the fluid packets.

Considering an approach very similar to that followed by Fan et al. (1978), the frequency of renewal of the fluid packets is assumed to be proportional to the number of times the fluid gets "deflected" by the insert in a given length and to the "amount" of fluid that gets deflected by the insert per unit mean residence time \( t_m \) of the fluid in the tube. The sketch of the insert in Figure 2.9 shows that the fluid gets deflected by the insert twice in one pitch \( L \), and therefore, in a given length \( L_s \), the fluid gets deflected \( 2L_s/L \) times. The proportion
of the total flow that gets deflected by the insert should be related to the fraction of the total cross-sectional area that is blocked by the projected area of the insert. This fraction is given by \( \frac{W \cos(\theta) D}{\pi D^2/4} \) which reduces to \( 4W \cos(\theta)/\pi D \). Therefore, the frequency of renewal can be written as

\[
s = C_1 \left[ \frac{2L_x}{L} \left( \frac{4W \cos(\theta)}{\pi D_i} \right) \left( \frac{1}{t_m} \right) \right]
\]

where \( C_1 \) is a proportionality constant.

Combining Equations 4.3 to 4.6 gives the following expression (Appendix F):

\[
q = - \left( T_s - T_i \right) \sqrt{\left( \frac{skpc_p}{\alpha} \right)} \coth \left( \frac{s^2}{\alpha} \right)^{\frac{1}{2}}
\]

For moderate Prandtl-number fluids such as air or gas, \( T_i \) may be replaced by \( T_b \). (This is not a good assumption for low-Prandtl number fluids where the energy transfer during the flight of eddies from the bulk stream to the wall region can be significant (Thomas, 1970b).) Thus, changing \( T_i \) to \( T_b \) and converting the heat-transfer rate to heat-transfer coefficient by using the usual definitions, the following equation is obtained:

\[
h = - \left( \frac{q}{T_b - T_s} \right) = - \left( \frac{q}{T_i - T_s} \right) = \sqrt{\left( \frac{skpc_p}{\alpha} \right)} \coth \left( \frac{s^2}{\alpha} \right)^{\frac{1}{2}}
\]
For sufficiently large values of $\delta$,

$$\coth \left( \frac{s \delta^2}{\alpha} \right)^{\frac{1}{2}} = 1$$

(4.10)

This assumption reduces Equation 4.9 to

$$h = \sqrt{\frac{skpc}{p}}$$

(4.11)

Equation 4.11 is identical to the equations developed by other users of the surface renewal/penetration model.

Substituting Equation 4.7 for $s$ in Equation 4.11 yields

$$h = C_1 \left( \frac{2I_s}{L} \left( \frac{4W}{D_i} \cos(\gamma) \right) \left( \frac{kp\gamma}{t_m} \right) \right)^{\frac{1}{2}}$$

(4.12)

To calculate $h$ using the basic surface renewal model by applying Equation 4.12, the mean residence time of the fluid packets in the tube must be known. Logically, the mean residence time, $t_m$, of the fluid in the tube must be inversely proportional to the mean velocity, $u_m$, of the fluid and proportional to the length of the tube, $L_s$. Thus

$$t_m = C_2 \left( \frac{L_s}{u_m} \right)$$

(4.13)

where $C_2$ is another proportionality constant. Substituting Equation 4.13 into Equation 4.12 and rearranging terms, the expression for the heat-transfer coefficient is obtained as
\begin{equation}
  h = C \left( \frac{8W \cos(\theta)}{L} \right) \left( \frac{kp_c u_m}{D_i} \right)^{\frac{1}{2}}
\end{equation}

where \( C = C_1 C_2 \).

Converting to dimensionless terms (\( \text{Nu}_D = hD_i/k, \text{Re}_D = \rho u_m D_i/\nu, \) and \( \text{Pr} = \mu c_p /k \)) the final form of the model is obtained:

\begin{equation}
  \text{Nu}_D = C \left( \frac{8W \cos(\theta)}{\pi L} \right)^{\frac{1}{2}} \text{Re}_D^{0.5} \text{Pr}^{0.5}
\end{equation}

D. Comparison of Model with Experimental Results and Discussion

Equation 4.15 indicates that an increase in width, \( W \), or a decrease in pitch, \( L \), or twist angle, \( \theta \), results in an increase in heat transfer. These trends compare well with the experimental results presented in Chapter II as shown in Figures 2.21, 2.23, 2.25, 2.26, 2.29, and 2.31. For example, the data in Figures 2.21 and 2.23 show a 6\% increase in Nusselt number for an increase of 17\% in width. Similarly, Figure 2.25 indicates a 15\% increase in Nusselt number for a 37\% decrease in pitch.

Since the constant \( C \) depends on the fluid properties and flow conditions, an exact numerical value for the constant \( C \) could not be obtained analytically by using the surface renewal model. Nevertheless, a linear regression analysis based on Equation 4.15 yielded a value of \( C = 2.06 \). The intercept was forced to go through the origin so as to
fit the model. The experimental results are compared in Figure 4.2 according to Equation 4.15. The line thus obtained is shown on Figure 4.2. The agreement is within ±20%. It is noted, however, that the deviation is greater at higher Nu_D, i.e., higher Re_D.

To show the variation of heat transfer coefficient with insert geometry, the experimental data and prediction are plotted on a logarithmic plot of Nusselt number as a function of Reynolds number in Figures 4.3-4.5. Figures 4.3, 4.4, and 4.5 show the effect of pitch, L, width, W, and twist angle, \( \gamma \), respectively. The figures also indicate that the Reynold number exponent of 0.5 in the model is too low. However, several users of the basic surface renewal model also use a exponent of 0.5 (Thomas and Mohler, 1971). The surface renewal/penetration theory assumes (i) the inrush of fluid penetrates down to the wall itself and (ii) the ejected fluid packets gets completely mixed with the bulk fluid and attains the bulk properties; but as described by Hagge and Junkhan (1974), several investigators argue that these assumptions do not hold in practice. To overcome these shortcomings, several methods—including empirical methods—have been used, where constants are added to the existing exponent of Re_D of Pr. These constants take care of the inaccuracies of the model and are dependent on fluid properties and flow geometry. Moreover, these constants have to be evaluated experimentally. Since the analytical model described by Equation 4.15 is very similar to the empirical correlation given by Equation 2.15 in Chapter II, it was considered not
FIGURE 4.2. Scatter plot of Equation 4.15 with experimental data
necessary to further carry out an empirical evaluation of constants in the analytical model.
FIGURE 4.3. Effect of insert pitch, L, on Nusselt number and comparison with Equation 4.15
FIGURE 4.4. Effect of insert width, $W$, on Nusselt number and comparison with Equation 4.15.
FIGURE 4.5. Effect of insert twist angle, $\gamma$, on Nusselt number and comparison with Equation 4.15
The present study investigated the enhancement of in-tube single phase convective heat transfer in gases by means of turbulence-producing bent-strip inserts. Different geometrical variations of a bent-strip insert were tested to obtain heat-transfer and pressure-drop data. Flow-visualization studies were conducted to view the flow patterns created by the insert and thus try to understand the mechanism of enhancement. Finally, the surface renewal/penetration theory was used to obtain an analytically based model for heat transfer. The following conclusions were drawn:

1. The present investigation is the first comprehensive study conducted on enhancement of gas-side heat transfer by using bent-strip inserts.

2. The heat-transfer and pressure-drop results for bent-strip inserts confirm that there is significant increase in heat transfer; however, a considerable pressure drop penalty is incurred. A maximum heat-transfer enhancement of about 300% was achieved with an accompanying pressure-drop increase of about 1800%. The heat-transfer increases are substantially higher than the enhancement level achieved by using wire-coil or twisted-tape inserts.
3. The performance of bent-strip inserts shows that high enhancement in heat-transfer coefficient cannot be achieved with moderate increase in pressure drop.

4. The thermal-hydraulic experiments indicate that $\text{Nu}_D$ and $f$ are directly dependent on the pitch and width of a bent-strip insert.

5. The experimental investigations reveal an area of lesser level of augmentation at the entrance region of the insert. An insert length of one and one-half times insert pitch is estimated to be necessary before developed enhanced heat transfer is attained. This is in contrast to an entrance region in an empty tube where heat transfer in the entrance region is higher than in the developed region of the empty tube.

6. The thermal-hydraulic results also point out that the core region of the insert is responsible for the major portion of the heat-transfer enhancement.

7. Empirical correlations are presented for heat transfer and friction factor that are applicable for the design of in-tube enhancement using bent-strip inserts for cooling of hot gases. In the case of heat transfer, different correlations are given for the entrance region and developed region of the flow.
8. For conditions investigated, performance evaluation studies indicate that under a constant pumping power constraint, favorable enhancement is available in the Reynolds number range of 3000 to 30,000. However, under a constant pressure drop constraint, favorable enhancement is available only in the lower Reynolds number range of about 3000 to 5000. Hence, bent-strip inserts are not recommended at higher Reynolds number flows and in cases where allowances cannot be made for high pressure drops.

9. Flow-visualization studies suggest the existence of a swirl flow at the wall and also points out that maximum shear stresses occurs near the contact points of the insert and the tube wall.

10. Flow patterns also confirm the existence of an entry length region corresponding to a distance of about one to two times the insert pitch.

11. The water/dye-injection flow-visualization technique suggests a surface penetration/surface renewal type mechanism at the wall.

12. A semiempirical model using surface renewal/surface penetration theory describes the heat transfer data with reasonable accuracy.
B. Recommendations for Future Work

1. Due to the high pressure-drop penalty, it may be advantageous to use bent-strip inserts in laminar and transition heat-transfer applications. Flow-visualization studies have indicated that even in laminar Reynolds number ranges, the bent-strip insert tends to "turbulize" the flow. Heat-transfer and pressure-drop data in the laminar and transition range would be a valuable extension to this study.

2. Using bent-strip inserts in liquids may have possible applications and should be investigated.

3. In high temperature, real-life applications, an accurate radiant heat-transfer model is required. The measurement of the heat-transfer coefficient from the gas to the insert will be required to obtain more accurate radiation results.

4. By attaining higher gas temperatures, the effect of bulk properties can be evaluated. The effect of variable properties will be useful in fire-tube boiler applications, where the operating temperatures are much higher.

5. Fundamental studies of velocity and temperature profiles in the tube would give more insight into the mechanism of enhancement. Also, the surface renewal frequency could be estimated by using surface film hot-wire anemometry. This would facilitate the development of an independent method to predict the heat transfer by analytical means.
VI. REFERENCES


VII. ACKNOWLEDGEMENTS

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DATA ACQUISITION PROGRAM FOR TWISTED TAPE
INSERTS IN FIRE TUBE BOILERS

JANUARY 1985

DEFINITION OF VARIABLES

Date........date of run
Operator........name of operator
Type........type of twisted tape
Run_no........run number for specific tape
Dataout(1)....air density at orifice, lbm/cu-ft
Dataout(2)....air viscosity at orifice, lbm/hr-ft
Dataout(3)....air mass flow rate, lbm/hr
Dataout(4)....average inlet air temperature, deg.F
Dataout(5)....average outlet air temperature, deg.F
Dataout(6)....air temperature drop, deg.F
Dataout(7)....specific heat of air, Btu/lbm-deg.F
Dataout(8)....heat transfer from air, Btu/hr
Dataout(9)....heat transfer to water thru #4, Btu/hr
Dataout(10)....heat transfer to water thru #3, Btu/hr
Dataout(11)....heat transfer to water thru #2, Btu/hr
Dataout(12)....heat transfer to water thru #1, Btu/hr
Dataout(13)....total heat transfer to water, Btu/hr
Dataout(14)....upstream temp drop across teflen, deg.F
Dataout(15)....downstream temp drop across teflen, deg.F
Dataout(16)....axial conduction from upstream, Btu/hr
Dataout(17)....axial conduction from downstream, Btu/hr
Dataout(18)....energy balance

Datain(1).....test sec. wall temp 1-top, deg.F
Datain(2).....test sec. wall temp 1-nearside, deg.F
Datain(3).....test sec. wall temp 1-farside, deg.F
Datain(4).....test sec. wall temp 2-top, deg.F
Datain(5).....test sec. wall temp 2-nearside, deg.F
Datain(6).....test sec. wall temp 2-farside, deg.F
Datain(7).....test sec. wall temp 3-top, deg.F
Datain(8).....test sec. wall temp 3-nearside, deg.F
Datain(9).....test sec. wall temp 3-farside, deg.F
Datain(10)....test sec. wall temp 4-top, deg.F
Datain(11)....test sec. wall temp 4-nearside, deg.F
Datain(12)....test sec. wall temp 4-farside, deg.F
Datain(13)....test sec. wall temp 5-top, deg.F
Datain(14)....test sec. wall temp 5-nearside, deg.F
Datain(15)......test sec. wall temp 5-farside, deg.P
Datain(16)......water inlet-4, deg.P
Datain(17)......water outlet-4, deg.P
Datain(18)......water inlet-3, deg.P
Datain(19)......water outlet-3, deg.P
Datain(20)......water inlet-2, deg.P
Datain(21)......water outlet-2, deg.P
Datain(22)......water inlet-1, deg.P
Datain(23)......water outlet-1, deg.P
Datain(24)......air inlet-1, deg.P
Datain(25)......air inlet-2, deg.P
Datain(26)......air inlet-3, deg.P
Datain(27)......air inlet-4, deg.P
Datain(28)......air outlet-1, deg.P
Datain(29)......air outlet-2, deg.P
Datain(30)......air outlet-3, deg.P
Datain(31)......air outlet-4, deg.P
Datain(32)......room temperature, deg.P
Datain(33)......temp across teflon-upstream-1, deg.P
Datain(34)......temp across teflon-upstream-2, deg.P
Datain(35)......temp across teflon-downstream-1, deg.P
Datain(36)......temp across teflon-downstream-2, deg.P
Datain(37)......atmospheric reference pressure, in.water
Datain(38)......static pressure-inlet heater box, in.water
Datain(39)......pressure drop-orifice, in.water
Datain(40)......static pressure-outlet orifice, in.water
Datain(41)......pressure drop-test section, in.water
Datain(42)......static pressure-downstream test sec., in.water
Datain(43)......atmospheric pressure, in.Hg
Datain(44)......water flow rate, lbm/hr

SegmentX(1)......outlet air temp in segment X
SegmentX(2)......specific heat of air in segment X
SegmentX(3)......average wall temp-upstream in segment X
SegmentX(4)......average wall temp-downstream in segment X
SegmentX(5)......overall heat transfer coeff. in segment X
SegmentX(6)......air-side heat transfer coeff. in segment X
SegmentX(7)......thermal conductivity of air in segment X
SegmentX(8)......Nusselt number of air in segment X
SegmentX(9)......viscosity of air in segment X
SegmentX(10)......Prandtl number of air in segment X
SegmentX(11)......Reynold number in segment X

ASSIGNMENT OF MULTIPLEXER CHANNELS TO ANALOG INPUTS

Mux Channel
Input
0......pressure transducer output
1......Datain(1)
1219

1229

1239

1249

1259

1269

1279

1289

1299

1399

1499

1599

1699

DIM Rdg(29)
DIM Dms$(3)
Dms$(1)="FLR12FL6M0T2" : dcv, range=10, auto=off, filter=off
Dms$(2)="HSM992L1M9S110STNT3Q" : 10 Readings and get mean
Dms$(3)="HSM992L1M9S120STNT3Q" : 20 Readings and get mean

1319

Control of the Printer

1349

1359

1369 Key=7@1
1379 DISP "SET PAPER AT TOP OF PAGE"
1389 PAUSE
1399 WRITE BIN Key:27,119

1419

Input of Run Type

1449

1469

1479 INPUT "DATE?", Dates
1489 INPUT "OPERATOR?", OperatorS
1499 INPUT "TYPE OF INSERT?", TypeS

1519

1529 Calculation of Radiation Area and Shape Factors

1549

1559

1569 Sigma=1.712E-9
1579 INPUT "Length Of Insert?", Length
1589 INPUT "Pitch of Insert?", Pitch
1599 INPUT "Width of Insert?", Width
1619 IF Pitch<9.5 THEN 1659
1629 Rad_con=Sigma/(1/9.964+1/7.1+1/9-1)*2*28.4375*Width/149.94
1639 Areall=28.4375*Width/2/144
1649 Areall=(Length-53.625)/17.875*28.4375*Width/2/144
1659 IF Pitch<>16 THEN 1699
1669 Rad_con=Sigma/(1/9.994+1/7.1+1/9-1)*2*18.8491*Width/149.94
1679 Areall=18.8491*Width/2/144
1689 Areall=(Length-53.625)/17.875*18.8491*Width/2/144
1699 IF Pitch<6 THEN 1739
1709 Rad_con=Sigma/(1/9.9+1/7.1+1/9-1)*2*25.2791*Width/149.94
1719 Areall=25.2791*Width/2/144
1729 Areall=(Length-53.625)/17.875*25.2791*Width/2/144
1739 IF (Pitch<9.5) AND (Pitch<16) AND (Pitch>6) THEN 1599
1749

1759

Operating conditions monitoring sequence

1779

1789

1799

1809
1810 Re-run: INPUT "ENTER RUN NUMBER?", Run_no
1820 INPUT "FLOW RANGE (1) 6900-29999 or (2) 3000-6990 ? 1 or 2 ",Or_key
1830 INPUT "ATMOSPHERIC PRESSURE, in IN.Hg?",Data_in(43)
1840 OUTPUT 789 USING Mux:32
1850 Mux: IMAGE ";C",22,"E"
1860 Dvm=1 : Flag for averaging 10 to 20 readings
1870 OUTPUT 722; Dms$(2)
1880 OUTPUT 722; Dms$(1)
1890 GOSUB Reading
1900 Rawdata(32)-Rdg
1910 Temp_data=Rawdata(32)*1000
1920 Data_in(32)=1.139*Temp_data+1.139*Temp_data^2-.0015*Temp_data^3
1930 OUTPUT 789 USING Mux:8
1940 OUTPUT 722; Dms$(2)
1950 OUTPUT 722; Dms$(1)
1960 FOR I=1 TO 2
1970 BEEP
1980 DISP "SET SCANNIVALVE AT CHANNEL ",I
1990 PAUSE
2000 GOSUB Reading
2010 Rawdata(I+36)=Rdg
2020 Data_in(I+36)=1.139*Rawdata(I+37)+1.139*Rawdata(I+36)
2030 NEXT I
2040 BEEP
2050 DISP "SET SCANNIVALVE AT CHANNEL 3"
2060 PAUSE
2070 Pflag=1
2080 Repeat_mass:!
2090 ON KEY #8 GOTO Monit
2100 GOSUB Reading
2110 Rawdata(39)=Rdg
2120 Data_in(39)=1.139*Rawdata(37)+1.139*Rawdata(39)
2130 IF Pflag=1 THEN INPUT "MAXIMUM PRESSURE REACHED? Yes=1, No=0",Yes_no
2140 IF (Pflag=1) AND (Yes_no=0) THEN Next
2150 IF (Pflag=1) AND (Yes_no=1) THEN INPUT "PRESSURE IN in.H2O?",Data_in(36+1)
2160 Next: DISP "TO MONITER TEMP: PAUSE;CONT Monit EXECUTE OR HIT k9"
2170 IF Pflag=1 THEN Pflag=0
2180 GOSUB Orifice
2190 Data_out(3)=Mass
2200 OUTPUT 16 USING Scrfmt1;Data_out(3)
2210 Scrfmt1: IMAGE "MASS FLOW RATE OF AIR = ",3D.3D," lbm/hr"
2220 GOTO Repeat_mass
2230 Monit:!
2240 PRINT CHR$(27)&CHR$(69)
2250 DISP "SET AIR TEMPERATURE"
2260 BEEP
2270 Repeat_temp: OUTPUT 722; Dms$(1)
2280 OUTPUT 722; Dms$(2)
2290 ON KEY #1 GOTO Take_data
2300 DISP "TO TAKE DATA PAUSE AND CONT Take_data EXECUTEOR HIT k1"
2310 OUTPUT 789 USING Mux:24
2320 GOSUB Reading
2330 Rawdata(24)=Rdg
2340 Temp_data=Rawdata(24)*1000
2350 Data_in(24)=1.139*Temp_data+1.139*Temp_data^2-.0015*Temp_data^3
2360 OUTPUT 16 USING Scrfmt2;Data_in(24)
2370 Scrfmt2: IMAGE "AVERAGE INLET AIR TEMPERATURE = ",3D.3D," deg.F"
2380 FOR I=1 TO 10
2390 OUTPUT 722; Dms$(1)
2400 OUTPUT 789 USING Mux:23
DATA COLLECTION AND CONVERTING TO TRUE VALUES

Thermocouple readings

Pressure Readings

DATA REDUCTION
Dataout(4) = (Datain(24) + Datain(25) + Datain(26) + Datain(27)) / 4
Dataout(5) = (Datain(28) + Datain(29) + Datain(30) + Datain(31)) / 4
Dataout(6) = Dataout(4) - Dataout(5)

Avg_air_temp = (Dataout(4) + Dataout(5)) / 2 + 459.67
Dataout(7) = 2 Pi * 45E-9 * Avg_air_temp * 2

GOSUB Orifice

Dataout(3) = Mass

Dataout(8) = Dataout(3) * Dataout(6) * Dataout(7)

Avg_air_temp = (Dataout(4) - Dataout(5)) / 2 + 459.67

Dataout(18) = (Datain(13) - Dataout(16) - Dataout(17)) / Dataout(8)

Husselt Number Calculation

 FOR I = 1 TO 4 STEP -1
 IF I = 4 THEN Sub(3) = Dataout(5)
 IF I <> 4 THEN Sub(3) = Segment(I*1, 1)
 IF I = 1 THEN Area = Aceall
 IF I <> 1 THEN Area = Aceal
 Sub(2) = Dataout(3)
 IF I = 1 THEN Sub(5) = Dataout(12) - Dataout(16)
 IF I = 3 THEN Sub(5) = Dataout(10)
 IF I = 4 THEN Sub(5) = Dataout(9) - Dataout(17)
 Sub(6) = Segment(I, 2) = (Datain(3*I-2) + Datain(3*I-1) + Datain(3*I)) / 3
 Sub(7) = Segment(I, 4) = (Datain(3*I+1) + Datain(3*I+2) + Datain(3*I+3)) / 3
 IF I = 1 THEN Sub(8) = 1.039343
 IF I = 3 THEN Sub(8) = 1.087407
 IF I = 4 THEN Sub(8) = 0.92701
 CALL Husselt(Sub(*), Corr_nu(I), Rad_con, Area, Qt(I), T_ins(I))
 Segment(I, 1) = Sub(1)
 Segment(I, 2) = Sub(4)
 Segment(I, 5) = Sub(8)
 Segment(I, 6) = Sub(9)
 Segment(I, 7) = Sub(10)
 Segment(I, 8) = Sub(11)
 Segment(I, 9) = Sub(12)
 Segment(I, 10) = Sub(13)
 Segment(I, 11) = Sub(14)
 NEXT I

 friction factor calculation
Density1 = (1.325 + Datain(43) - 0.97527 * Datain(42)) / (Dataout(4) + 459.67)
Density2 = (1.325 + Datain(43) - 0.97527 * (Datain(42) - Datain(41))) / (Dataout(5) + 459.67)
Delp_mom = (25.71876 * Datain(3)) * 2 * (1/Density1 - 1/Density2) / (369.2 * 32.174)
Delp_tot = Datain(41) * 5.2024 + Delp_mom
Cf = 31.432.11228 * Delp_tot * Datain(43) / (Dataout(3) * 2 * Dataout(4) + Dataout(5) + 92.81)

OUTPUT 16 USING Scrfmt5; Dataout(18)
Scr fmt5: IMAGE "ENERGY BALANCE = ",0D,4D
INPUT "IS ENERGY BALANCE OK? Yes=1, No=0", Yes_no
IF Yes_no=0 THEN Repeat_mass

PRINTING OUT THE RAW AND REDUCED DATA

WRITE BIN Key; 32, 32
WRITE BIN Key; 27, 49
WRITE BIN Key; 32, 32, 32, 32, 32, 32, 32, 32, 32, 32, 32, 32, 32, 32, 32, 32
WRITE BIN Key; 32, 32, 32, 32, 32, 32, 32, 32, 32, 32, 32, 32, 32, 32, 32, 32
WRITE BIN Key; 32, 32, 32
WRITE BIN Key; 32, 32, 32, 32, 32, 32, 32, 32, 32, 32, 32, 32, 32, 32, 32, 32
WRITE BIN Key; 32, 32
WRITE BIN Key; 27, 49
EOL Key; CHR$(9)
WRITE BIN Key; 13
OUTPUT Key USING Pfmt1; Type$, Run no
OUTPUT Key USING Pfmt2; Date$, Operator$
Pline: IMAGE L, "---"
Pline1: IMAGE L, "---"
OUTPUT Key USING Pline1;
WRITE BIN Key; 9
OUTPUT Key; "RAW DATA"
OUTPUT Key USING Pline1;
WRITE BIN Key; 9
OUTPUT Key; "TEMPERATURE DATA"
WRITE BIN Key; 9
OUTPUT Key USING ",K", "Location Scanner Reading Temperature"
OUTPUT Key USING Pline1;
WRITE BIN Key; 9
OUTPUT Key USING ",K", "Channel Volts deg.F"
WRITE BIN Key; 9
OUTPUT Key; "REDUCED RUN DATA"
OUTPUT Key USING Pline1;
WRITE BIN Key; 9
OUTPUT Key USING Pline1;
OUTPUT Key USING Pline1;
Ptemps(1) = "Test.Sec.wall 1-top 1" 1
Ptemps(2) = "Test.Sec.wall 1-ms 2"
Ptemps(3) = "Test.Sec.wall 1-fs 3"
Ptemps(4) = "Test.Sec.wall 2-top 4"
Ptemps(5) = "Test.Sec.wall 2-ms 5"
Ptemps(6) = "Test.Sec.wall 2-fs 6"
Ptemps(7) = "Test.Sec.wall 3-top 7"
4150 PtempS(8)="Test.Sec.wall 3-ns 8 "
4160 PtempS(9)="Test.Sec.wall 3-fs 9 "
4170 PtempS(10)="Test.Sec.wall 4-top 10 "
4180 PtempS(11)="Test.Sec.wall 4-ns 11 "
4190 PtempS(12)="Test.Sec.wall 4-fs 12 "
4200 PtempS(13)="Test.Sec.wall 5-top 13 "
4210 PtempS(14)="Test.Sec.wall 5-ns 14 "
4220 PtempS(15)="Test.Sec.wall 5-fs 15 "
4230 PtempS(16)="Water inlet-4 16 "
4240 PtempS(17)="Water outlet-4 17 "
4250 PtempS(18)="Water inlet-3 18 "
4260 PtempS(19)="Water outlet-3 19 "
4270 PtempS(20)="Water inlet-2 20 "
4280 PtempS(21)="Water outlet-2 21 "
4290 PtempS(22)="Water inlet-1 22 "
4300 PtempS(23)="Water outlet-1 23 "
4310 PtempS(24)="Air inlet-1 24 "
4320 PtempS(25)="Air inlet-2 25 "
4330 PtempS(26)="Air inlet-3 26 "
4340 PtempS(27)="Air inlet-4 27 "
4350 PtempS(28)="Air outlet-1 28 "
4360 PtempS(29)="Air outlet-2 29 "
4370 PtempS(30)="Air outlet-3 30 "
4380 PtempS(31)="Air outlet-4 31 "
4390 PtempS(32)="Room temperature 32 "
4400 PtempS(33)="Teflon up stream-1 33 "
4410 PtempS(34)="Teflon up stream-2 34 "
4420 PtempS(35)="Teflon down stream-1 35 "
4430 PtempS(36)="Teflon down stream-2 36 "
4440 Pfmt3: IMAGE L,28A,5D,6X,5D,6D
4450 PpresS(1)="Atmospheric reference 1 "
4460 PpresS(2)="St. inlet heater box 2 "
4470 PpresS(3)="Pres. drop orifice 3 "
4480 PpresS(4)="St. down str. orifice 4 "
4490 PpresS(5)="Pres. drop test sec. 5 "
4500 PpresS(6)="St. up str. test sec. 6 "
4510 PkeyS(1)="Atmospheric pressure in. Hg "
4520 PkeyS(2)="Water flow rate lmb/hr "
4530 Pfmt4: IMAGE L,48A,5D,5D
4540 Pfmt4L: IMAGE L,48A,5D,5D
4550 PredS(0)="Air density at orifice lmb/cu.ft "
4560 PredS(1)="Air viscosity at orifice lmb/hr-ft "
4570 PredS(2)="Air mass flow rate lmb/hr "
4580 PredS(3)="Average inlet air temperature deg.F "
4590 PredS(4)="Average outlet air temperature deg.F "
4600 PredS(5)="Air temperature drop deg.F "
4610 PredS(6)="Specific heat of air Btu/lbm-deg.F "
4620 PredS(7)="Heat transfer from air Btu/hr "
4630 PredS(8)="Heat transfer to water thru # 4 Btu/hr "
4640 PredS(9)="Heat transfer to water thru # 3 Btu/hr "
4650 PredS(10)="Heat transfer to water thru # 2 Btu/hr "
4660 PredS(11)="Heat transfer to water thru # 1 Btu/hr "
4670 PredS(12)="Total heat transfer to water Btu/hr "
4680 PredS(13)="Upstream temp. drop across teflon deg.F "
4690 PredS(14)="Downstream temp. drop across teflon deg.F "
4700 PredS(15)="Axial conduction from upstr. end Btu/hr "
4710 PredS(16)="Axial conduction from dnstr. end Btu/hr "
4720 PredS(17)="Energy balance "
4730 FOR I=1 TO 18
4740 OUTPUT Key USING Pfmt3;PtempS(I),Rawdata(I),Datain(I)
139

4758 OUTPUT Key USING Pfmt41:PredS(I),Dataout(I)
4765 NEXT I:
4770 PsegS(1,1)="Inlet air temperature at 1" deg.F
4780 PsegS(1,2)="Heat tran. coeff.-1 Btu/sq.ft-hr-deg.F"
4790 PsegS(1,3)="Nusselt number-1"
4800 PsegS(1,4)="Prandtl number-1"
4810 PsegS(1,5)="Reynolds number-1"
4820 PsegS(2,1)="Inlet air temperature at 2" deg.F
4830 PsegS(2,2)="Heat tran. coeff.-2 Btu/sq.ft-hr-deg.F"
4840 PsegS(2,3)="Nusselt number-2"
4850 PsegS(2,4)="Prandtl number-2"
4860 PsegS(2,5)="Reynolds number-2"
4870 PsegS(3,1)="Inlet air temperature at 3" deg.F
4880 PsegS(3,2)="Heat tran. coeff.-3 Btu/sq.ft-hr-deg.F"
4890 PsegS(3,3)="Nusselt number-3"
4900 PsegS(3,4)="Prandtl number-3"
4910 PsegS(3,5)="Reynolds number-3"
4920 PsegS(4,1)="Inlet air temperature at 4" deg.F
4930 PsegS(4,2)="Heat tran. coeff.-4 Btu/sq.ft-hr-deg.F"
4940 PsegS(4,3)="Nusselt number-4"
4950 PsegS(4,4)="Prandtl number-4"
4960 PsegS(4,5)="Reynolds number-4"
4970 Pcorr_nus(1)="Nu*F(L/D)*(Tw/Tb)^-.45; (1)"
4980 Pcorr_nus(2)="Nu*F(L/D)*(Tw/Tb)^-.45; (2)"
4990 Pcorr_nus(3)="Nu*F(L/D)*(Tw/Tb)^-.45; (3)"
5000 Pcorr_nus(4)="Nu*F(L/D)*(Tw/Tb)^-.45; (4)"
5010 Pcorr_nus(5)="Radiation Correction (1) Btu/hr"
5020 Pcorr_nus(6)="Radiation Correction (2) Btu/hr"
5030 Pcorr_nus(7)="Radiation Correction (3) Btu/hr"
5040 Pcorr_nus(8)="Radiation Correction (4) Btu/hr"
5050 Pcorr_nus(9)="Insert Temperature (1) deg.F"
5060 Pcorr_nus(10)="Insert Temperature (2) deg.F"
5070 Pcorr_nus(11)="Insert Temperature (3) deg.F"
5080 Pcorr_nus(12)="Insert Temperature (4) deg.F"
5090 FOR I=1 TO 3
5100 OUTPUT Key USING Pfmt3;Ptemp$(6*(I-1)+19),Rawdata(6*(I-1)+19),Datain(6*(I-1)+19)
5110 OUTPUT Key USING Pfmt41;PsegS(I,1),Segment(I,1)
5120 OUTPUT Key USING Pfmt3;Ptemp$(6*(I-1)+20),Rawdata(6*(I-1)+20),Datain(6*(I-1)+20)
5130 OUTPUT Key USING Pfmt41;PsegS(I,2),Segment(I,6)
5140 OUTPUT Key USING Pfmt3;Ptemp$(6*(I-1)+21),Rawdata(6*(I-1)+21),Datain(6*(I-1)+21)
5150 OUTPUT Key USING Pfmt41;PsegS(I,3),Segment(I,8)
5160 OUTPUT Key USING Pfmt3;Ptemp$(6*(I-1)+22),Rawdata(6*(I-1)+22),Datain(6*(I-1)+22)
5170 OUTPUT Key USING Pfmt41;Pcorr_nus(I),Corr_nu(I)
5180 OUTPUT Key USING Pfmt3;Ptemp$(6*(I-1)+23),Rawdata(6*(I-1)+23),Datain(6*(I-1)+23)
5190 OUTPUT Key USING Pfmt41;PsegS(I,4),Segment(I,10)
5200 OUTPUT Key USING Pfmt3;Ptemp$(6*(I-1)+24),Rawdata(6*(I-1)+24),Datain(6*(I-1)+24)
5210 OUTPUT Key USING Pfmt41;PsegS(I,5),Segment(I,11)
5220 NEXT I
5230 OUTPUT Key USING Pline:
5240 OUTPUT Key USING Pfmt41;PsegS(4,1),Segment(4,1)
5250 WRITE BIN Key;9
5260 OUTPUT Key USING "$,K;" PRESSURE DATA"
WRITE BIN Key:9

For I=1 TO 3

Output Key USING Pfmt4;Pcorc_nu$(I+4).Tins(I)-459.67

NEXT I

RECORDING OF RAW AND REDUCED DATA ON DISK

FILE NAME IS "F8,0"

BEEP

DISP "DATA READY TO BE RECORDED ON DISK. IF OK CONT"

PAUSE

DISP "INSERT DISK INTO DISK DRIVE 0 AND CONT"

PAUSE

Typel$=""

FOR I=1 TO LEN(TypeS)

IF TypeS[I,1]="" THEN Out

Typel$=Typel$&TypeS[I,1]

NEXT I

Out: NEXT I

FileS="FR"&VALS(Run_no)&TypeS

OUTPUT 16;"FILE NAME IS ",FileS

INPUT "IS FILENAME OK? Yes=1, No=0 ",Yes_no

IF Yes_no=0 THEN Name

IF LEN(FileS)>6 THEN Name

GOTO Out1

Out1: LINPUT "ENTER THE CORRECT FILENAME PLEASE",FileS

Out: CREATE FileS,5
ASSIGN 11 TO FILE$  
PRINT $1;DATE$;OPERATOR$;TYPE$;RUN_no;RAWdata(*);DATAin(*);DATAout(*)  
CHECK READ $1  
PROTECT FILE$;"Filetube"  
OUTPUT Key USING Pfat$;FILE$  
PRINT "FILE: IMAGE L, "DATA RECORDED ON DISK: FILE NAME IS ", 6A  
WRITE BIN Key;12  
BEEP  
INPUT "ANOTHER RUN? Yes=1, No=0", Yes_no  
IF Yes_no»1 THEN Rerun  
DISP "BYE FOR NOW"  
STOP  
END  
!
!
! SUBROUTINE TO CALCULATE THE MASS FLOW RATE OF AIR  
!
Orifice:  
Dataout(1)=(1.3259*DATAin(43)+.97527*DATAin(38))/(DATAin(32)+459.67)  
Dataout(2)=2.6866E-3*(DATAin(32)+459.67)-1.5/(DATAin(32)+658.37)  
Orifice data1:  
DATA 12.2525,7,15000,0.2612,25000,0.2601,35000,0.2596,50000,0.2591,75000,  
02558  
DATA 100000,0.2582,15000,0.2579  
Orifice data2:  
DATA 17.2525,7,15000,0.2612,25000,0.2601,35000,0.2596,50000,0.2591,75000,  
02558  
DATA 100000,0.2582,15000,0.2579  
IF Or_key=1 THEN RESTORE Orifice data1  
IF Or_key=2 THEN RESTORE Orifice data2  
READ Re_cons,Or_num  
FOR I=1 TO Or_num  
READ Re_or(I),Cons_or(I)  
NEXT I  
Re=Re_or(1)  
Cons=Cons_or(1)  
Constant: Mass=12932.9505*CONS*SQRT(Dataout(1)*Datain(39))  
Re_check=Re_cons*Mass/Dataout(2)  
IF ABS(Re_check-Re)<1 THEN RETURN  
FOR I=1 TO 6  
IF Re_check<Re_or(I) THEN Extrapolate  
IF Re_check>Re_or(I) THEN Interpolate  
NEXT I  
Interpolate: Cons=Cons_or(I)+(Re_check-Re_or(I))*(Cons_or(I+1)-Cons_or(I))/  
(Re_or(I+1)-Re_or(I))  
Extrapolate: Cons=Cons_or(I)+(Re_check-Re_or(I))*(Cons_or(I+1)-Cons_or(I))/  
(Re_or(I+1)-Re_or(I))  
Re=Re_check  
GOTO Constant  
STOP  
END  
!
SUBROUTINE TO TAKE 10/20 READINGS AND FIND THE MEAN USING THE VOLTOMETER INTERNAL PROGRAM MODE

SET TIMEOUT 7;2000 ! If it hangs up
ON INT #7,1 GOSUB Clear_dvm
OUTPUT 722;"XI"
St_ct=0 ! Timer set to zero
STATUS 722;Status
St_ct=St_ct+1 ! Timer counter
IF St_ct=5 THEN REMOTE 722! If it misses the program execution end notice
IF St_ct>800 THEN 6550
IF Status<>66 THEN 6520
OUTPUT 722;"SO0REM"
ENTER 722;Rdg
RETURN
Clear_dvm: !
DISP "CLEARING VOLTOMETER"
CLEAR 722
REMOTE 722
IF Dvm=1 THEN OUTPUT 722;Dmm$(2)
IF Dvm=2 THEN OUTPUT 722;Dmm$(3)
OUTPUT 722;Dmm$(1)
RETURN

SUBROUTINE TO CALCULATE Nu, Re and Pr FOR EACH SEGMENT

SUB Nusselt(Sub(*),Corr_nu,Rad_con,Area,Qr,T_ins)
Sub(l)-Sub(3) " ~ ~
Iterate: Avg air_temp=(Sub(l)+Sub(3))/2+459.67
Sub(4)=.2231+.72E-5*Avg_air_temp-2.93E-9*Avg air_temp^2
IF ABS(Sub(1)-Sub(3))/Sub(4)/Sub(2) < T1 THEN Done
Sub(1)-Sub(3)+Sub(5)/Sub(4)/Sub(2)
GOTO Iterate
Done: Tb=(Sub(l)+Sub(3))/2+459.67
Len_con=Sub(8)
Sub9)-Sub(8)*Sub(5)*LOG((Sub(3)-Sub(7))/(Sub(1)-Sub(6)))/.990256/(Sub(3)-Sub(7)-(Sub(1)-Sub(6))))
Sub9)=1/(1/Sub(8)-.994783)
T_ins=Tb-Rad_con/Sub(9)*T_ins^4-((Sub(3)+Sub(6))/2+459.62)^4)
IF ABS(T_ins-T_ins) < .1 THEN Done
T_ins=T_ins
GOTO 6990
Done:
Qc=Sub(5)-Qr
Qc=Sub(5)-Qc*LOG((Sub(3)-Sub(7))/(Sub(1)-Sub(6)))/.990256/(Sub(3)-Sub(7)-(Sub(1)-Sub(6))))
Ha=1/(1/Sub(8)-.994783)
IF ABS(Ha-Sub(9)) < .01 THEN Done2
6900 Sub(9)=Ha
6900 GOTO 6880
7000 Done2: Sub(9)=Ha
7010 Sub(10)=1.3466E-2+2.2323E-5*(Sub(1)+Sub(3))/2
7020 Sub(11)=2.67*Sub(9)/12/Sub(10)
7030 Sub(12)=2.6286E-3*Avg_air_temp^1.5/(Avg_air_temp+198.7)
7040 Sub(13)=Sub(12)*Sub(4)/Sub(10)
7050 Sub(14)=5.722425*Sub(2)/Sub(12)
7060 Corr_nu=Sub(11)*((Sub(6)+Sub(7))/2+459.67)/((Sub(1)+Sub(3))/2+459.67)^.4
7070 SUBEND
7080 !
IX. APPENDIX B: MEASUREMENT OF INSERT TEMPERATURE TO VERIFY
ASSUMPTIONS MADE TO ESTIMATE RADIANT HEAT TRANSFER FROM THE INSERT

A. Measurement of Insert Temperature

As mentioned in the data-reduction procedure, the radiation heat transfer from the insert to the tube was estimated by using the following assumptions:

1. The convective heat transfer coefficient at the wall, $h_a$, is equal to the insert convective heat-transfer coefficient, $h_i$.
2. Under steady-state conditions, the convective heat transfer from the hot gas to the insert is equal to the radiation from the insert to the tube wall.
3. Axial heat conduction along the insert length is negligible.
4. Radiation from the hot gas to the wall or to the insert is negligible.

To verify whether these assumptions are valid, eleven 28 gage chromel/alumel thermocouples were attached to insert A6 at various axial locations to measure the insert temperature. Five of these thermocouples were attached at locations corresponding to the ends of the four calorimeter segments of the test section. The other six thermocouples were placed at different locations in the insert length corresponding to segment 3 of the test section. This was thus done to obtain a radial and axial average insert temperature in segment 3.

To attach the thermocouple to the insert at each location, a 0.125 in. diameter hole was drilled partway into the insert. The
thermocouple bead was then placed in the hole and was filled with silver solder. The lead wire was brought along the length of the insert for a few inches and then axially along the side of the insert towards the downstream end of the test section. Once all the thermocouples were attached to the insert, the insert was pushed into the test section. The eleven thermocouple lead wires were brought out from the test section between the teflon and asbestos spacer at the downstream test section end and connected to the thermocouple reference junction.

To permit measurement of the eleven extra thermocouples, a Hewlett Packard Scanner HP 3488A was used in addition to the other instrumentations described earlier. The standard heat-transfer experiments were conducted, and in addition the insert temperatures were also measured for 14 different air-side Reynolds numbers.

B. Results and Verification of Assumptions

Using the simple radiation model given in Equation 2.6, the radiant energy transfer from the insert to the wall in segment 3 was calculated. Figure B.1 compares the measured insert temperature and the estimated insert temperature on the basis of the assumptions given earlier and using Equation 2.8 given below.

\[
q_{r,4} = \frac{\sigma(T_{i,4}^4 - T_{s,4}^4) A_i}{\frac{1}{F_{i-s}} + \left(\frac{1}{\varepsilon_i} - 1\right) + \left(\frac{1}{\varepsilon_s} - 1\right)\left(\frac{A_i}{A_s}\right)}
\]  

(2.8)

The figure indicates that the actual insert temperature is about 22°F
FIGURE B.1. Comparison of estimated and measured insert temperatures
higher than the estimated insert temperature. The apparent
discontinuity in the data occur due to variation in inlet air
temperature. In Figure B.2, the estimated and the measured components
of radiant heat transfer are plotted as functions of Reynolds number.
The figure shows that the estimation procedure based on the assumptions
underpredicts the radiant component of energy transfer by about 30%.
Again, the discontinuity in the data occurs due to the variation in
inlet air temperature. This points out that radiation from the insert
is strongly dependent on the air temperature. To obtain a better
estimate of radiation heat transfer by using the same model, a more
accurate estimate of the insert convective heat-transfer coefficient (and
thereby insert temperature) is necessary. This is also confirmed by
Beckermann (1984), who indicates that the variations in insert
emissivity, \( \varepsilon_i \), and shape factor, \( F_{i-s} \), do not affect the radiant heat
transfer appreciably. However, he points out that the radiant component
of heat transfer is very much dependent on insert temperature.

Although there is a discrepancy of about 30\% in the radiant heat-
transfer component due to the assumptions made, this error affects the
convective heat transfer component in this investigation by only about
1\% to 2\% as seen in Figures B.3. In the uncertainty analysis given in
Appendix C, the uncertainty in the radiant component was assumed to be
50\%. With this pessimistic assumption, the uncertainty in Nusselt
number worked out to be about 6\%. Therefore, in this experimental
investigation, the radiant transfer from insert to the wall is not a
major component of the energy transfer. The convective heat transfer
FIGURE B.2. Comparison of radiant heat transfer evaluated by estimation and measurement of insert temperature.
FIGURE B.3. Comparison of convective heat transfer evaluated by estimation and measurement of insert temperature
coefficients evaluated in study represent only the convective component of heat transfer from the hot gas to tube wall. Nevertheless, in actual fire-tube boiler type applications where the gas temperature could be as high as 1500°F, the radiant transfer from the insert would be considerably higher. This would further enhance the heat transfer from the hot gas to the tube wall in addition to the convective heat transfer enhancement obtained by using bent-strip inserts.
X. APPENDIX C: SAMPLE CALCULATION AND ERROR ANALYSIS

Given below are the equations used for the calculation of experimentally determined quantities necessary for the determination of the Nusselt number in one segment of the calorimeter as well as of the friction factor for the whole test section. Also included is an uncertainty analysis of each calculated quantity obtained by using the method of Kline and McClintock (1953).

The expression for the uncertainty $W$ in a calculated result found from a linear function of variables is

$$W_Z = \left[ \sum_{i=1}^{n} \left( \frac{\partial Z}{\partial X_i} \right)^2 \right]^{\frac{1}{2}} \quad (C.1)$$

where $X_i$ is any of $n$ parameters of which quantity $Z$ is a function.

In the following calculations, the properties of water were calculated by using equations given by Morcos (1974). Morcos had given uncertainties in using his equations, and these values were used in the error analysis. It is also assumed that the uncertainties in the properties of air is ±1%. The measurement uncertainties used in this investigation are based primarily on manufacturers' specifications and in some cases from the calibration data. The following are the uncertainties in measured quantities:

- air temperature = ±3°F
- water temperature = ±0.2°F
- wall temperature = ±0.2°F
water flow rate = \pm 0.75 \text{lbf/hr}

pressure = \pm 0.01 \text{ in. H}_2\text{O}

The air density \( \rho_a \) at the orifice is calculated by using the equation

\[
\rho_a = \frac{P_{\text{atm}} + P_{\text{st}}}{R T_{\text{atm}}} \tag{C.2}
\]

which with conversion factors for compatible units becomes

\[
\rho_a = \frac{70.731 P_{\text{atm}} + 5.2024 P_{\text{st}}}{53.35 T_{\text{atm}}} \text{ lbf/ft}^3 \tag{C.3}
\]

The uncertainty in \( \rho_a \) is given by

\[
W_{\rho} = \frac{1}{53.35 T_{\text{atm}}} \left[ \left( 70.731 W_{P_{\text{atm}}} \right)^2 + \left( 5.2024 W_{P_{\text{st}}} \right)^2 \right]^{\frac{1}{2}} + \left( \frac{70.731 P_{\text{atm}} + 5.2024 P_{\text{st}}}{T_{\text{atm}}} \right)^2 W_{T_{\text{atm}}} \tag{C.4}
\]

\( P_{\text{atm}} = 28.985 \text{ in. Hg} \)

\( P_{\text{st}} = 0.41 \text{ in. H}_2\text{O} \)

\( T_{\text{atm}} = 78.96 ^\circ\text{F} \)

\( \rho_a = 0.07197 \text{ lbf/ft}^3 \)

\( W_{\rho_a} = 0.00012 \text{ lbf/ft}^3 \)
The mass-flow rate of air is calculated by using the following equation:

\[ m_a = C_D \rho_a \left( \frac{\Delta p_{or}}{\rho_a} \right) \]  \hspace{1cm} (C.5)

which with conversion factors for compatible units become

\[ m_a = 12932.95 \ C_D \sqrt{(\rho_a \Delta p_{or})} \] \hspace{1cm} (C.6)

Then, the uncertainty in \( m_a \) is given by

\[
\begin{align*}
W_{m_a} &= 12932.95 \left[ \left( \sqrt{(\Delta p_{or})} W_{C_D} \right)^2 + \left( \frac{C_D}{2} \sqrt{\frac{\Delta p_{or}}{\rho_a}} W_{\rho_a} \right)^2 \\
&+ \left( \frac{C_D}{2} \frac{\rho_a}{\Delta p_{or}} W_{\Delta p_{or}} \right)^2 \right]^{1/2} 
\end{align*}
\] \hspace{1cm} (C.7)

\( C_D = 0.02599 \)
\( \Delta p_{or} = 1.249 \text{ in. } H_2O \)
\( m_a = 100.75 \text{ lbm/hr} \)
\( W_{m_a} = 2.06 \text{ lbm/hr} \)

The heat gained by water in a particular segment (in this case the fourth segment) of the calorimeter is

\[ q_{w,4} = m_w c_{p,w,4} (T_{w2,4} - T_{w1,4}) \] \hspace{1cm} (C.8)

and the uncertainty in heat gained by water is given by
$$W_{q_{w,4}} = \left[ (m_w c_{pw,4} \Delta T_{w,4})^2 + (c_{pw,4} \Delta T_{w,4} W_{m_w})^2 \right] \frac{1}{2}$$

where $\Delta T_{w,4}$ is the temperature drop ($T_{w2,4} - T_{w1,4}$).

$m_w = 80.34 \text{ lbm/hr}$

$c_{pw,4} = 1.0005 \text{ BTU/lbm-°F}$

$\Delta T_{w,4} = 10.27°F$

$q_{w,4} = 825.44 \text{ BTU/hr}$

$W_{q_{w,4}} = 33.06 \text{ BTU/hr}$

The relation used to calculate heat gained by conduction by the first segment in the calorimeter is

$$q_{cd} = k_t A_t \left( \frac{\Delta T_t}{\Delta X} \right)$$

and the uncertainty in $q_{cd}$ is given by

$$W_{q_{cd}} = \left[ \left( \frac{k_t \Delta T_t}{\Delta X} \right)^2 + \left( \frac{A_t \Delta T_t}{\Delta X} W_{k_t} \right)^2 + \left( \frac{k_t \Delta T_t}{\Delta X} W_{A_t} \right)^2 \right] \frac{1}{2}$$

$k_t = 1.77 \text{ BTU/hr-ft-°F}$

$A_t = 0.05 \text{ ft}^2$

$\Delta T_t = 54.27°F$
\[ \Delta X = 1.0 \text{ in.} \]
\[ q_{cd} = 58.17 \text{ BTU/hr} \]

Assuming \( W_k, W_A, W_{\Delta X} \ll W_{\Delta T} \)

\[ W_{cd} = 0.62 \text{ BTU/hr} \]

Therefore, the actual heat gained by water from hot air in segment 4 is

\[ q_{wc,4} = q_{w,4} - q_{cd} \quad (C.12) \]

and the uncertainty in \( q_{wc,4} \) is given by

\[ W_{q_{wc,4}} = W_{q_{w,4}} + W_{q_{cd}} \quad (C.13) \]

\[ q_{wc,4} = 767.28 \text{ BTU/hr} \]
\[ W_{q_{c,4}} = 33.68 \text{ BTU/hr} \]

The outlet air temperature for the first segment in the calorimeter was calculated by using the following equation:

\[ T_{a1,4} = T_{a2,4} - \frac{q_{wc,4}}{c_{pa,4} \cdot m_a} \quad (C.14) \]

The uncertainty in \( T_{a2,4} \) is given by
The log-mean temperature difference is obtained by using the standard equation given below.

\[
\Delta T_{\text{LMTD}} = \frac{\Delta T_1 - \Delta T_2}{\ln \left( \frac{\Delta T_1}{\Delta T_2} \right)} = \frac{(T_{a1,4} - T_{t,4}) - (T_{a2,4} - T_{t,4})}{\ln \left( \frac{T_{a1,4} - T_{t,4}}{T_{a2,4} - T_{t,4}} \right)}
\]  

(C.16)

The uncertainty in \( \Delta T_{\text{LMTD}} \) is given by the expression

\[
W_{\Delta T_{\text{LMTD}}} = \left[ \left( \frac{\ln \left( \frac{\Delta T_1}{\Delta T_2} \right) - \frac{\Delta T_1 - \Delta T_2}{\Delta T_1}}{\ln \left( \frac{\Delta T_1}{\Delta T_2} \right)} \right)^2 \right]^{\frac{1}{2}} W_{\Delta T_1}
\]
The overall heat-transfer coefficient is calculated by using the expression

\[ U = \frac{q_{wc,4} - q_{r,4}}{A_s \Delta T_{LMTD}} \]  

(C.18)

and the uncertainty in \( U \) is given by expression

\[ W_U = \left[ \left( \frac{1}{A_s \Delta T_{LMTD}} - W_{(q_{wc,4} - q_{r,4})} \right)^2 + \left( \frac{-(q_{wc,4} - q_{r,4})}{A_s \Delta T_{LMTD}} W_{A_s} \right)^2 \right]^{\frac{1}{2}} \]  

(C.19)

\[ A_s = 0.98 \text{ ft}^2 \]
\[ q_{r,4} = 22.31 \text{ BTU/hr} \]
\[ U = 6.95 \text{ BTU/hr-ft}^2-\text{°F} \]
Assuming a 50% uncertainty in $q_{r,4}$,

$W_U = 0.43 \text{ BTU/hr-ft}^2\cdot\text{°F}$

The average heat-transfer coefficient in the first segment is calculated by

$$h_{a,4} = \left[\frac{1}{U} - \frac{D_i \ln\left(\frac{D_t}{D_i}\right)}{2k_s}\right]^{-1} \tag{C.20}$$

and the uncertainty in $h_a$ is given by

$$W_{h_{a,4}} = \left[\left(\frac{1}{U} - \frac{D_i \ln\left(\frac{D_t}{D_i}\right)}{2k_s}\right)^{-2} \left(-\frac{1}{U^2}\right) W_U \right]^2$$

$$+ \left\{\left(\frac{1}{U} - \frac{D_i \ln\left(\frac{D_t}{D_i}\right)}{2k_s}\right)^{-2} \frac{\ln\left(\frac{D_t}{D_i}\right)}{2k_s} W_{D_t} \right\}^2$$

$$+ \left\{\left(\frac{1}{U} - \frac{D_i \ln\left(\frac{D_t}{D_i}\right)}{2k_s}\right)^{-2} \frac{D_i}{2k_s D_t} W_D \right\}^2$$

$$+ \left\{\left(\frac{1}{U} - \frac{D_i \ln\left(\frac{D_t}{D_i}\right)}{2k_s}\right)^{-2} \frac{D_i \ln\left(\frac{D_t}{D_i}\right)}{2k_s} W_{D_i} \right\}^2$$

$$+ \left\{\left(\frac{1}{U} - \frac{D_i \ln\left(\frac{D_t}{D_i}\right)}{2k_s}\right)^{-2} \left(-\frac{1}{U^2}\right) W_{U} \right\}^2 \tag{C.21}$$
D_i = 2.67 in.
D_t = 2.80 in.
k_s = 25 BTU/hr-ft-°F
h_a,4 = 6.98 BTU/hr-ft²-°F
W_h_a,4 = 0.43 BTU/hr-ft²-°F

As explained in the data reduction procedure,

\[ q_{r,4} = \frac{\sigma A_i (T_{i,4} - T_{s,4})}{1/F_{i-s} + (1/\varepsilon_i - 1) + (1/\varepsilon_s - 1)(A_i/A_s)} \quad (C.22) \]

\[ q_{ci,4} = h_{i,4} A_i (T_{b,4} - T_{i,4}) = q_{r,4} \quad (C.23) \]

\[ \sigma = 0.1712 \times 10^{-8} \text{ BTU/hr-ft}^2\text{-°R} \]
F_{i-s} = 0.964
\[ \varepsilon_i = 0.7 \]
\[ \varepsilon_s = 0.9 \]
A_i = 0.2484 ft²
A_s = 1.0413 ft²
T_{b,4} = 185.24°F = 644.91°R
T_{s,4} = 74.18°F = 553.85°R
T_{i,4} = 632.03°R = 172.36°F
q_{r,4} = 22.31 BTU/hr

The Nusselt number is obtained by using the standard definition and is given by
The uncertainty in $\text{Nu}_{D,4}$ is obtained by using the following expression:

$$
W_{\text{Nu}_{D,4}} = \left[ \left( \frac{D_i}{k_{a,4}} \right)^2 + \left( \frac{h_{a,4}}{k_{a,4}} \cdot W_{h_{a,4}} \right)^2 + \left( \frac{h_{a,4} D_i}{k_{a,4}^2} \cdot W_{k_{a,4}} \right)^2 \right]^{\frac{1}{2}}
$$

(C.25)

$k_{a,4} = 0.0176$ BTU/hr-ft-°F

$\text{Nu}_{D,4} = 88.18$

$W_{\text{Nu}_{D,4}} = 5.58$

The airside Reynolds number in the fourth segment is calculated by using

$$
\text{Re}_{D,4} = \frac{4m_a}{\pi u_{a,4} D_i}
$$

(C.26)

and the uncertainty in $\text{Re}_{D,4}$ is given by

$$
W_{\text{Re}_{D,4}} = \frac{4}{\pi} \left[ \left( \frac{1}{u_{a,4} D_i} \cdot W_{m_a} \right)^2 + \left( \frac{-m_a}{u_{a,4} D_i} \cdot W_{D_i} \right)^2 + \left( \frac{-m_a}{u_{a,4}^2 D_i} \cdot W_{u_{a,4}} \right)^2 \right]^{\frac{1}{2}}
$$

(C.27)

$u_{a,4} = 0.05103$ lbm/ft-hr

$\text{Re}_{D,4} = 11298$

$W_{\text{Re}_{D,4}} = 290$
This process of calculation is repeated by using the inlet air temperature at one segment as the outlet air temperature for the next segment (i.e., \( T_{a2,3} = T_{a1,4} \)). The second and third segments have no conduction correction; however, heat gain by the coolant in the first segment requires a conduction correction similar to that used for the fourth segment. By using a similar calculating procedure as given in Equations C.8 to C.27, the following values of Nusselt numbers, Reynolds numbers, and the respective uncertainties were calculated:

\[
\begin{align*}
\text{Nu}_{D,3} &= 82.02, \quad \text{Nu}_{D,3}^* = 4.68 \\
\text{Re}_{D,3} &= 10882, \quad \text{Re}_{D,3}^* = 280 \\
\text{Nu}_{D,2} &= 76.62, \quad \text{Nu}_{D,2}^* = 4.56 \\
\text{Re}_{D,2} &= 10447, \quad \text{Re}_{D,2}^* = 269 \\
\text{Nu}_{D,1} &= 70.80, \quad \text{Nu}_{D,1}^* = 4.50 \\
\text{Re}_{D,1} &= 9989, \quad \text{Re}_{D,1}^* = 258
\end{align*}
\]

The pressure drop across the calorimeter is calculated in two steps. First, by knowing the inlet and outlet air temperature, the pressure rise due to temperature drop is calculated by

\[
Ap_c = \frac{1}{g_c} \left( \frac{m_a}{A} \right)^2 \left( \frac{1}{\rho_{a1}} - \frac{1}{\rho_{a2}} \right) \tag{C.28}
\]
and the uncertainty in $\Delta p_c$ is given by the following expression:

$$W_{\Delta p_c} = \frac{1}{A_x g_c} \left\{ \left[ 2 m_a W_m \left( \frac{1}{\rho_{a1}} - \frac{1}{\rho_{a2}} \right) \right]^2 + \left[ \frac{-m_a^2}{\rho_{a1}^2} W \rho_{a2} \right]^2 \right\}^{\frac{1}{2}}$$

(C.29)

$$\rho_{a1} = 0.0487 \text{ lbm/ft}^3$$
$$\rho_{a2} = 0.0611 \text{ lbm/ft}^3$$
$$g_c = 32.174 \times 3600^2 \text{ ft/hr}^2$$
$$\Delta p_c = 0.013 \text{ in. H}_2\text{O}$$
$$W_{\Delta p_c} = 0.003 \text{ in. H}_2\text{O}$$

The actual friction pressure drop is calculated by adding the measured pressure drop to the pressure rise due to momentum:

$$\Delta p_f = \Delta p + \Delta p_c$$

(C.30)

The uncertainty in $\Delta p_f$ is given by

$$W_{\Delta p_f} = W_{\Delta p} + W_{\Delta p_c}$$

(C.31)

$$\Delta p = 0.203 \text{ in. H}_2\text{O}$$
$$\Delta p_f = 0.216 \text{ in. H}_2\text{O}$$
$$W_{\Delta p_f} = 0.013 \text{ in. H}_2\text{O}$$
The Fanning friction factor is calculated by using the expression

\[ f = \frac{8 \Delta p_f}{g_c \left( \frac{\rho_{am}}{2} \right)^2 u_m^2} \frac{D_i}{L_t} = \frac{\pi^2 \Delta p_f \rho_{am} D_i^5}{16 \frac{m_a^2}{L_t}} \]  
(C.32)

The uncertainty in \( f \) is given by

\[
W_f = \frac{\pi^2}{16} \left[ \left( \frac{\rho_{am} D_i^5}{m_a^2 L_t} W_{\Delta p_f} \right)^2 + \left( \frac{\Delta p_f D_i^5}{m_a^2 L_t} W_{\rho_{am}} \right)^2 + \left( \frac{\Delta p_f \rho_{am} S D_i^4}{m_a^2 L_t} W_{D_i} \right)^2 \right. \\
+ \left. \left( \frac{-\Delta p_f \rho_{am} D_i^5}{m_a^2 L_t^2} W_{L_t} \right)^2 + \left( \frac{-2 \Delta p_f \rho_{am} D_i^5}{m_a^3 L_t} W_{\rho_{am}} \right)^2 \right]^{\frac{1}{2}} \]  
(C.33)

\[ \rho_{am} = 0.0548 \text{ lbm/ft}^3 \]
\[ L_t = 71 \text{ in.} \]
\[ f = 0.071 \]
\[ W_f = 0.004 \]
XI. APPENDIX D. CALCULATIONS OF ANGLES REQUIRED FOR FABRICATING THE INSERTS

Inserts A1, AN, and PA were made by a commercial manufacturer. To fabricate similar inserts with different pitches and widths, the relationship between pitch, \( L \), twist angle, \( \gamma \), longitudinal angle, \( \phi \), bending angle, \( \alpha \), and bending line angle, \( \theta \), were required.

Figure D.1 shows the three-dimensional view of the insert at the bend. Figure D.2 shows the plan view and Figure D.3 is view of the section ZZ marked in Figure D.2. The vertical plane passing through BFC in Figure D.1 is shown in Figure D.4.

Considering Figure D.2,

\[
BA = 2BE = \frac{W_A}{\sin(\theta)} \tag{D.1}
\]

where \( W_A \) is the apparent width, which is equal to \( W \sin(\gamma) \).

Again, in Figure D.2,

\[
EG = BE \tan(\theta)
\]

Substituting for \( BE \) from Equation D.1,

\[
EG = \left(\frac{W_A}{2\sin(\gamma)}\right) \tan(\theta) = \frac{W_A}{2\cos(\theta)} \tag{D.2}
\]

In Figure D.2,

\[
BG = BE / \cos(\theta)
\]

substituting for \( BE \) from Equation D.1,
FIGURE D.1. Isometric view of the insert at a bend

FIGURE D.2. Plan view of insert at a bend

\[ BG = \frac{W_A}{(2\cos\theta \sin\theta)} \]  \hspace{1cm} (D.3)

From Figure D.4,

\[ FG = BG \tan\beta = \frac{W_A \tan\beta}{(2\cos\theta \sin\theta)} \]  \hspace{1cm} (D.4)
FIGURE D.3. View of section ZZ

FIGURE D.4. View of vertical plane passing through BFC

Considering Figure D.3,

\[ \tan(\mu) = \frac{FG}{EG} \]

Substituting for FG and EG from Equations D.4 and D.2, respectively,

\[ \tan(\mu) = \frac{W_A \tan\beta/(2\cos\theta \sin\theta)}{W_A/(2\cos\theta)} = \frac{\tan\beta}{\sin\theta} \]

(D.5)
But $\mu = 90-\alpha/2$ and $\beta = 90-\beta/2$, which gives the relationship between $\alpha$, $\phi$, and $\theta$; thus

$$\cot(\alpha/2) = \frac{\cot(\phi/2)}{\sin(\theta)}$$

or

$$\tan(\alpha/2) = \sin(\theta) \tan(\phi/2) \quad (D.6)$$

Figure D.5 shows the vertical plane passing through AD in Figure D.1. Point D in Figure D.5 is the same point D in Figure D.4.

![Figure D.5. View of vertical plane passing through AD](image)

FIGURE D.5. View of vertical plane passing through AD

Considering Figure D.5,

$$\tan(\gamma) = \frac{HD}{AD} \quad (D.7)$$

From Figure D.4,

$$HD = BD \tan\beta$$
and from Figure D.2,

$$AD = BD \tan \theta$$

Substituting for $HD$ and $AD$ in Equation D.7,

$$\frac{\tan(\gamma)}{\tan(\theta)} = \frac{\tan(\beta)}{\tan(\alpha/2)}$$

But since $\beta = 90 - \phi/2$, the relationship between $\gamma$, $\theta$, and $\phi$ is given as

$$\tan(\theta) = \frac{1}{\tan(\gamma) \tan(\phi/2)} \quad (D.8)$$

The design procedure is as follows. For a particular pitch, $L$, the longitudinal angle $\phi$ can be calculated from

$$\frac{\phi}{2} = \tan^{-1}(L/2D)$$

For the value of $\phi$ and a given twist angle, $\gamma$, the bending line angle, $\theta$, is calculated using Equation D.6. Using this value of $\theta$ and the value of $\phi$, the bending angle, $\alpha$, is calculated by using Equation D.8.

To fabricate the inserts, markings at distances of $(L/2)/\sin(\phi/2)$ on a metal strip are made at angle of $\theta$ with the strip edge (as shown by line BA in Figure D.2). Then, the metal strip is bent at that line through an angle of $180 - \alpha$. It should be noted that at every other marking, the strip is bent in the opposite direction.
XII. APPENDIX E: HEAT-TRANSFER AND PRESSURE-DROP DATA

Original data is available from the Heat Transfer Laboratory, Department of Mechanical Engineering, Iowa State University, Ames.

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XIII. APPENDIX F: MULTIPLE LINEAR REGRESSION USING MATRICES

To obtain the values of the parameters in the correlation equations, a matrix approach to multiple linear regression was used. Taking the logarithm of the correlation equations converts them to linear equations with multiple independent variables of the following form:

\[ Y_i = \beta_0 + \beta_1 X_{1i} + \beta_2 X_{2i} + \ldots + \beta_p X_{pi} + \varepsilon_i, \quad i = 1, \ldots, n. \]  

(F.1)

The same model for all the \( n \) observations can be written in matrix form by defining the following matrices:

\[
[Y] = \begin{bmatrix}
Y_1 \\
Y_2 \\
\vdots \\
Y_n
\end{bmatrix}, \quad [X] = \begin{bmatrix}
1 & X_{11} & X_{21} & \ldots & X_{p1} \\
1 & X_{12} & X_{22} & \ldots & X_{p2} \\
\vdots & \vdots & \vdots & \ddots & \vdots \\
1 & X_{1n} & X_{2n} & \ldots & X_{pn}
\end{bmatrix}, \quad [\beta] = \begin{bmatrix}
\beta_1 \\
\beta_2 \\
\vdots \\
\beta_n
\end{bmatrix}, \quad [\varepsilon] = \begin{bmatrix}
\varepsilon_1 \\
\varepsilon_2 \\
\vdots \\
\varepsilon_n
\end{bmatrix}
\]

Then, the model in matrix notation is

\[
[Y] = [X][\beta] + [\varepsilon] \quad \text{(F.2)}
\]

To estimate the parameters, the sum of squares is to be minimized

\[
SS = ([Y] - [X][\beta])^T ([Y] - [X][\beta]) \quad \text{(F.3)}
\]

where \([B]\) is estimate of \([\beta]\). The normal equation can be written as
The least square estimates of \( [B] \) can be obtained by inverting the matrix \( [X]^T [X] \) and is

\[
[B] = ([X]^T [X])^{-1} ([X]^T [Y])
\]  

(F.5)

The estimate of the variance is given by

\[
\]  

(F.6)

Also, the variances and covariances of the b's are easily expressed in matrix notation by

\[
[V](B) = ([X]^T [X]) \, \sigma_E^2
\]  

(F.7)

If \( [C] \) is defined as \( ([X]^T [X])^{-1} \) and the estimated variance \( s_E^2 \) is used as the estimate of \( \sigma_E^2 \), the variance of \( b_h \) is

\[
s_{b_h}^2 = c_{hh} \, s_E^2, \quad h = 1, 2, \ldots, p
\]  

(F.8)

where \( c_{hh} \) are the elements in matrix \( [C] \).

Based on the assumption of normal distribution, 100\% percent confidence interval for \( b_h \) is given by

\[
U = b_h \pm t \left( \frac{(1+\gamma)/2}{n-p-1} \right) s_{b_h}, \quad h = 1, 2, \ldots, p
\]  

(F.9)
The correlation coefficient, \( R^2 \), is given by the following equation

\[
R^2 = \frac{[B]^T[X][Y] - n(ZY_i)^2}{[Y]^T[Y] - n(ZY_i)^2}
\]  

where \( ZY_i \) is given by the (1,1) element in the matrix \([X]^T[X]\).

The equations in this appendix were by solved using the Heat Transfer Laboratory HP 9845B Desktop Computer which has the capability to handle matrix arithmetic without any involved programming.
XIV. APPENDIX G. MATHEMATICAL SOLUTIONS USED IN THE SURFACE RENEWAL/PENETRATION MODEL

The solution of

\[ \frac{\partial T}{\partial t} - \alpha \frac{\partial^2 T}{\partial y^2} = 0 \]  

with the following initial and boundary conditions:

- \( T = T_i \) at \( t = 0 \) and \( y > 0 \)
- \( T = T_w \) at \( t > 0 \) and \( y = 0 \)
- \( T = T_i \) at \( t > 0 \) and \( y = \delta \)

is obtained by using the method of separation of variables.

First, the variables are transformed into nondimensional variables by defining

\[ T = \frac{T - T_w}{T_i - T_w} \quad \text{and} \quad Y = \frac{y}{\delta} . \]

The differential equation can then be written as

\[ \frac{\partial T}{\partial t} - \alpha \frac{\partial^2 T}{\partial Y^2} = 0 \]  

and the initial and boundary conditions become

\[ T[Y,0] = 1 \]
\[ T[0,t] = 0 \]
To obtain a homogeneous boundary condition, a new variable \( W \), is defined as \( W = T - Y \). The differential equation becomes

\[
\frac{\partial W}{\partial t} = \frac{\alpha}{\delta^2} \frac{\partial^2 W}{\partial Y^2}
\]

with initial and boundary conditions

\[
W(t = 0, Y > 0) = 1 - Y
\]
\[
W(t > 0, Y = 0) = 0
\]
\[
W(t > 0, Y = 1) = 0
\]

To solve the above equation, the method of separation of variables is used. Defining

\[
W(t,Y) = U(t) V(Y)
\]

the differential equation becomes

\[
V(Y) \frac{\partial U(t)}{\partial t} = \frac{\alpha}{\delta^2} U(t) \frac{\partial^2 V(Y)}{\partial Y^2}
\]

\[
\frac{U'}{U} = \frac{V''}{V}
\]

In the above equation, the left member is clearly independent of \( Y \), and therefore, the right member also is independent of \( Y \). Similarly the left and right members must be independent of \( t \).
Therefore,

\[
\frac{U'}{(\alpha/\delta^2)U} = \frac{V''}{V} = \xi = \text{a constant}
\]

For \( \xi > 0 \) the solutions turn out to be trivial.

Therefore,

for \( \xi < 0 \), i.e. \( \xi = -\beta^2 \)

\[V'' + \beta^2 V = 0, \ V(0) = V(1) = 0\]

\[V = A \cos(\beta Y) + B \sin(\beta Y)\]

\[V(0) = 0 \rightarrow A = 0\]

\[V = B \sin(\beta Y)\]

\[V(1) = B \sin(\beta) = 0\]

Here, \( B \neq 0 \), as it would imply a trivial solution.

Therefore,

\[\sin(\beta) = 0, \ \text{or} \ \beta = n\pi, \ \text{where} \ n = 0,1,2 \ldots\]

\[V = B \sin(n\pi Y)\]

\[= B_n \sin(n\pi Y), \ n = 0,1,2 \ldots\]

The solution of

\[
\frac{U'}{(\alpha/\delta^2)U} = -\beta^2
\]
\[
\begin{align*}
\frac{U'}{U} &= -\frac{\alpha \beta^2}{\delta^2} \\
\ln U &= -\frac{\alpha \beta^2 t}{\delta^2} + C_1 \\
U &= C \exp\left(-\frac{\alpha \beta^2 t}{\delta^2}\right)
\end{align*}
\]

Combining \(U\) and \(V\)

\[
W = UV = D_n \sin(n \pi Y) \exp\left(-\frac{\alpha \beta^2 t}{\delta^2}\right), \quad n = 0, 1, 2, \ldots
\]

where \(D_n = C B_n\),
or

\[
W = \sum_{n=1}^{\infty} D_n \sin(n \pi Y) \exp\left(-\frac{\alpha \beta^2 t}{\delta^2}\right)
\]

since each value of \(n\) is a solution.

For the initial condition, \(t = 0\)

\[
W = 1 - Y = \sum_{n=1}^{\infty} D_n \sin(n \pi Y) \exp\left(-\frac{\alpha \beta^2 t}{\delta^2}\right)
\]

Using Fourier series
\[ D_n = \frac{2}{1} \int_0^1 (1 - Y) \sin(n\pi Y) dY \]

\[ = 2 \int_0^1 \sin(n\pi Y) dY - 2 \int_0^1 Y \sin(n\pi Y) dY \]

\[ = 2 \left( -\frac{\cos(n\pi Y)}{n \pi} \right) \bigg|_0^1 - 2 \left( -\frac{\cos(n\pi Y)}{n \pi} \right) \bigg|_0^1 + 2 \left( \frac{\cos(n\pi Y)}{n^2 \pi^2} \right) \bigg|_0^1 \]

\[ = 2 \left[ \frac{-\cos(n\pi) + 1}{n \pi} \right] - 2 \left[ \frac{-\cos(n\pi) + 1}{n \pi} \right] + 2 \left[ \frac{\sin(n\pi)}{n^2 \pi^2} \right] \bigg|_0^1 \]

\[ = 2/n \pi \]

\[ W = \sum_{n=1}^{\infty} \frac{2}{n \pi} \sin(n\pi Y) \exp \left( -\frac{\alpha \beta^2 t}{\delta^2} \right) \quad (G.7) \]

Substituting back for \( W \)

\[ T - Y = \sum_{n=1}^{\infty} \frac{2}{n \pi} \sin(n\pi Y) \exp \left( -\frac{\alpha \beta^2 t}{\delta^2} \right) \]

\[ \frac{T - T_w}{T_i - T_w} = \sum_{n=1}^{\infty} \frac{2}{n \pi} \sin(n\pi Y) \exp \left( -\frac{\alpha \beta^2 t}{\delta^2} \right) \]
\[
\frac{T - T_w}{T_i - T_w} = \frac{y}{\delta} + \sum_{n=1}^{\infty} \frac{2}{n\pi} \sin(n\pi y) \exp \left( -\frac{\alpha n^2 \pi^2 t}{\delta^2} \right) \tag{G.8}
\]

The instantaneous heat-transfer rate is given by

\[
q_i(t) = -k \frac{dT}{dy} \bigg|_{y=0} \tag{G.9}
\]

The mean heat transfer rate is

\[
q = \int_0^\infty q_i(t) P(t) \, dt \tag{G.10}
\]

where

\[
P(t) = s \exp(-st) \tag{G.11}
\]

Substituting G.8 into G.9

\[
q_i(t) = -k(T_i - T_w) \left[ \frac{1}{\delta} + \frac{2}{\delta} \sum_{n=1}^{\infty} \frac{1}{n} \cos\left( \frac{n\pi y}{\delta} \right) \frac{n\pi}{\delta} \exp\left( -\frac{\alpha n^2 \pi^2 t}{\delta^2} \right) \right]_{y=0}
\]

\[
= -k(T_i - T_w) \left[ \frac{1}{\delta} + \frac{2}{\delta} \sum_{n=1}^{\infty} \exp\left( -\frac{\alpha n^2 \pi^2 t}{\delta^2} \right) \right]
\]

\[
q = \int_0^\infty \frac{k}{\delta} (T_i - T_w) \left[ \frac{1}{\delta} + \sum_{n=1}^{\infty} \exp\left( -\frac{\alpha n^2 \pi^2 t}{\delta^2} \right) se^{-st} \right] \, dt
\]
\[ q = - \frac{k}{s\delta} (T_i - T_w) \int_0^\infty \left[ e^{-st} + 2 \sum_{n=1}^\infty \exp\left(-s + \frac{an^2\pi^2}{\delta^2}\right) \right] dt \]

\[ q = - \frac{k}{s\delta} (T_i - T_w) \left[ \frac{e^{-st}}{-s} \right]_0^\infty + 2 \sum_{n=1}^\infty \frac{\exp\left(-s + \frac{an^2\pi^2}{\delta^2}\right)}{-s + \frac{an^2\pi^2}{\delta^2}} \]

\[ q = - \frac{ks}{\delta} (T_i - T_w) \left[ \frac{1}{s} + 2 \sum_{n=1}^\infty \frac{\delta^2}{s\delta^2 + an^2\pi^2} \right] \]

\[ q = - \frac{ks}{\delta} (T_i - T_w) \left[ \frac{1}{s} + 2 \sum_{n=1}^\infty \frac{\delta^2/\alpha}{s\delta^2/\alpha + n^2\pi^2} \right] \]

\[ q = - \frac{ks}{\delta} (T_i - T_w) \left[ \frac{1}{s} + \frac{2s\delta^2}{s\alpha} \sum_{n=1}^\infty \frac{1}{s\delta^2/\alpha + n^2\pi^2} \right] \]

Letting \( Z^2 = s\delta^2/\alpha \) and substituting \( Z \) into the above equation,

\[ q = - \frac{k}{\delta} (T_i - T_w) \left[ 1 + 2Z^2 \sum_{n=1}^\infty \frac{1}{Z^2 + n^2\pi^2} \right] \]

\[ q = - \frac{k}{\delta} (T_i - T_w) Z \left[ \frac{1}{Z} + 2Z \sum_{n=1}^\infty \frac{1}{Z^2 + n^2\pi^2} \right] \]

But,
\[
\coth(Z) = \frac{1}{Z} + 2Z \sum_{n=1}^{\infty} \frac{1}{Z^2 + n^2\pi^2}
\]

Therefore,

\[
q = -\frac{k}{\delta} (T_i - T_w) Z \coth(Z)
\]

\[
q = -\frac{k}{\delta} (T_i - T_w) \left(\frac{s\delta^2}{\alpha}\right)^\frac{1}{4} \coth\left(\frac{s\delta^2}{\alpha}\right)
\]

\[
q = -\frac{k}{\delta} (T_i - T_w) \left(\frac{s\delta^2}{k}\right)^\frac{1}{4} \coth\left(\frac{s\delta^2}{\alpha}\right)
\]

\[
q = -\left(T_i - T_w\right) \left(\frac{sk\mathcal{P}}{p}\right)^\frac{1}{4} \coth\left(\frac{s\delta^2}{\alpha}\right) \quad \text{(G.12)}
\]