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Abstract

Inelastic neutron scattering measurements were carried out on the Ti50Pd50-xCrx alloy, which has the potential for being a high temperature shape memory material. For x=0, the transformation temperature is ~800K and for the composition studied (x=8 at.%) TM ~ 400K. The majority of the measurements were performed in the parent, β-phase, up to 873K. Most of the phonons propagating along the three symmetry directions $[\zeta00]$, $[\zeta\zeta\zeta]$, and $[\zeta\zeta0]$ were well defined with the exception of the $[\zeta\zeta0]$ transverse acoustic mode with displacements along the $[-\zeta\zeta0]$ corresponding to the C'=1/2(C11-C12) elastic constant. These phonons are well defined for small ζ, but for $\zeta > 0.15$ they are strongly overdamped near the transition temperature, but become better defined at higher temperatures. An elastic peak develop in the cubic phase at ζ=0.22 and increases in intensity as TM is approached. However, the dispersion curves show no anomaly at this particular wavevector, in marked contrast to the lattice dynamic studies of other systems exhibiting Martensitic transformations.

Disciplines

Metallurgy

Comments

PHONON PRECURSORS TO THE HIGH TEMPERATURE MARTENSITIC TRANSFORMATION IN Ti$_{50}$Pd$_{42}$Cr$_8$

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Topic: Martensitic Transformation

Abstract. Inelastic neutron scattering measurements were carried out on the Ti$_{50}$Pd$_{50-x}$Cr$_x$ alloy, which has the potential for being a high temperature shape memory material. For x=0, the transformation temperature is ~800K and for the composition studied (x=8 at.%) T$_M$ ~ 400K. The majority of the measurements were performed in the parent, β-phase, up to 873K. Most of the phonons propagating along the three symmetry directions [ζ00], [ζζζ], and [ζζ0] were well defined with the exception of the [ζζ0] transverse acoustic mode with displacements along the [-ζζ0] corresponding to the C'=1/2(C$_{11}$-C$_{12}$) elastic constant. These phonons are well defined for small ζ, but for ζ > 0.15 they are strongly overdamped near the transition temperature, but become better defined at higher temperatures. An elastic peak develop in the cubic phase at ζ=0.22 and increases in intensity as T$_M$ is approached. However, the dispersion curves show no anomaly at this particular wavevector, in marked contrast to the lattice dynamic studies of other systems exhibiting Martensitic transformations.

1. INTRODUCTION

Shape memory alloys (SMA) generate tremendous interest, in both, the scientific and technological area. The applications of SMA are nearly limitless and an excellent review was recently published [1]. However, most of the applications are within ± 100 C of room temperature, since this corresponds to the transformation temperatures of most SMA. There is an interest in SMA with transition temperatures approaching 1000 °C since it would open an entirely new area of applications in engine components. Three prime candidates for high temperature SMA are Ti-Pd, Nb-Ru, Ta-Ru alloys which have transformation temperatures near 800 K [2], 1400 K [3], and 1700 K [3], respectively. In this report we study the precursor phonon effects of the Martensitic transformation of a TiPd alloy by inelastic neutron scattering. A preliminary report has previously been published [4].

TiPd has been know to undergo displacive Martensitic transformation near 800 K from an ordered B2 (CsCl-type structure (Pm3m = O$_h^1$) to a 2H (B19) Martensite on cooling[2]. Upon alloying with various transition metals the transformation temperatures decreases and various intermediate incommensurate phases and different types of Martensites appear. The most extensive study of the phase diagram was by Mateeva[5] and reported by Enami [6]. Most of the work focused on the structure of the Martensitic phase, However a detailed study of TiPd-Cr investigated the modulated cubic structure that appears upon cooling by means of electron diffraction (ED) and high resolution transmission electron microscopy (HRTEM) for various Cr substitution [7].

3. RESULTS AND DISCUSSIONS

We chose to study the Ti$_{50}$Pd$_{42}$Cr$_8$ alloy with a Ms temperature around 400K, because it is a more accessible temperature than the 50-50 alloy. The mosaic of the crystal in the cubic Austenite phase was better than 25'. The cubic lattice parameter was 3.162Å at T=450 K. The transformation temperatures were determined by carefully monitoring the intensity cubic [100] Bragg peak during a slow cooling and heating of the crystals: M$_s$ = 407±5K and A$_f$= 415±5K. The neutron experiments were performed at the NIST Center for Neutron Research using the BT9 triple axis spectrometer operating in the fixed Ef=14.7meV mode. Pyrolytic graphite (PG) was used for monochromator and analyzer with a PG filter placed after the sample to eliminate higher order neutrons. The collimation chosen was 40'-40'-40'-80', which yielded an energy resolution of ~1.0meV. Most of the measurements were made in the [H,K,0] scattering plane centered around the (1,2,0) Brillouin zone center. This is a very intense Bragg peak of the B2 structure because of Ti and Pd have opposite signs of their scattering lengths.

The phonon dispersion curves measured along the [ζ00] and [ζζ 0] directions in the cubic phase are shown in Fig. 1. Most measurements were made at T= 673K, except the TA2 mode, which was measured at 873K The error bars shown for most of the curves are the estimates of the accuracy, except for the TA2 mode where the error bars represent the width of the phonon response.
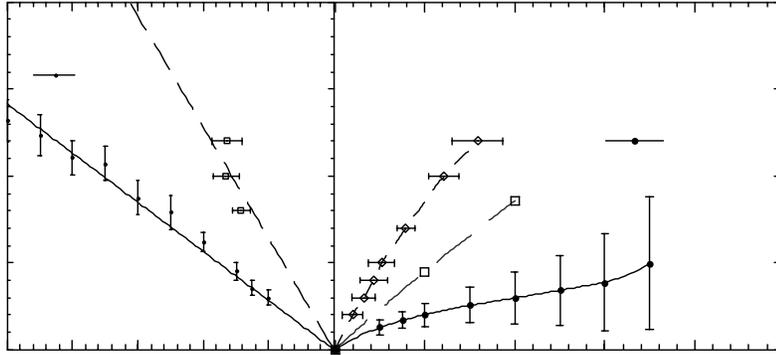

Figure 1: Measured dispersion curve along [00ζ] and [ζζ 0] directions in $Ti_{50}Pd_{42}Cr_8$. Except for the TA2 branch the measurements were performed at 673K. The error bars on TA2 correspond to the full-width-at-half maximum.

The slope of the dispersion curves for ζ ➤ 0 determines the velocity of sound and, knowing the density, the elastic constants. In Table I we list the velocity of sound determined from the limiting slopes. From this

Table I: Velcoity of sound (10^4 cm/sec)
measured at 673 K

	LA	TA1	TA2*
[001]	49±5	22±2	
[110]	44±5	23.±3	12±3

*measured at 873K

we calculate the elastic constants for $Ti_{50}Pd_{42}Cr_8$, using the theoretical density of 7.64 gm/ cm^3;

$$C_{11} = 18.8 \times 10^{11} \text{ dynes/cm}^2$$
$$C_{44} = 3.8 \times 10^{11} \text{ dynes/cm}^2$$
$$C_{12} = 16.7 \times 10^{11} \text{ dynes/cm}^2$$
$$C' = 1/2 (C11-C12) = 1.06 \times 10^{11} \text{ dynes/cm}^2$$

The anomalous branch is the TA2 branch. In the limit of ζ ➤ 0 this corresponds to C'. The atomic displacements associated with his mode correspond to a sliding of [110] planes along the perpendicular [1-10] direction. Zener showed long ago that this mode can lead to instability in the bcc-type lattice. Nearly allbcc-based martensites exhibit strong temperature dependent anomalies and diffuse scattering associated with this branch [8].

Let us now discuss the temperature dependent behavior of this transverse mode and the associated diffuse scattering. Figure 2 shows intensity contours of the elastic scattering measured along the transverse [110] direction from the (1,2,0) Brillouin zone center for Ms<T=450K. There is a ridge of scattering with a pronounced peak at ζ =(0.22,0.22,0) plus diffuse scattering near ζ =0 that is strongly temperature dependent.
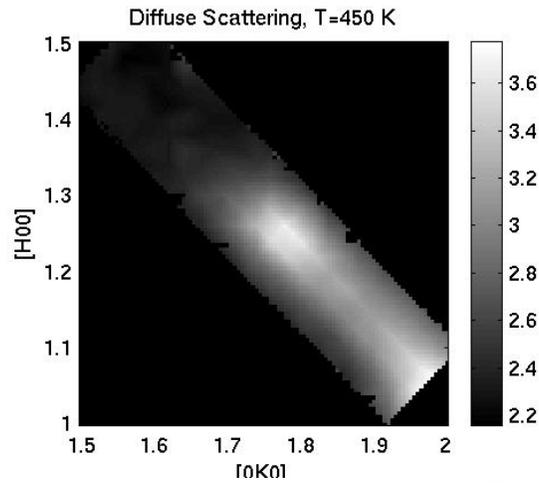
Figure 2. Intensity contours along the transverse direction at T-450K measured from the (1,2,0) Brillouin zone center.

A scan along this direction from the (1,2,0) zone center to the (0,3,0) zone center is shown in Fig 3a
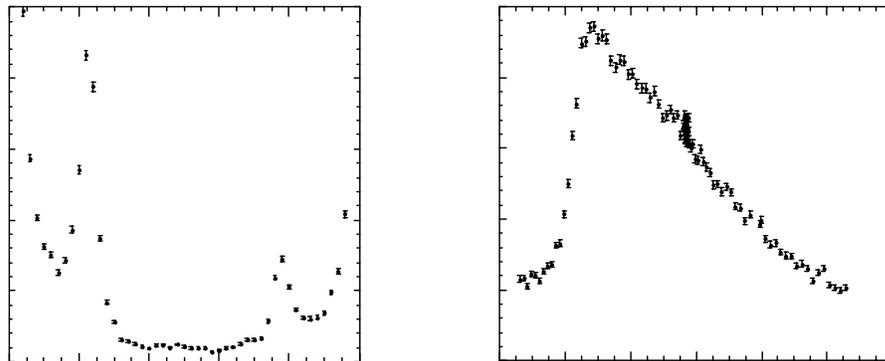

Figure 3. (a) An elastic scan measured from (1,2,0) to (0,3,0) at T=450K. (b) Temperature dependence of q=0.22 diffuse peak.

Two peaks are seen at the reduced value of ζ =(0.22,0.22,0) which implies that a structural modulation exists within the cubic phase. Fig 3b shows the intensity of the peak as a function of temperature. It's position remains the same and the peak intensity increases upon cooling and exhibits a maximum at M_s, but is still present in the low temperature martensite phase but at a reduced level. The linewidth at high temperatures is nearly 5 times the instrumental resolution and it is nearly temperature independent. Its width would correspond to correlations along [110] of about 15 unit cells. The ζ =.22 value of the modulation is smaller than ζ =.29 measured in electron diffraction studies [7]. This discrepancy could be due to uncertainties in composition In this series of measurements on a single crystal we were unable to determine the low temperature structure, but it is reported to be B19 (2H) [2]

We also performed inelastic measurements of the TA2 branch at several temperatures. Fig 4 shows scans at two temperatures for two different ζ values. In Fig. 4a, ζ = 0.1, the peak is well defined at the highest temperature of the measurement, T=873K. As Ms is approached, at T=410K the peak has broadened considerable and reduced its energy and a strong elastic peak developed. Figure 4b shows the spectra along the same direction taken at ζ = 0.25. For the high temperature, the peak is very broad clearly shows a maximum. This indicates how quickly the broadening of the peak increases with q which is also shown in Fig. 1 by the error bars on the TA2 branch. On cooling, a peak is no longer present and an intense elastic peak has developed that is shown in the contour plot of Fig. 2. A
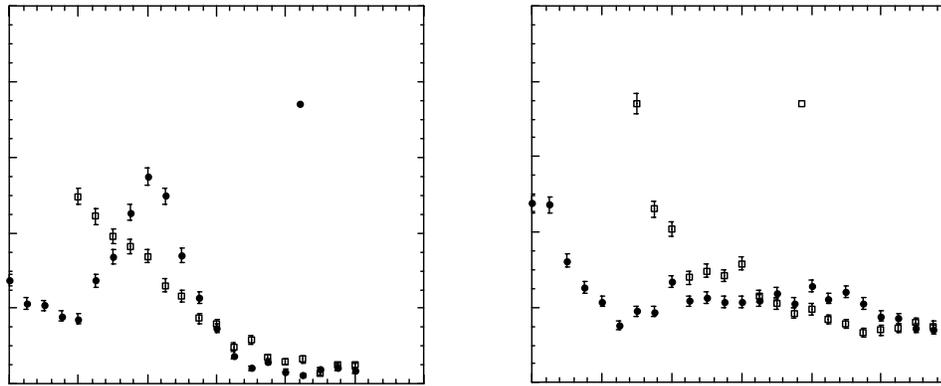
Firgure 4. Inelastic spectra measured at Ms<<T=873K and Ms~T=410K for [ζζ0] TA2 for (a) ζ=0.1 and (b) ζ=0.25.

surprising feature in this TiPd systems, when compared to other martensites, is that the phonon dispersion curve does not exhibit any pronounced anomaly at the same q-value as the appearance of the elastic diffuse peak. The line-broadening extends throughout the Brillouin zone.

In summary i)we deduced the elastic constants from the measurement of the phonon dispersion curves in $Ti_{50}Pd_{42}Cr_8$, ii) measured a modulation of ζ=.22 in the parent cubic phase and iii) observed an anomalous broadening of the [ζζ0] TA2 phonon branch which is strong temperature and ζ dependent.

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