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Exploring the role of electronic structure on photo-catalytic behavior of carbon-nitride (C3N4) polymorphs

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Abstract
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Keywords
Density-functional theory, graphene, semiconductor, band gap, photocatalysis

Disciplines
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Comments

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Keywords: Density-functional theory, graphene, semiconductor, band gap, photocatalysis
1. Introduction

Conjugated carbon-nitride (C$_3$N$_4$) polymers have drawn broad interdisciplinary attention as metal-free and visible-light-responsive photo-catalysts for solar-energy conversion and environmental remediation [1–3], especially for their appealing electronic dispersion, high physicochemical stabilities, and earth-abundant elements. Though experimental works on graphitic C$_3$N$_4$ with its photocatalytic performance. Among other experimental techniques, combination of graphitic-C$_3$N$_4$ with three dimensional hydrogel framework, exfoliation into nanosheets, construction of graphitic-C$_3$N$_4$ based heterojunction systems are widely investigated [4,5]. Besides experimental trials, a plethora of first principles calculations are also aimed to improve the photocatalytic activity of graphitic-C$_3$N$_4$ and its derivatives [6,7]. Notably, doping (e.g., by S or Si) or hybridization with other photocatalyst enhances the photo-reactive performance of graphitic-C$_3$N$_4$ [8–9]. The large band gap tunability makes C$_3$N$_4$ derivatives potential candidates for CO$_2$ capture, control of pollutants [10], water splitting, or energy-storage devices [10,11].

Carbon-nitride is found in amorphous and various crystalline forms [12,13], with distinctly different characteristics (similar to two crystalline forms of carbon: graphite and diamond). However, first-principles electronic-structure and optical properties of the known C$_3$N$_4$ polymorphs are sparse. Using density-functional theory (DFT) methods, we investigate the structural, electronic, and optical properties of C$_3$N$_4$ polymorphs, specifically: (i) α, (ii) β, (iii) cubic, (iv) γ (spinel), (v) $t$−$g$(AA) [AA-stacked graphitic-triazine], (vi) $t$−$g$(AB) [AB-stacked graphitic-triazine], and (vii) $h$−$g$ [AA-stacked graphitic-heptazine].

The LDA or GGA exchange-correlation (XC) functionals [14,15] are known to provide a reasonable account of structural properties, but show limitations and often miserably fail in predicting band gaps [16–18]. To address C$_3$N$_4$ polymorphs, we use our recently developed van Leeuwen-Baerends (vLB) corrected LDA (LDA+vLB) exchange [19] implemented both in tight-binding linearized muffin-tin orbital (TB-LMTO) [20,21] and in full-potential $N^\text{th}$-order muffin-tin orbital (FP-NMTO) [22]. The LDA+vLB is similar to a modified-LDA [23] and mimics the orbital-dependent, exact-exchange potentials [24] and notably requires a computational cost similar to the semi-local XC, like LDA or GGA. Notably, the advantage of our approach lies in the fact that the LDA+vLB greatly improves the band gap relative to traditional (semi)local functionals, with speed of semi-local functionals and accuracy of advanced hybrid-functional (e.g. HSE) or GW functionals [19,24–28]. Our LDA+vLB calculated band-structure of 2-dimensional (2D) graphene and bulk-Si show good agreement to experiments [29] (see supplement), and with hybrid-functionals (HSE) [30] as implemented in the Quantum-Espresso (QE) package [31,32].

Considering the fact that most semi-local functionals underestimate band-energies and band gap [16,18], we employ self-consistent QE-HSE to calculate optical properties, work functions, and photocatalytic behavior of C$_3$N$_4$ polymorphs. We calculate work-functions of well-studied 2D-graphene [33] and bulk-Si [34] to showcase agreement between prediction and experiments. Our study of C$_3$N$_4$ polymorphs finds better optical and photocatalytic behavior of γ-C$_3$N$_4$ compared to other polymorphs. However, we find γ-phase exhibits a dynamical instability, i.e., unstable phonon modes with imaginary frequencies. Therefore, we theoretically assess the structural stability versus pressure to find the dynamical stability of γ-phase, and then fully detail the unstable and stable phonon dispersion. Under pressure, the dynamically stable γ-C$_3$N$_4$ also satisfies the Born structural stability criteria for the elastic constants of the structure. Our notable finding is that the dynamically stable γ-C$_3$N$_4$ (under applied hydrostatic pressure) also shows much improved photocatalytic behavior.

In this systematic study, we established the robustness of our methodology and then we explored the relevant properties of C$_3$N$_4$ polymorphs such as band gap, optical properties, and photocatalytic behavior. Our study reveals the possibility of γ phase as better photocatalytic material than graphitic phases, which are traditionally found to be a suitable photo-catalytic material. Beside this, the scheme can further be applied to the appropriate derivatives of C$_3$N$_4$ polymorphs and other emergent semiconducting materials.

2. Computational Method

FP-NMTO electronic-structure: Choice of a minimal basis set is always tricky, but tight-binding and FP-NMTO [22] methods provide that platform. The FP-NMTO handles the orbital (l-dependent) and $m$-dependent downfoldings independently, which is very useful for $sp^3$-hybridized systems, where $p_z$ orbitals
behaves differently than $p_x$ and $p_y$. As most semi-local functionals fail to estimate the correct band gap of semiconducting materials, we predict band gaps of C$_3$N$_4$ polymorphs using an in-house developed LDA+vLB exchange with von-Barth-Hedin correlation for solids [19,25], as implemented in FP-NMTO [22]. The LDA+vLB potential for given atomic densities ($\rho$) can be written as

$$v_{\alpha\beta}(r) = [v_{\alpha\beta}^{LDA}(r) + v_{\alpha\beta}^{vLB}(r)] + v_{\alpha\beta}^{LDA}(r),$$

with $v_{\alpha\beta}^{vLB}(r) = -\beta z^{1/3}(r) \frac{z^2}{1 + 3\beta z \sinh^{-1}(z)}$ \hspace{1cm} (2)

Here, $z = |\nabla \rho(r)|/|\rho|^{1/3}(r)$ and $\beta = 0.05$.

The LDA+vLB (here onwards we use LDA+vLB instead of FP-NMTO-LDA+vLB) provides improvement to the exchange-potential [27] and better band gaps than other semi-local functionals as gradient correction naturally provides the self-interaction correction (SIC) to the LDA without need of SIC-LDA [36]. Most semi-local functionals fail due to the wrong asymptotic behavior at $r \to 0$ and $r \to \infty$ limits, whereas LDA+vLB produces asymptotically-correct, Coulomb-like $(-1/r)$ behavior [19,37].

The LDA+vLB energies are soft-consistently converged to $10^{-6}$ Ry/atom. We use Anderson method [38] to mix charge densities. For $k$-space integration via the tetrahedron method, we use $k$-mesh of $(6 \times 6 \times 6)$ for $\alpha$; $(6 \times 6 \times 10)$ for $\beta$; $(4 \times 4 \times 4)$ for both cubic and spinnel (y); $(6 \times 6 \times 8)$ for t-$g$ (AA); and $(6 \times 6 \times 4)$ for t-$g$(AB) phases.

**QE-HSE Optical & Work-Function Properties:** We performed band-structure calculations of C$_3$N$_4$ polymorphs using HSE as implemented in the QE plane-wave basis [31,32] and compare that with our LDA+vLB predictions (also see supplement). The LDA+vLB predicted band gaps show very good agreement (see Results) with hybrid-functional and experiments. In addition, the self-consistent QE-HSE is used to calculate the complex dielectric tensor $\epsilon_{\alpha\beta}(\omega)$ within the random phase approximation to analyze the work function and photocatalytic behavior of C$_3$N$_4$ polymorphs [31,32]. For details see supplement (Table S1).

The dielectric tensor $\epsilon_{\alpha\beta}(\omega)$ can be defined as:

$$\epsilon_{\alpha\beta}(\omega) = 1 + \frac{4\pi^2}{3\hbar^2} \sum_{n_\alpha n_\beta} \left( \frac{\langle k_{\alpha} | \phi_{n_\alpha} | \phi_{n_\beta} \rangle \langle \phi_{n_\beta} | \phi_{n_\alpha} \rangle}{(E_{k_{\alpha}} - E_{k_{\beta}})^2} + \frac{\langle k_{\alpha} | \phi_{n_\alpha} | \phi_{n_\beta} \rangle \langle \phi_{n_\beta} | \phi_{n_\alpha} \rangle}{(E_{k_{\beta}} - E_{k_{\alpha}})^2} \right), \hspace{1cm} (3)$$

where $\Gamma$ is an adiabatic (inter-broadening) parameter tending to zero. To retain a finite lifetime of all excited-states, we have introduced small positive value of $\Gamma$ to produce an intrinsic broadening to all excited states. The imaginary part of $\epsilon'_{\alpha\beta}$ is found first and real part $\epsilon''_{\alpha\beta}$ is calculated using Kramers-Kronig relation.

$$\epsilon''_{\alpha\beta}(\omega) = 1 + \frac{2}{\pi} \int_0^{\infty} \frac{\omega \epsilon''_{\alpha\beta}(\omega')}{\omega^2 - \omega'^2} d\omega' .$$

Real and imaginary parts of the dielectric function are used to calculate optical conductivity, refractive index and absorption coefficient [39].

**Dielectric tensor:**

$$\epsilon_{\alpha\beta} = \epsilon'_{\alpha\beta} + i\epsilon''_{\alpha\beta} \hspace{1cm} (4)$$

**Optical Conductivity:**

$$Re[\sigma_{\alpha\beta}(\omega)] = \frac{\omega \epsilon''_{\alpha\beta}(\omega)}{4\pi} \hspace{1cm} (5)$$

**Refractive Index:**

$$\mu_{na} = n_{oa} + ik_{a} \hspace{1cm} (6)$$

$$n_{oa}(\omega) = \sqrt{\frac{\epsilon_{oa}(\omega)\epsilon_{\alpha\alpha}(\omega)}{2}} ; k_{oa}(\omega) = \sqrt{\frac{\epsilon_{oa}(\omega)\epsilon_{\alpha\alpha}(\omega)}{2}}$$

**Absorption Coefficient:**

$$A_{oa}(\omega) = \frac{2\omega k_{oa}(\omega)}{c} \hspace{1cm} (7)$$

**Work-function calculation:** For photocatalytic activity, the valence-band maximum (VBM) and conduction-band minimum (CBM ) position are equally important as band gaps. Therefore, the accurate calculation of work-function ($\Phi$) becomes important, which is defined as the energy required to remove an electron from a slab surface. The $\Phi$ is calculated by taking the difference of vacuum energy ($E_{\text{vac}}$) with respect to Fermi-energy ($E_F$), i.e., $\Phi = V_{\text{vac}} - E_F$ [40]. For semi-infinite surfaces (with sufficient number of layers that reproduces bulk behavior), $\Phi$ is the difference of vacuum energy for slab and the Fermi energy of the bulk ($\Phi = V_{\text{vac}} - E_{\text{bulk}}$). This needs two different calculations, (a) bulk, and (b) slabs (semi-infinite surface). The electrostatic potential within the interstitial region in (a) and (b) works as a common parameter, which should be same for (a) bulk and (b) slab in order to compare the energy levels. The matching bulk and slab values standardizes the energy levels. Also, the presence of microscopic fluctuations of electrostatic potentials in the interstitial region require macroscopic averaging of average planar potential with a window size equal to the inter-slab separation. The average planar potential is the average over the plane parallel to the surface-plane and the direction of microscopic averaging (of planar potential) is perpendicular to that plane [41,42]. Here, the difference of the macroscopic averages for bulk and slab ($\Delta \Phi = V_{\text{bulk}} - V_{\text{slab}}$) works as a constant correctional shift to $E_{\text{bulk}}$ that brings equivalence of energy levels of these two separate self-consistent calculations. Finally, $\Phi$ can be evaluated as $\Phi = V_{\text{vac}} - E_{\text{bulk}} + (V_{\text{bulk}} - V_{\text{slab}})$. This is how, we
can get rid of the quantum-size effect in designing slabs for work-function calculation. Furthermore, for the semi-infinite systems, the band gap is not much dependent on the surface, and for intrinsic semiconductors the $E_F^{bulk}$ lies at the mid-point of VBM and CBM; therefore, the exact position of energy bands with respect to vacuum energy can be identified. Each slab is created with four layers and inter-slab vacuum is of 15 Å. We have provided validation test of our approach on 2D-graphene [33] and bulk-Si [34].

**Phonons for dynamical stability:** To check the dynamical stability in terms of phonons for γ-C$_3$N$_4$, we use first-principles density-functional-perturbation theory (DFPT). The DFPT is a straightforward approach that permits phonons to be calculated by perturbing atomic positions on $2 \times 2 \times 2$ supercell (112 atom/cell) of the original 14 atom/cell. All the atomic coordinates are relaxed up to $10^{-6}$eV/Å. The atoms are displaced by 0.01 Å from their equilibrium positions to calculate force constants. We use these force constants to calculate phonon dispersion along the high-symmetry direction of the Brillouin zone of γ-phase [44].

3. Results and Discussion

We discuss band-structure of 2D graphene [45], and bulk-Si [29] calculated using FP-NMTO-LDA+vLB to showcase the applicability of our approach and compared with other theory and experiments. We then study allotropes of C$_3$N$_4$, both non-graphitic and graphitic phases, for structural and electronic properties.

**Graphene:** With a 2D hexagonal structure, graphene is very well-studied form of carbon and considered as the building block of many carbon allotropes [46, 47]. The calculated band-structure (Fig. 1 left-panel) shows zero band gap at the K-point with Dirac (Fermionic) linear dispersion. Here, the zero band gap is attributed to the sub-lattice symmetry of the 2D graphene.

**Bulk Si:** The predicted band gap of bulk-Si (1.25 eV) using LDA+vLB (Fig. 1 right-panel) is also in good agreement with hybrid-functional (1.24 eV) and experiments (1.17 eV) [29]. Notably, we reproduce the indirect (Γ−X) nature of Si band gap [25], which is not found by standard semi-local XC, for example LDA/GGA [14, 15].

3.1. C$_3$N$_4$ polymorphs

**Non-Graphitic Phases – structural property:** An account of α, β, γ (spinel), and cubic phases of C$_3$N$_4$ has been presented earlier [12, 18]. The structures are shown in Fig. 2. The α-C$_3$N$_4$ phase (Fig. 2(a)) consists of the corner sharing CN$_4$ tetrahedron along with pyramidal NC$_3$ in spheroidal cavities, suggesting that C and N are $sp^3$ and $sp^2$ hybridized, respectively [12]. The α-phase has 4 formula units (f.u.) per cell with 28 atoms. The β-C$_3$N$_4$ (Fig. 2(b)) has 6-, 8- and 12-membered rings of alternating C and N atoms [18]. The β-phase has two inequivalent N atoms: (I) one N has three equidistant nearest carbon atoms with the C-N-C angles ~120°, which shows that the C and N are $sp^2$ hybridized; and (II) other N is bonded to three non-planar carbon atoms in $sp^2$-$sp^3$ fashion. The γ-phase (Fig. 2(c)) has a spinel structure, which is known as cubic-modification of boron-nitride (c-BN) [48]. Two carbon atoms bond octahedrally to six nitrogen atoms, while the third carbon atom bonds tetrahedrally to four nitrogen atoms [49]. The tetrahedrally and the octahedrally bonded C-N are arranged alternatively in a 1:2 ratio. These are sequentially connected to one another at the N–corners. The cubic-C$_3$N$_4$ (Fig. 2(d)) is a high-pressure modification of Willemite mineral (Zn$_2$SiO$_4$) [12], formed by replacing the O with N and Zn and Si with C of Zn$_2$SiO$_4$ structure, so N and C are $sp^3$ hybridized.

The structural properties of C$_3$N$_4$ polymorphs are summarized in Table. 1. The PBE, PBEsol and LDA+vLB calculated structural parameters are in good agreement [50, 51]. Our predicted lattice constants are similar to those of experiment for α (a=6.5680 Å, c=4.7060 Å) and β phases (a=6.4017 Å, c=2.4041 Å) [50, 52]. We also evaluate bulk moduli by fitting lattice constant versus total energies to the Birch-Murnaghan equation of state [53, 54], which indicates that C$_3$N$_4$ undergoes uniform compression under applied hydrostatic pressure. The agreement between LDA+vLB and PBEsol calculated numbers establishes effectiveness and the utility of vLB-modified XC functional [19]. The electronic and optical property calculations are carried out using optimal lattice parameters
Table 1: Comparison of lattice constants and bulk moduli of C$_3$N$_4$ polymorphs calculated using vLB, PBE and PBEsol functionals. Experimental lattice constant of $\beta$ phase [50] are compared with calculated values. Bulk moduli calculated with vLB and PBEsol are in good agreement.

<table>
<thead>
<tr>
<th>Polymorphs</th>
<th>$\alpha$</th>
<th>$\beta$</th>
<th>$\gamma$</th>
</tr>
</thead>
<tbody>
<tr>
<td>XC→</td>
<td>vLB</td>
<td>PBE</td>
<td>PBEsol</td>
</tr>
<tr>
<td>$c$ (Å)</td>
<td>4.6296</td>
<td>4.7382</td>
<td>4.7099</td>
</tr>
<tr>
<td>$V$ (Å$^3$)</td>
<td>165.766</td>
<td>175.242</td>
<td>172.041</td>
</tr>
<tr>
<td>$R_0$ (GPa)</td>
<td>452.676</td>
<td>419.362</td>
<td>453.841</td>
</tr>
</tbody>
</table>

Figure 2: Crystal structure of carbon-nitride (C$_3$N$_4$) polymorphs: (a) $\alpha$-phase (trigonal; space group: P3$_1$c (159)); (b) $\beta$-phase (hexagonal; P6$_3$/m, 176); (c) $\gamma$-phase (spinel; Fd3m, 227); (d) C-phase (cubic; I43d, 220). See phase stability in Fig. S3.

Non-Graphitic Phases – electronic structure: Band structure is one of the most stringent tests to detail the materials physics. For example, Si, calcite and Cu have similar electron densities, but they have very different physical and electronic properties. This drives us to understand the band structure versus energy of C$_3$N$_4$ in detail, shown in Fig. 3 with a zero of energy at the VBM (E$_V$) and CBM (E$_C$). The LDA+vdW band gaps for $\alpha$, $\beta$, $\gamma$ and cubic (C) phases are 5.81, 5.32, 1.81 and 4.23 eV, respectively. The band gaps of $\alpha$ and $\beta$ phases are indirect in nature, while $\gamma$ phase shows a direct gap. We find good agreement between our predictions and more advanced DFT techniques (e.g., HSE and G$_0$W$_0$), see Table 2.

The projected and total density of states (DOS) of N and C are shown in Fig. 3 for each of the C$_3$N$_4$ polymorphs. In Fig. 3(a), for the $\alpha$ phase the valence bands (VBs) below $-4$ eV phase are dominated by both N and C, while the bands from $-4$ eV to E$_F$ are mostly from N. This suggests that the electrons of N are loosely bound than those of C. The $\alpha$ phase has a wider band gap, and steeper VB and CB edges, compared to the others. The VB maxima and CB minima in $\alpha$ phase are at $\Gamma$ and M point, respectively. Similar to $\alpha$, we find that $\gamma$ ($\Gamma$ to $A$ to $\Gamma$) and C ($\Gamma$ to $\Gamma$ to $H$) are indirect band gap semiconductors with no contribution from C bands near E$_F$. However, the $\gamma$ phase is a direct ($\Gamma$-$\Gamma$) band gap with minor contribution from C bands near E$_F$. The electronic structures of $\alpha$, $\beta$, $\gamma$ and cubic polymorphs split to form the separate VB, which reflects the presence of smaller asymmetric part of the potential. The presence of asymmetric potential is the reason behind the stronger mixing of low lying N and C states.

Graphitic Phases – structural property: Notably, graphitic phases of C$_3$N$_4$ are considered for next generation visible-light-driven metal-free, non-toxic, earth-abundant semiconductors, which can find applications in energy conversion, hydrogen evolution, sensing and imaging. Earth-abundant graphitic C$_3$N$_4$ possesses excellent electronic band structures, electron-rich properties, and higher stability. The sp$^2$ hybridized planar graphitic C$_3$N$_4$ (Space group: P6m2; #187) in Fig. 4 can be viewed as graphite whose C lattices are partially
Table 2: Calculated band gaps for C₃N₄ polymorphs using (LDA, PBE, vLB, HSE) and comparison with other theories [16–18, 55] (LDA, PBE, HSE, GW, mBJ) and experiments [1, 8, 16, 59]. HSE calculations are done with plane-wave basis [31, 32]. Band gaps using vLB [19], and HSE (this work) shows good agreement with existing HSE, GW and experiments. We also include non-self-consistent GW band gap for γ-C₃N₄ phase calculated with QE package. Band gaps from HSE (1.95 eV), GW (2.01 eV) and vLB+LDA (1.81 eV) are in good agreement.

<table>
<thead>
<tr>
<th>C₃N₄ Polymorphs</th>
<th>Expt.</th>
<th>This Work</th>
<th>Others</th>
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<tbody>
<tr>
<td></td>
<td>LDA</td>
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<td>vLB</td>
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<tr>
<td>α</td>
<td>–</td>
<td>5.27</td>
<td>3.81</td>
</tr>
<tr>
<td>β</td>
<td>–</td>
<td>3.27</td>
<td>3.31</td>
</tr>
<tr>
<td>γ</td>
<td>–</td>
<td>1.19</td>
<td>1.14</td>
</tr>
<tr>
<td>cubic</td>
<td>–</td>
<td>2.86</td>
<td>2.95</td>
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<tr>
<td>t − g(AA)</td>
<td>3.1</td>
<td>[59]</td>
<td>2.79</td>
</tr>
<tr>
<td>t − g(AB)</td>
<td>–</td>
<td>2.66</td>
<td>2.67</td>
</tr>
<tr>
<td>h − g</td>
<td>1.8 − 2.95</td>
<td>[12, 16, 60]</td>
<td>–</td>
</tr>
</tbody>
</table>

Figure 4: Crystal structure of relaxed graphitic-C₃N₄: (a) t − g(AA), (b) t − g(AA) unit cell, and (c) h − g(AA), (d) t − g(AB) unit cell C₃N₄. For relaxed t − g phase, bond angles and lengths vary from (118−124°) and (1.33−1.47 Å), respectively, while, in relaxed h − g phase, bond angles and lengths vary as (117.6−122.2°) and (1.325−1.468 Å).

We prepare triazine with (AA; AB) stacking and heptazine with AA stacking, respectively. The large cavities in t-g(AA), t-g(AB) and h-g(AA) allow neighboring atoms to relax and lead to increased bond lengths and angles compared to other phases (Fig. 4). For t-g(AA), the equilibrium lattice constants are (a, c) = 4.786, 3.758 Å, where bond lengths and angles vary from (1.327−1.463 Å) and (118−124°), respectively. The calculated interlayer spacing of 3.505 Å and lattice constant of t-g(AB) (a, c) = (4.784, 7.007) Å is in good agreement with experiment ((a, c) = (5.041, 6.576) Å) [61]. The bond lengths vary from 1.326−1.468 Å, while the angles range from 117.6−122.2°. Due to larger cavity size compared to t-g(AA) and t-g(AB), the h-g (AA) structure shows larger degree of relaxation, both in bond lengths and angles as shown in Fig. 4(d).

Graphitic Phases – electronic structure: Experimentally thin films of graphitic-C₃N₄ can be produced on different substrates. By controlling the synthesis conditions, the band gap can be tuned from 2.65 to 3.1 eV, which falls in the range of visible light [1, 16, 59]. The LDA+vLB band gap of 3.02 eV for AA (trigine) and 2.95 eV for AB (heptazine) stackings are in close agreement with experimental gap (3.1 eV) [59]. The QE-HSE predicted band gap for AA and AB stackings is 3.21 eV and 3.39 eV, respectively. LDA+vLB predicted band gaps are in excellent agreement with HSE, other theory [56, 57], and other experiments. The t-g(AA) shows indirect (K−A) band gap, whereas t-g(AB) is direct. The LDA+vLB predicted indirect (Γ−M) band gap (2.71 eV) of h-g-C₃N₄ is in good agreement with HSE (2.71 eV), other theory [56, 57], and experiments 2.67−2.95 eV [1, 8, 16, 60].

To identify active sites in C₃N₄ polymorphs, we plot the CBM and VBM charge densities (Fig. 5, right panel). The CBM of t − g(AA), t − g(AB) and h − g(AA) C₃N₄ are mainly composed of N-2p and C-2p states, while the VBM is dominated by N-2p states. Also, strongly localized CBM and VBM states in charge density plots suggest the possibility of low photo-absorption efficiencies, i.e., e−−h+ excitation under visible-light irradiation occurring on the edge of N and C atoms.

In Table 2 we calculated the band gaps from LDA, PBE, vLB, and HSE (present work) with other theory and experiments [1, 16, 18, 56, 59]. The LDA+vLB and QE-HSE predicted band gap for cubic phase is in good agreement with GW, while band gap for α and β phase is higher than GW.

3.2. Effective mass – C₃N₄ polymorphs

The dynamical activity of the charge carriers in semiconductor depends on its mobility, which is inversely proportional to the effective mass (m*). Near the band edges, effective mass is described by m* = ℏ² (∂²E/∂k²), where ∂²E/∂k² is the band-curvature at VBM for holes and
Figure 5: Electronic band-structures, density of states and charge densities (VBM & CBM states) of relaxed graphitic-C\(_3\)N\(_4\) – (a) t–g(AA), (b) t–g(AB), and (c) h–g(AA). The t–g(AA), t–g(AB) and h–g(AA) phases exhibit band gaps of 3.02 eV (K–A; indirect), 2.95 eV (Γ–Γ; direct) and 2.71 eV (Γ–M; indirect), respectively. C-p (red) and N-p (blue) projected density of states (PDOS) are shown. Charge-density isosurface plots are set at 0.035 electrons/Å\(^3\).

Figure 6: QE-HSE calculated (a) Absorption spectra, (b) optical conductivity, and (c) refractive index vs. photon energy for C\(_3\)N\(_4\) polymorphs. (a\(_1\)) Absorption spectra, (b\(_1\)) optical conductivity, and (c\(_1\)) refractive index are also shown in the visible region of light.

3.3. Optical Properties:

Solar energy output is mostly dominated by: (I) ultraviolet (∼5%), visible (∼45%), and infrared (∼50%) region of electromagnetic wave spectrum [64]. The visible-light photocatalysis therefore offers the best opportunity to utilize maximum solar energy [65]. Except for γ and h – g, the C\(_3\)N\(_4\) polymorphs possess relatively wide band gaps (see Table 2). The optimal gap of γ and h – g C\(_3\)N\(_4\) [66] makes them promising polymeric semiconductors suited for visible-light absorption.

In Fig. 6, we plot absorption coefficient, optical conductivity, and refractive index vs photon energy. Plots of effective number of electrons participating in interband transition (\(n_{eff}\)) shows that the (0–15 eV) range of photon energy is justified (Fig. S4). However, our region of interest is visible light from 1.59 eV (780 nm) to 3.18 eV (390 nm), in Fig. 6 (a\(_1\), b\(_1\), c\(_1\)). The optical spectra (obtained from the imaginary part of the dielectric function) contains the information of character and number of occupied and unoccupied bands. These together decide the accuracy of optical properties at higher excita-

<table>
<thead>
<tr>
<th>Polymorphs</th>
<th>(m^*)</th>
<th>(e^*)</th>
<th>(h^*)</th>
</tr>
</thead>
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</tr>
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<tr>
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</tr>
<tr>
<td>cubic</td>
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<td>-0.065</td>
<td></td>
</tr>
<tr>
<td>t – g(AA)</td>
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<td>-0.044</td>
<td></td>
</tr>
<tr>
<td>t – g(AB)</td>
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<td>-0.026</td>
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</tr>
<tr>
<td>h – g(AA)</td>
<td>0.045</td>
<td>-0.081</td>
<td></td>
</tr>
</tbody>
</table>

Table 3: Effective mass (\(m^*\)) of C\(_3\)N\(_4\) polymorphs in electron rest mass (\(m_0\)) units.

at CBM for electrons. The \(m^*\) for C\(_3\)N\(_4\) polymorphs is shown in Table 3. The calculated \(m^*\) in graphitic phase is in good agreement with the available results [62, 63] and we find electron \(m^*\) of γ phase is much lower than graphitic phases. This indicates that the electron mobility in γ phase is easier. Interestingly, between the two types of stacking in t-graphitic phases, the A-B alternate stacking provides lower electron and hole effective mass and higher mobility than the AA phase. We attribute this to the lower symmetry of AB phase.
tion energies.

The absorption coefficient of a material determines the spatial region in which most of the light is absorbed. The absorption edge extent of the $\gamma$ and $h - g$ phases is relatively larger than other polymorphs due to narrower gaps of 1.81 and 2.71 eV, respectively [66]. The onset of absorption for $\gamma$ and $h - g$ phases are located at roughly 2.0 and 3.0 eV, respectively, very close to the their band gaps. Also, the structural change induced by relaxation in Fig. 4(d) for $h - g$ are manifested by more pronounced changes in the optical absorption in Fig. 5(a) & (a1) (AA) as compared to $t$-$g$ (AA) and $t$-$g$ (AB). The degree of relaxation and band gap decides the absorption-onset (AA) as compared to $t$-$g$ (AA) and $t$-$g$ (AB). Interestingly, both $\gamma h$-g phase BE (AO) and band-edge absorption (BE), for example, for $h$-$g$ phase BE/AO is at 3.0 eV smaller than $t$-$g$ (AA) and $t$-$g$ (AB). Also, the degree gap materials provide better photons to electron-hole pair conversion useful for efficient electro-optical devices – the foremost reason behind higher optical absorbance in $\gamma$ phase, as shown in Fig. 6(a).

The accurate description of the band gap and band-positions from hybrid-functionalss assure the reliability of our prediction of optical properties. We compared the QE-HSE calculated optical conductivity of C$_3$N$_4$ polymorphs in Fig. 6(b). Our results shows that the conductivity starts with a gap, indicating the semiconducting character of each polymorph. The optical conductivity is zero below a certain energy value, which is consistent with the respective band gaps in Table 2 of C$_3$N$_4$ polymorphs: $\gamma$-phase a promising candidate for photovoltaic application. Refractive index is shown in Fig. 6(c) as a function of photon energy (wavelength: in Fig. 7(a), $a_t$, $b_1$, $c_1$). Physically the calculated optical gap corresponds to the photon energy at which the imaginary part of the refractive index, $k$, becomes non-zero. We found that the refractive index for C$_3$N$_4$ polymorphs vary from 1.4 to 1.9 for the visible range of light 390 – 780 nm, showing sensitivity of the refractive indices to crystal structure. The optical absorption spectra are directly connected to the imaginary part of the refractive index, see Eq. (4).

3.4. Work Function

Monolayer graphene and bulk-Si: The work function, which limits the performance of devices, is an important characteristics of semiconductors that should be taken into considerations. First, we cross-validate QE-HSE method by calculating work-function of well-studied 2D-graphene, and bulk-Si. The $3 \times 3 \times 1$ and $2 \times 2 \times 1$ supercells are chosen for accurate calculations of 2D-graphene and bulk-Si, respectively. The inter-slab vacuum of 15 Å is chosen along (001) to avoid periodic image problem. In Fig. 7, we show the calculated work-function of 4.39 eV for 2D-graphene. The $\phi_{BE}$ is in good agreement with $\phi_{Exp}$=4.56 eV [33]. In Fig. 7, the conduction-band minima (CBM) and valence-band maxima (VBM) for Si are plotted considering $E_{vac}$ = 0. The calculated work-function of 4.43 eV for Si is in very good agreement with experiments (4.87 eV) [34].

C$_3$N$_4$ polymorphs: At ambient conditions (i.e., at normal (room) temperature and one atmospheric pressure), the thermodynamic voltage for water splitting is 1.23 V [67, 68]. But, to supply the required photovoltage, a photo-electrochemical cell with single illuminated electrode needs gaps greater than 1.6 eV [69]. In Table 2, we show that the gap in $\gamma$ and graphic phases is suitable for photoabsorption. To initiate water-redox reaction in semiconductors, optimal gaps, as well as band-edge (conduction and valence) positions of VBM and CBM states are important. The valence-band maxima should be more positive ($E_{H_2O/O_2}$=1.23, 0.81 V for pH=0, 7) with respect to normal hydrogen electrode (NHE) than the water oxidation level, whereas the conduction-band minima should be more negative ($E_{H+/H_2}$=0, -0.41 V for pH=0, 7 vs. NHE) than the...
H-production potential:

\[
H_2O + 2h^+ \rightarrow 2H^+ + \frac{1}{2}O_2
\]

\[
2H^+ + 2e^- \rightarrow H_2
\]

(8)

In Fig. 8, we illustrate valence and conduction band positions and work function for \(t - g(AA), t - g(AB), h - g\), and \(\gamma\) phase for (100), (110) and (111) planes. The VBM for graphitic phase is positive with respect to \(H_2O/2 - p_H = 0\) level while CBM are negative compared to \(H^+/2 - p_H = 0\) level in the entire \(pH\)-range (0 to 7). The work function of the triazine phases depends largely on the choice of stacking: alternate (AB) stacking increases the band gap and also enhances the work function potential. The band-position in both t-g (AA) and t-g (AB) makes them good candidate for photocatalyst -C3N4. Among the three illu- 
minated surface conditions of \(\gamma\) phase, i.e., (100), (110), (111) (see Fig. S5 in supplement), the VB and CB band range of (110) surface is favorable at \(pH=0\). Our study suggest \(\gamma\)-C3N4 as a potential candidate for photocata- 
ytic application, a new and efficient photocatalysts from spinel group \[73\].

3.5. Pressure effect on key properties of \(\gamma\)-C3N4:

Optimal band gap (1.95 eV) for \(\gamma\)-phase indicates its suitability for optical and photocatalytic application. Yet, most C3N4-poly morphs suffer from structural instabilities \[74\] \[75\], i.e., all these have phonons with imaginary frequencies at zero applied pressure. It is noteworthy that metastable \(\gamma\)-C3N4 is already been experimentally observed by Andrade et. al. \[76\]. The measurements found the emission signals for \(\gamma\)-C3N4, but signals disappear after few scanning tunneling microscopy (STM) scans. As the experiments were performed at atmospheric conditions, where oxygen is always present and STM scans can generate high-energy photons, which possibly oxidizes the \(\gamma\)-C3N4 phase. This indicates towards metastability of \(\gamma\)-C3N4 at ambient conditions (i.e., at room temperature and one atmospheric pressure). Therefore, we investigate the pressure effect on \(\gamma\)-phase. In Fig. 9 we plot phonons at ‘zero’ pressure (top) and at 275 GPa (bottom). The 0 GPa phonons show imaginary frequencies that indicates a dynamical instability in \(\gamma\)-phase. However, under hydrostatic pressure (above 275 GPa) all imaginary frequencies in Fig. 9 disappear, i.e., above a critical pressure \(\gamma\)-phase become dynamically stable.

The dynamically stable \(\gamma\)-phase also satisfies the Born stability criteria \[77\], i.e., (a) \(C_{11} - C_{12} > 0\); (b)
C_{11} + 2C_{12} > 0; (c) C_{44} > 0, where C_{11} = 436.5 GPa, C_{12} = 363.09 GPa, and C_{44} = 446.5 GPa. The strong covalent bonding gives large bulk-, elastic-, and shear-modulus of 387.6 GPa, 682.0 GPa and 282.6 GPa, respectively.

To point out the effect of pressure on band-structure, we plot electronic band-structures of γ-C_{3}N_{4} in Fig. 10(a)&(b). For with and without pressure cases, γ-C_{3}N_{4} shows good agreement for conduction band minima and valence band maxima positions, which is important both for calculation of optical properties and band alignment in determining photocatalytic behavior.

The effect of pressure clearly reflects on optical properties of γ-C_{3}N_{4} (Fig. 11). We found an increase in the extent of absorption edge for the γ phase compared to ambient (one atmospheric) pressure (Fig. 11(a)). The absorption edges at ~300 nm for γ phase corresponds to the band gap energy of 3.43 eV. First, the optical absorption range for γ phase increases under pressure, and, secondly, the optical conductivity of γ-C_{3}N_{4} (Fig. 11(b)) also changes from that at “zero pressure” (Fig. 6(b)). A slight change in conductivity compared to “zero-pressure” suggests increased photocatalytic behavior of γ-phase. Refractive index of dynamically stable γ-phase in Fig. 11(c) reduces compared to “zero pressure” case. The calculated optical gap corresponds to the photon energy at which the imaginary part of the refractive index, k, becomes non-zero. We find the refractive index for γ-C_{3}N_{4} changes 2.5 to 2.05 for the visible light range.

In Fig. 12, we compare the VB and CB positions and work function for most favorable surface (110) of γ phase with(out) pressure. Clearly, the photocatalytic range of γ (110)-surface under pressure (275 GPa) increases. We also showcase the effect of water on (110) surface on γ phase and find slight reduction in CB band-edge compared to surface with no water. The CB band-edge becomes slightly less positive compared to pure (110) surface, whereas, the γ (110) phase is now more favorable under pressure due to increased photocatalytic range (pH=0 to 7) in both with and without water case.

The pressure profile for the structural stability of γ-C_{3}N_{4} is significantly higher, however, recent studies presented novel and more practical synthesis routes for new or known high-pressure phases under predictable nonhydrostatic loading [78]. It has been exemplified for silicon that the hydrostatic pressure of 76 GPa could be lowered by 21 times to 3.7 GPa under uniaxial loading. We believe that the similar idea could be employed in future for synthesizing new high-pressure carbon-
nitride materials. To add further discussion on pressure and point out its importance, recently, pressure has helped discover new material(s) and mechanisms \[79, 80\]. Although, the stability of $\gamma$-$C_3N_4$ is an issue \[81, 82\], the use of hydrostatic pressure in this work is an effort to provide theoretical understanding of metastable (structurally unstable) $\gamma$-$C_3N_4$.

Figure 12: At 0 GPa (dynamically unstable) and 275 GPa (stable), the conduction and valence band positions of favorable (110) surface with and without water of $\gamma - C_3N_4$-polymorph with respect to water reduction and oxidation potentials.

4. Conclusion

We present a detailed study of both structural and electronic-structure of $C_3N_4$ polymorphs using a localized-basis (FP-LMTO-LDA++) and a plane-wave basis (QE-HSE) method. The structural properties and band gap of $C_3N_4$-polymorphs predicted from LDA++HSE agrees well with our more accurate QE-HSE results, and available experiments. We did cross-validation tests performed on two diverse cases: (a) 2D-graphene, (b) bulk-Si. The calculated band gaps (graphene, 0 eV; bulk Si, 1.25 eV indirect) agree well with experiments, which shows the predictive reliability and accuracy. To point out important differences between our approach and other calculations done for $C_3N_4$ in literature are two-fold, (a) most calculations are performed using LDA/GGA functionals, which are known to underestimate the band gap; and (b) hybrid or GW functionals were not used in systematic way to study of relevant governing factors for photocatalytic performance of carbon-nitriles. We performed optical-property, and work-function calculation using hybrid-functional (HSE) as implemented within Quantum-Espresso package. Beside lattice parameters and band gaps, work-function calculations are found to be in very good agreement with the available experiments. Our calculations show optimal band gap of 1.95 eV as well as higher optical conductivity and carrier mobility for $\gamma$-$C_3N_4$ than other polymorphs, which makes $\gamma$-$C_3N_4$ a suitable candidate for photocatalytic applications. The work function in $\gamma$-$C_3N_4$ shows an orientation dependence with (110) surface as most favorable orientation over (100) and (111). In spite of favorable electronic-structure and optical properties, $\gamma$-phase shows a dynamical instability (unstable phonons) at standard temperature and pressure. We show that $\gamma$-$C_3N_4$ is structurally stabilized under hydrostatic pressure (275 GPa) and leads to a significant improvement in photocatalytic behavior with respect to water reduction and oxidation potentials. We could also identify the active sites in $C_3N_4$-polymorphs from charge-density plots as strongly localized states, usually suggesting low photo-absorption efficiencies. Our study offers $\gamma$ phase as a new candidate to explore for optical and photocatalytic applications, which may further ignite interest in $C_3N_4$ polymorphs within the materials and chemistry community.

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