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## Helical magnetic ordering in Sr(Co<sub>1-x</sub>Ni<sub>x</sub>)(<sub>2</sub>)As-2

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## Helical magnetic ordering in Sr(Co<sub>1-x</sub>Ni<sub>x</sub>)<sub>2</sub>As<sub>2</sub>

### Abstract

SrCo<sub>2</sub>As<sub>2</sub> is a peculiar itinerant magnetic system that does not order magnetically, but inelastic neutron-scattering experiments observe the same stripe-type antiferromagnetic (AF) fluctuations found in many of the Fe-based superconductors along with evidence of magnetic frustration. Here we present results from neutron diffraction measurements on single crystals of Sr(Co<sub>1-x</sub>Ni<sub>x</sub>)<sub>2</sub>As<sub>2</sub> that show the development of long-range AF order with Ni doping. However, the AF order is not stripe type. Rather, the magnetic structure consists of ferromagnetically aligned (FM) layers (with moments laying in the layer) that are AF arranged along *c* with an incommensurate propagation vector of (0 0  $\tau$ ), i.e., a helix. Using high-energy x-ray diffraction, we find no evidence for a temperature-induced structural phase transition that would indicate a collinear AF order. This finding supports a picture of competing FM and AF interactions within the square transition-metal layers due to flat-band magnetic instabilities. However, the composition dependence of the propagation vector suggests that far more subtle Fermi surface and orbital effects control the interlayer magnetic correlations.

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### Comments

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**Helical magnetic ordering in  $\text{Sr}(\text{Co}_{1-x}\text{Ni}_x)_2\text{As}_2$** J. M. Wilde<sup>1,2</sup>, A. Kreyssig<sup>1,2</sup>, D. Vaknin<sup>1,2</sup>, N. S. Sangeetha<sup>1,2</sup>, Bing Li<sup>1,2</sup>, W. Tian,<sup>3</sup> P. P. Orth<sup>1,2</sup>, D. C. Johnston,<sup>1,2</sup> B. G. Ueland<sup>1,2</sup> and R. J. McQueeney<sup>1,2</sup><sup>1</sup>Ames Laboratory, U.S. DOE, Ames, Iowa 50011, USA<sup>2</sup>Department of Physics and Astronomy, Iowa State University, Ames, Iowa 50011, USA<sup>3</sup>Neutron Scattering Division, Oak Ridge National Laboratory, Oak Ridge, Tennessee 37831, USA

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$\text{SrCo}_2\text{As}_2$  is a peculiar itinerant magnetic system that does not order magnetically, but inelastic neutron-scattering experiments observe the same stripe-type antiferromagnetic (AF) fluctuations found in many of the Fe-based superconductors along with evidence of magnetic frustration. Here we present results from neutron diffraction measurements on single crystals of  $\text{Sr}(\text{Co}_{1-x}\text{Ni}_x)_2\text{As}_2$  that show the development of long-range AF order with Ni doping. However, the AF order is not stripe type. Rather, the magnetic structure consists of ferromagnetically aligned (FM) layers (with moments laying in the layer) that are AF arranged along  $c$  with an incommensurate propagation vector of  $(0\ 0\ \tau)$ , i.e., a helix. Using high-energy x-ray diffraction, we find no evidence for a temperature-induced structural phase transition that would indicate a collinear AF order. This finding supports a picture of competing FM and AF interactions within the square transition-metal layers due to flat-band magnetic instabilities. However, the composition dependence of the propagation vector suggests that far more subtle Fermi surface and orbital effects control the interlayer magnetic correlations.

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The Fe-based superconductors and their parent compounds [1–4] are prime examples of intertwined structural, magnetic, and electronic ground states that can be sensitively tuned by chemical substitution [5–11]. Phenomena emerging from these compounds such as spin and electronic nematic phases [12,13], magnetic frustration [14,15], and magnetostructural volume-collapse transitions [16–19], and their interrelationships with superconductivity are central issues in condensed-matter physics. Such properties can often be interpreted in terms of either itinerant spin-density-wave (SDW) type or local-moment magnetism [13,20–22] since the compounds are placed somewhere in between these two standard descriptions. Many of these phenomena extend in unique ways to the structurally related  $\text{ACo}_2\text{As}_2$ ,  $A = \text{Ca}, \text{Sr}, \text{Ba}, \text{Eu}$ , cobalt arsenides. While the cobalt-arsenide materials are not found to be superconducting, they harbor signatures of weakly itinerant ferromagnetism (FM) [23,24], unusual spin fluctuations [25,26], and magnetic frustration [26,27], which are tied to flat-band-driven Stoner instabilities [24,28].

Among the cobalt arsenides, tetragonal  $\text{SrCo}_2\text{As}_2$  is unique. While no long-range magnetic order is found, neutron-scattering measurements find evidence for antiferromagnetic (AF) stripe-type spin fluctuations similar to those associated with superconducting pairing in the Fe-based superconductors rather than the expected FM fluctuations [25]. In principle, moving the  $\text{SrCo}_2\text{As}_2$  system closer to stripe-AF order through appropriate chemical substitution may realize superconductivity. However, to date no cobalt-arsenide material has demonstrated long-range stripe-AF order [25,27,29,30].

Recently, the electron-doped materials  $\text{Sr}_{1-x}\text{La}_x\text{Co}_2\text{As}_2$  [31] and  $\text{Sr}(\text{Co}_{1-x}\text{Ni}_x)_2\text{As}_2$  [32,33] have been shown to develop long-range AF order for extremely low substitutions of  $x \approx 1\%$  as shown by the magnetic phase diagram in Fig. 1(a)

for the Ni-doped series. A diagram of the chemical unit cell is given in Fig. 1(b). Here, we study the Ni-doped compounds using neutron and high-energy x-ray diffraction to determine the microscopic details of the AF ground state. Rather than the stripe-AF order one may expect from the spin-fluctuation spectrum of  $\text{SrCo}_2\text{As}_2$  [24,25,34], we find that the series develops *incommensurate* AF order consisting of FM-aligned transition-metal layers where the in-plane ordered magnetic moment ( $\mu \perp c$ ) forms a helix propagating perpendicular to the layers (i.e., along  $c$ ). These results support recent evidence that  $\text{SrCo}_2\text{As}_2$  possesses frustrated magnetic interactions driven by flat-band instabilities that place the system on the border between itinerant two-dimensional FM (2D-FM) and stripe-AF [24,34]. Far more subtle, composition-dependent variations in interlayer interactions and magnetic anisotropy, as recently observed in  $\text{Sr}_{1-x}\text{Ca}_x\text{Co}_2\text{As}_2$ , are at play in determining the details of the layer stacking (such as helical, A-type, or more complex collinear magnetic order) [27].

Detailed sample growth and characterization data of  $\text{Sr}(\text{Co}_{1-x}\text{Ni}_x)_2\text{As}_2$  have recently been described in Ref. [32]. Whereas no AF order is found for  $x = 0$  down to  $T = 50$  mK [34], magnetization and electronic-transport data indicate that small amounts of Ni doping trigger AF order. The dome of AF order spans  $x \approx 0.013$ – $0.25$  with a maximum Néel temperature of  $T_N \approx 27$  K. We performed neutron and x-ray diffraction on  $x = 0.06(1)$  and  $0.15(2)$  single-crystal samples with  $T_N = 20$  and  $25$  K, respectively, as shown in Fig. 1(a).

Neutron diffraction measurements were performed on  $\approx 18$  mg single-crystal samples using the FIE-TAX diffractometer at the High Flux Isotope Reactor, Oak Ridge National Laboratory. Samples were sealed in an Al can containing He exchange gas which was subsequently attached to the cold

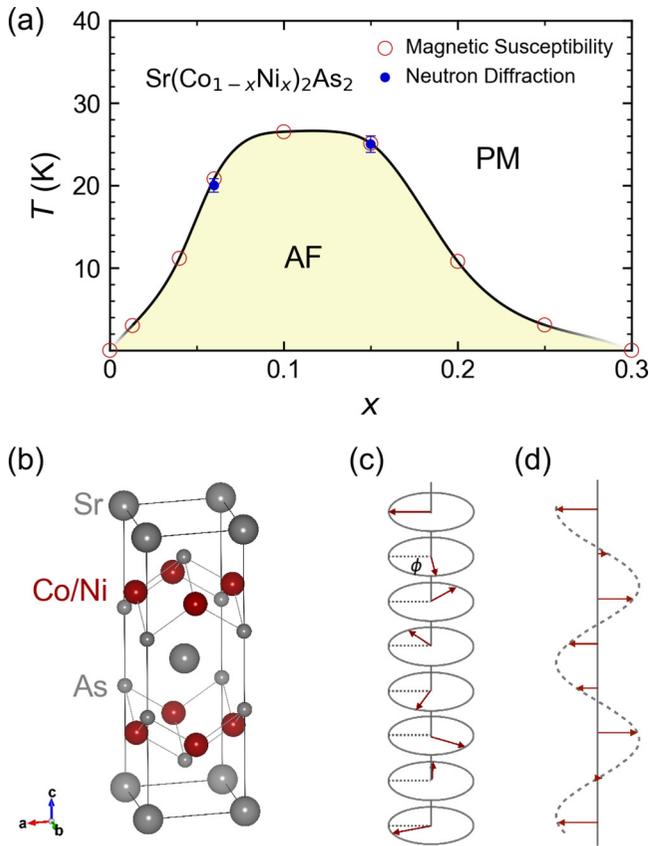


FIG. 1. (a) Magnetic phase diagram for  $\text{Sr}(\text{Co}_{1-x}\text{Ni}_x)_2\text{As}_2$  showing the antiferromagnetically ordered (AF) and paramagnetic (PM) phases. Solid (open) symbols correspond to neutron diffraction (magnetic susceptibility [32]) data. (b) The chemical unit cell, in which magnetic moments exist for the Co/Ni sites. (c) Possible helical-AF arrangement of the magnetic moments with a turn angle of  $\phi$  between ferromagnetically aligned layers. (d) Possible collinear spin-density-wave type order, which cannot be differentiated from helical-AF order shown in (c) using neutron diffraction data.

head of a closed-cycle He refrigerator. The beam collimators placed before the monochromator, between the the monochromator and sample, between the sample and analyzer, and between the analyzer and detector were  $40'-40'-40'-80'$ , respectively. FIE-TAX operates at a fixed incident energy of 14.7 meV using two pyrolytic graphite (PG) monochromators. In order to significantly reduce higher harmonics in the incident beam, PG filters were mounted before and after the second monochromator. The scattering data are described using reciprocal lattice units of  $\frac{2\pi}{a}$  for  $H$  and  $K$  and  $\frac{2\pi}{c}$  for  $L$  within the tetragonal  $\text{ThCr}_2\text{Si}_2$ -type structure (space group  $I4/mmm$ ), where  $a \approx 3.94 \text{ \AA}$  and  $c \approx 11.61 \text{ \AA}$ . A detailed dependence of the lattice parameters on  $x$  is given in Ref. [32]. The samples were aligned with their  $(H H L)$  reciprocal-lattice planes coincident with the spectrometer's scattering plane.

High-energy x-ray diffraction measurements were performed on station 6-ID-D at the Advanced Photon Source, Argonne National Laboratory. Measurements were made using 100 keV x rays, with the incident beam's direction normal to the  $(H H L)$  and  $(H K 0)$  reciprocal-lattice planes. Diffraction

patterns were recorded using a MAR345 area detector. Unlike laboratory sources, high-energy x rays ensure that the bulk of the sample is probed. By rocking the sample through small angular ranges about the axes perpendicular to the incident beam, we obtain an image of the reciprocal-lattice planes normal to the incident beam's direction [35].

We initially made neutron diffraction measurements at positions consistent with a stripe-AF propagation vector,  $\mathbf{q}_{\text{stripe}} = (\frac{1}{2} \frac{1}{2} L)$ , and found that no magnetic Bragg peaks occur at  $\mathbf{q}_{\text{stripe}}$  for either  $x = 0.06$  or  $0.15$ . Rather, scans made along  $(0 0 L)$  and  $(1 1 L)$  revealed weak magnetic Bragg peaks at positions incommensurate with the chemical lattice. Their positions are described by an AF propagation vector of  $(0 0 \tau)$ , with  $\tau = 0.58(1)$  and  $0.52(1)$  for  $x = 0.06$  and  $0.15$ , respectively. The highest intensity Bragg peaks occur at  $(0 0 L \pm \tau)$ ,  $L$  even, and based on the sensitivity of neutrons to the component of  $\boldsymbol{\mu}$  perpendicular to the scattering vector  $\mathbf{Q}$ , we find that  $\boldsymbol{\mu}$  lies within the  $\mathbf{ab}$  plane for both compositions. This result is consistent with magnetization data [32].

Detailed scans along  $(0 0 L)$  at  $T = 4 \text{ K}$  are presented in Fig. 2, in which arrows point from the structural to the magnetic Bragg peaks. The widths of the magnetic Bragg peaks are resolution limited, which attests to the presence of long-range AF order, and scans along  $(1 1 L)$  yield similar results. From measurements made at multiple temperatures, we find for both compositions that, within the resolution of our experiments, no significant change in  $\tau$  or the widths of the magnetic Bragg peaks occur with temperature. The temperature dependence of the integrated intensities of the  $(0 0 2 -\tau)$  Bragg peaks for  $x = 0.06$  and  $0.15$  are shown in Figs. 2(i) and 2(j), and are consistent with a second-order AF transition with  $T_N = 20$  and  $25 \text{ K}$ , respectively. These values agree with  $T_N$  obtained from magnetization data [32], as shown in Fig. 1(a). Overall, the results indicate that instead of the expected stripe-AF order, long-range incommensurate AF order develops in  $\text{SrCo}_2\text{As}_2$  upon electron doping via partial replacement of Co with Ni.

Analysis of the diffraction data establishes that the magnetic structures for both values of  $x$  consist of 2D-FM layers that are AF arranged along  $\mathbf{c}$  and controlled by  $\tau$ . Interestingly, for both compositions,  $(0 0 \tau)$  is close to the commensurate value of  $(0 0 \frac{1}{2})$  found for  $\text{Sr}_{1-x}\text{Ca}_x\text{Co}_2\text{As}_2$ ,  $0.5 \lesssim x \lesssim 0.8$  [27]. In that case, the magnetic order is described by a doubling of the conventional body-centered unit cell along  $\mathbf{c}$ , containing four ferromagnetic Co layers that are stacked  $\leftarrow \leftrightarrow \rightarrow$  along  $\mathbf{c}$  or follow a  $90^\circ$ -clock model.

In a similar vein, analysis of our neutron diffraction data for  $\text{Sr}(\text{Co}_{1-x}\text{Ni}_x)_2\text{As}_2$  determines that the magnetic structure for  $x = 0.06$  and  $0.15$  is either a collinear SDW, where the magnitude of  $\boldsymbol{\mu}$  varies sinusoidally along  $\mathbf{c}$ , or a noncollinear helix where the direction of  $\boldsymbol{\mu}$  varies along  $\mathbf{c}$ . These AF structures are shown in Figs. 1(c) and 1(d). Although we cannot use the neutron diffraction data alone to distinguish between these two possible AF structures due to the presence of domains, combining our neutron diffraction results with a molecular-field analysis of the anisotropic magnetic susceptibility [32] and our high-energy x-ray diffraction data favors the helical-AF structure.

As shown in Fig. 1(c), the AF stacking of the FM-aligned layers in the helical-AF structure can be parametrized by a

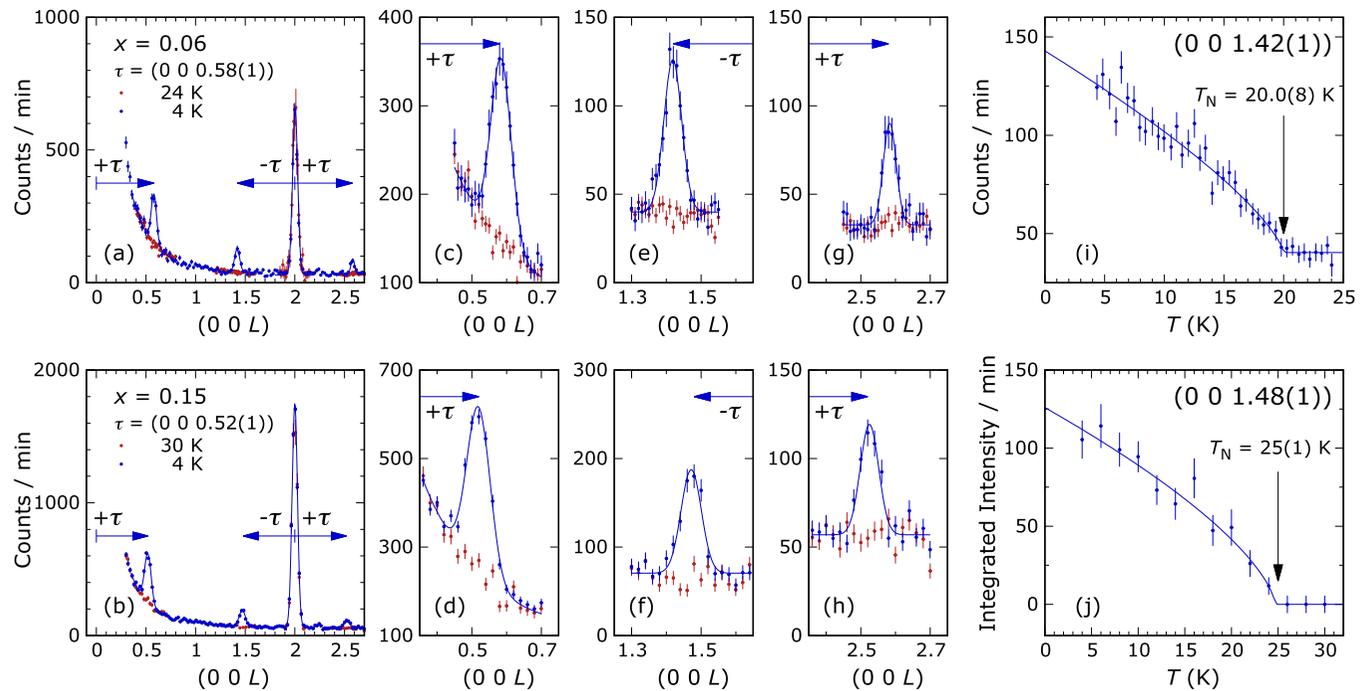


FIG. 2. Neutron diffraction data for  $\text{Sr}(\text{Co}_{1-x}\text{Ni}_x)_2\text{As}_2$  single crystals. (a) and (b) Scans along  $(0\ 0\ L)$  for  $x = 0.06$  and  $0.15$ , respectively. (c)–(h) Expanded views of the magnetic Bragg peaks. (i) and (j) Temperature dependence of the magnetic order parameter for  $x = 0.06$  and  $0.15$ , respectively, with lines as guides to the eye.

turn angle  $\phi$  between each layer. The values of  $\tau$  determined from our neutron diffraction data yield  $\phi = 104^\circ(2)$  and  $94^\circ(2)$ , respectively, whereas the molecular-field analysis of the anisotropic susceptibility [32] finds that  $\phi = 93^\circ$  and  $70^\circ$  for  $x = 0.06$  and  $0.15$ , respectively. Surprisingly, the trend in the dependence of the turn angle,  $\phi$ , on the Ni concentration,  $x$ , is correctly given by the local-moment model analysis of the anisotropic susceptibility [32] of this itinerant magnetic system, and the values of  $\phi$  are somewhat close. We further find from the neutron diffraction data that  $\mu = 0.13(2)$  and  $0.20(2)\mu_B/(\text{Co}+\text{Ni})$  for  $x = 0.06$  and  $0.15$ , respectively, for the helical-AF structure. These values are in quite good agreement with the values for the saturation moment of  $\mu_{\text{sat}} = 0.117(1)$  and  $0.165(1)\mu_B/(\text{Co}+\text{Ni})$  determined for  $x = 0.06$  and  $0.15$ , respectively, from magnetization measurements via considerations using an itinerant FM model [32]. The existence of left- and right-handed helical domains would not affect the intensity of the neutron diffraction peaks.

High-energy x-ray diffraction data were taken on a single-crystal sample of  $x = 0.05$  at  $T = 35$  and  $5$  K to search for any structural anomalies associated with the AF ordering. These data are shown in Fig. 3. The 2D images of the  $(H\ H\ L)$  plane shown in Fig. 3(a) and the detailed cut along  $(1\ 1\ L)$  in Fig. 3(b) show no additional Bragg peaks indicative of a superstructure or a charge-density wave, which is expected to accompany a SDW. In addition, no splitting of Bragg peaks indicative of an orthorhombic lattice distortion are observed as shown in Fig. 3(c). Such a distortion is typically expected for stripe-AF order or a collinear SDW with  $\mu \perp \mathbf{c}$ . Thus, analysis of these data favors the presence of helical-AF order.

Changes to the structure that accompany chemical substitutions are expected to affect (or reflect) the changing magnetic

interactions. In particular,  $\text{CaCo}_2\text{As}_2$  has a value for  $c$  that is  $\approx 13\%$  smaller than that for  $\text{SrCo}_2\text{As}_2$  and exists in the collapsed-tetragonal (cT) phase [36].  $\text{CaCo}_2\text{As}_2$  shows A-type AF order with  $\mu \parallel \mathbf{c}$  [30], the occurrence of which is tied to partially flat electronic bands lying closer to the Fermi energy  $E_F$  than for paramagnetic  $\text{SrCo}_2\text{As}_2$  [28]. The value of  $c$  changes with  $x$  for both  $\text{Sr}(\text{Co}_{1-x}\text{Ni}_x)_2\text{As}_2$  [32,33] and  $\text{Sr}_{1-x}\text{Ca}_x\text{Co}_2\text{As}_2$  [36], and, interestingly, both series show regions of AF order with  $\mu \perp \mathbf{c}$  for ratios of  $c/a$  corresponding to their uncollapsed-tetragonal (ucT) phases [32,36]. The ucT phase is characterized by weaker As-As covalent bonding than in the cT phase [37]. However, the suppression of helical-AF order in  $\text{Sr}(\text{Co}_{1-x}\text{Ni}_x)_2\text{As}_2$  with increasing  $x$  [32] and the emergence of AF order in  $\text{Sr}_{1-x}\text{Ca}_x\text{Co}_2\text{As}_2$  do not appear to coincide with the ucT to cT crossover [36]. Further, the mechanism controlling the direction of  $\mu$  is unclear.

Connections between the unique itinerant magnetic frustration and novel spin fluctuations in cobalt arsenides to partially flat electronic bands have recently been made using density-functional-theory calculations, angle-resolved photoemission spectroscopy, and elastic and inelastic neutron-scattering experiments [27,28,32–34]. The flat bands originate from the transition metals'  $3d_{x^2-y^2}$  orbitals [28,32] and form a sharp peak in the electronic density of states (DOS) close to  $E_F$ . The proximity of the peak to  $E_F$  is operative in determining the magnetism within the transition-metal planes via the Stoner mechanism [28]. Thus, the capability to tune the *intralayer* magnetism requires an understanding of the effects of carrier doping, disorder (and the accompanying smearing of the DOS), and structural modulations (especially those affecting the degree of As-As covalent bonding [37]). All of these

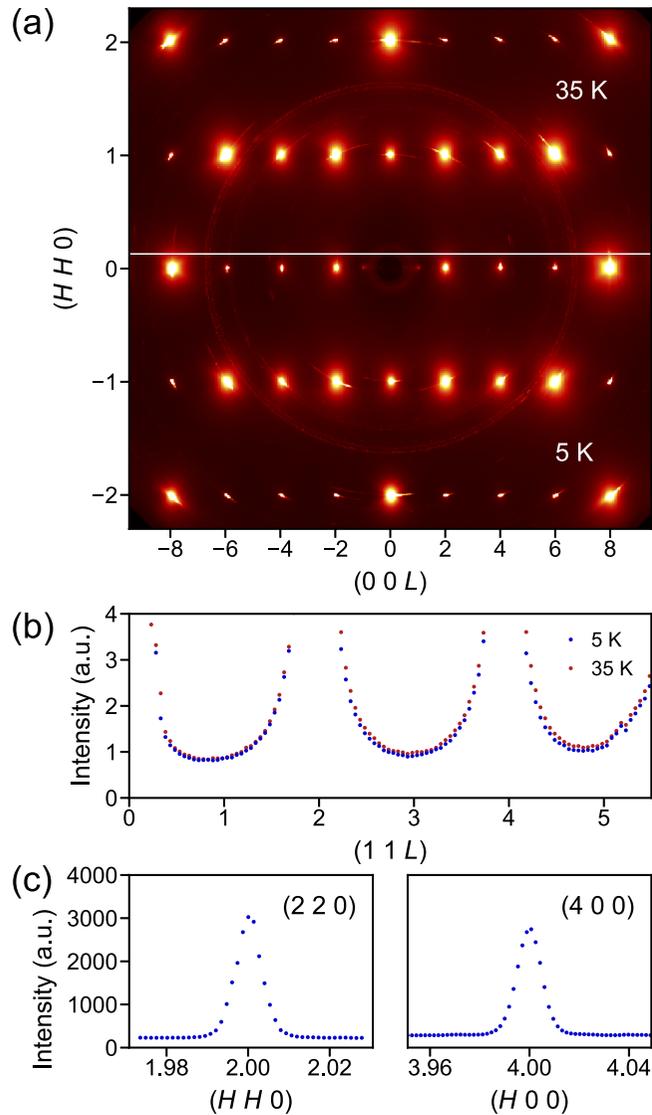


FIG. 3. High-energy x-ray diffraction data for  $\text{Sr}(\text{Co}_{1-x}\text{Ni}_x)_2\text{As}_2$  for  $x = 0.05$ . (a) Image plots of data for the  $(H H L)$  plane at  $T = 35$  and  $5$  K with intensities color coded. Two additional peaks close to the origin are due to the sample holder and are present at all temperatures. (b) Cut along  $(1 1 L)$  demonstrating the lack of CDW superstructure peaks. (c) Cuts through the  $(2 2 0)$  and  $(4 0 0)$  Bragg peaks showing no splitting indicative of a potential orthorhombic lattice distortion.

effects can change the nature of the magnetism within the Co layers. In this respect, the change from an intralayer stripe-AF instability in  $\text{SrCo}_2\text{As}_2$  to a FM instability with small Ni substitution highlights the balance of several competing effects.

The nature by which the 2D FM-aligned layers in AF ordered  $\text{Sr}(\text{Co}_{1-x}\text{Ni}_x)_2\text{As}_2$  stack along  $\mathbf{c}$  involves even more subtle *interlayer* interactions and their susceptibility to chemical substitutions. Similar complexity in the stacking of FM layers is reported for  $\text{Sr}_{1-x}\text{Ca}_x\text{Co}_2\text{As}_2$  [27]. In that case, the cascade of different AF stackings observed for high Ca concentrations is understood on the basis of a one-dimensional classical local-moment Heisenberg model known

as the  $J_0$ - $J_1$ - $J_2$  model [38,39], which is similar to the axial next-nearest-neighbor Ising model [40–42]. For sufficiently large single-ion magnetic anisotropy, the  $J_0$ - $J_1$ - $J_2$  model predicts FM ( $\tau = 0$ ), A-type AF ( $\tau = 1$ ,  $\uparrow\downarrow\uparrow\downarrow$ ), or phases with  $\leftarrow\leftarrow\rightarrow\rightarrow$  or  $90^\circ$ -clock AF structures ( $\tau = \frac{1}{2}$ ). Which phase exists depends on the ratio of the nearest-neighbor and next-nearest-neighbor interlayer interactions,  $J_1$  and  $J_2$ , respectively.

For smaller values of magnetic anisotropy, these models admit incommensurate helical-AF phases, which are single- $\mathbf{Q}$  for the case of zero magnetic anisotropy. Thus, in comparison to  $\text{Sr}_{1-x}\text{Ca}_x\text{Co}_2\text{As}_2$ ,  $\text{Sr}(\text{Co}_{1-x}\text{Ni}_x)_2\text{As}_2$  appears to fall into the regime of small magnetic anisotropy. This is supported by the small value of the spin-flop fields [39] seen for some of the AF ordered Ni-doped compounds [32] in comparison to those for  $\text{Sr}_{1-x}\text{Ca}_x\text{Co}_2\text{As}_2$  [27,36].

Summarizing, our neutron diffraction results for  $\text{Sr}(\text{Co}_{1-x}\text{Ni}_x)_2\text{As}_2$  show that its AF phase does not have stripe-AF order, but rather consists of FM-aligned transition-metal layers with  $\boldsymbol{\mu} \perp \mathbf{c}$ . Unlike  $\text{Sr}_{1-x}\text{Ca}_x\text{Co}_2\text{As}_2$ , the AF arrangement along  $c$  is incommensurate with the chemical lattice and characterized by  $(0 0 \tau)$ , with  $\tau = 0.58(1)$  and  $0.52(1)$  for  $x = 0.06$  and  $0.15$ , respectively. The agreement between  $\tau$  and a molecular-field analysis of anisotropic magnetic susceptibility data, as well as the absence of peaks corresponding to a charge-density wave and the lack of an orthorhombic lattice distortion in our high-energy x-ray diffraction data, tends to support a helical-AF structure rather than a collinear sinusoidally modulated SDW, both of which are consistent with our neutron diffraction data. Applying the helical-AF model to our neutron diffraction data yields  $\mu = 0.13(2)$  and  $0.20(2)\mu_B/(\text{Co}+\text{Ni})$  for  $x = 0.06$  and  $0.15$ , respectively. Our results highlight that, in addition to the highly-tunable intralayer FM driven by the proximity of flat electronic bands to a van Hove singularity, the details of the AF stacking of the FM-aligned layers involve subtle interlayer magnetic interactions which are highly susceptible to doping.

It has been previously emphasized that a local-moment model for the itinerant AF  $\text{CaCo}_{2-y}\text{As}_2$  accurately predicts the observed nearly-perfect magnetic frustration inferred from inelastic neutron-scattering measurements [26]. Here, we have shown that its correspondence holds more generally as previously anticipated [26]. The helical AF structure predicted using a local-moment model for the magnetic susceptibility measurements in Ref. [32] has been confirmed here not only qualitatively but also semiquantitatively using magnetic neutron diffraction measurements.

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