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Abstract

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Keywords

Mossbauer effect, hyperfine parameters, ferromagnetism, $\text{RFe}_2\text{Zn}_{20}$

Disciplines

Condensed Matter Physics

Comments

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Study of ^{57}Fe Mössbauer effect in $R\text{Fe}_2\text{Zn}_{20}$ ($R = \text{Lu}, \text{Yb}, \text{Gd}$).

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Abstract.

We report measurements of ^{57}Fe Mössbauer spectra for $R\text{Fe}_2\text{Zn}_{20}$ ($R = \text{Lu}, \text{Yb}, \text{Gd}$) from ~ 4.5 K to room temperature. The obtained isomer shift values are very similar for all three compounds, their temperature dependence was analyzed within the Debye model and resulted in an estimate of the Debye temperatures of 450-500 K. The values of quadrupole splitting at room temperature change with the cubic lattice constant a in a linear fashion. For $\text{GdFe}_2\text{Zn}_{20}$, ferromagnetic order is seen as an appearance of a sextet in the spectra. The ^{57}Fe site hyperfine field for $T \rightarrow 0$ was evaluated to be ~ 2.4 T.

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1. Introduction

The RT_2Zn_{20} (R = rare earth, T = transition metal) large family of dilute, rare earth containing intermetallic compounds was discovered almost two decades ago, [1] but only a decade later [2, 3] the physical properties of the members of this family started to be of interest for the broader physics community. The studies of this family brought to light six new Yb-based heavy fermion compounds [3], nearly ferromagnetic Fermi liquid behavior in YFe_2Zn_{20} and $LuFe_2Zn_{20}$ [2, 4] and physics of dilute, local magnetic moments in nearly ferromagnetic Fermi liquid. [2, 5] For example the difference in the ordering temperatures and the nature of the magnetic ordered state in the GdT_2Zn_{20} (T = Fe, Ru, Os, Co, Rh, and Ir) was understood in the model of Heisenberg moments embedded in a nearly ferromagnetic, highly polarizable Fermi liquid, combined with the electronic band filling that strongly affect this polarizability. [6] Similar approach, with an addition of the effects of the crystal magnetic field, can be used to understand enhanced Curie temperatures in the RFe_2Zn_{20} (R = rare earth) series. [4, 7, 8, 9]

The RT_2Zn_{20} family was mostly studied using bulk measurements of physical properties. Only few publications [7, 10, 11] reported results of the measurements by local probe techniques. One of such techniques is Mössbauer spectroscopy that allows for the determination of local electronic and magnetic properties. [12] Of the RT_2Zn_{20} family, ⁵⁷Fe Mössbauer spectroscopy was recently used to study $DyFe_2Zn_{20}$ and YFe_2Zn_{20} . [11] The data presented for YFe_2Zn_{20} at 299 K and 78 K indicated a paramagnetic state of this compound, for $DyFe_2Zn_{20}$ the data at 8 K demonstrated presence of the hyperfine field consistent with a long range magnetic order, [2, 4, 5, 6] however a complex combination one one unknown subspectrum and two main subspectra were used to fit the data. [11]

In this work we use ⁵⁷Fe Mössbauer spectroscopy to study the properties of $GdFe_2Zn_{20}$, the compound that has the highest Curie temperature in the RT_2Zn_{20} family [2, 4, 6] and has no possible complications of crystal electric field effects since Gd^{3+} has Hund's rule ground state multiplet J with zero angular momentum. Additionally, we perform measurements on paramagnetic $LuFe_2Zn_{20}$, to compare with the published data on YFe_2Zn_{20} and evaluate the steric effects on the hyperfine parameters, and on the heavy fermion member of the family: $YbFe_2Zn_{20}$.

2. Experimental methods

Single crystalline samples of RFe_2Zn_{20} (R = Lu, Yb, Gd) were grown out of excess Zn using standard solution growth techniques.[13] Initial ratios of starting elements (R:Fe:Zn) were 2:4:94. The constituent elements were placed in an alumina crucible, sealed in a quartz ampule under $\sim 1/3$ atmosphere Ar, heated in a box furnace to 900-1150° C and then slowly cooled to 600° C over ~ 100 h. More details on single crystal growth of these materials are provided in Refs. [2, 3].

Mössbauer spectroscopy measurements were performed using a SEE Co.

conventional constant acceleration type spectrometer in transmission geometry with an ⁵⁷Co(Rh) source, which had an initial (8-9 months before the measurements) intensity 50 mCi, kept at room temperature. For the absorber, the RFe₂Zn₂₀ (*R* = Lu, Yb, Gd) single crystals were powdered using mortar and pestle and mixed with a ZG grade BN powder to ensure homogeneity. The absorber holder comprised two nested white Delrin cups. The absorber holder was locked in a thermal contact with a copper block with a temperature sensor and a heater, and aligned with the γ - source and detector. The absorber was cooled to a desired temperature using a Janis model SHI-850-5 closed cycle refrigerator (with vibration damping). The driver velocity was calibrated by α -Fe foil and all isomer shifts (IS) are quoted relative to the α -Fe foil at room temperature. A commercial software package MossWinn [14] were used to analyze the Mössbauer spectra in this work.

3. Results and discussion

3.1. LuFe₂Zn₂₀ and YbFe₂Zn₂₀

The Fe atoms occupy a single 16*d* site with $\bar{3}m$ (trigonal) point symmetry in the cubic CeCr₂Al₂₀ type structure (*Fd* $\bar{3}m$, *Z* = 8). [1] The non-spherical environment of the Fe atoms (Fig. 1) ensures a nonzero electric field gradient (EFG) at the ⁵⁷Fe site. Indeed, the ⁵⁷Fe Mössbauer spectra of LuFe₂Zn₂₀ and YbFe₂Zn₂₀ taken at different temperatures (Fig. 2) show similar spectral shapes with a clear quadrupole splitting. Each of these spectra can be fitted with a single doublet, consistent with Fe occupying a single, unique crystallographic site. The spectra are very similar for both compounds and are very weakly temperature-dependent.

The temperature dependent hyperfine parameters for LuFe₂Zn₂₀ and YbFe₂Zn₂₀ are shown in Fig. 3. These parameters are very similar for both compounds. The linewidth is basically temperature independent. The isomer shift increases with decrease in temperature. The main contribution to the temperature dependence of the isomer shift comes from the second order Doppler shift and is usually described by the Debye model:

$$IS(T) = IS(0) - \frac{9 k_B T}{2 M c} \left(\frac{T}{\Theta_D} \right)^3 \int_0^{\Theta_D/T} \frac{x^3 dx}{e^x - 1}, \quad (1)$$

where *c* is the velocity of light, *M* is the mass of the ⁵⁷Fe nucleus, and *IS*(0) is the temperature-independent part, i.e. the chemical shift. The obtained values of the Debye temperature, Θ_D are 524 ± 47 K and 487 ± 17 K for LuFe₂Zn₂₀ and YbFe₂Zn₂₀ respectively.

The quadrupole splitting increases slightly, by $\sim 1.5 - 3\%$ on cooling from room temperature to the base temperature.

Temperature dependence of the relative spectral area (SA) in absence of phase

transitions can be described within the Debye model as well.

$$f = \exp \left[\frac{-3E_\gamma^2}{k_B \Theta_D M c^2} \left\{ \frac{1}{4} + \left(\frac{T}{\Theta_D} \right)^2 \int_0^{\Theta_D/T} \frac{x dx}{e^x - 1} \right\} \right], \quad (2)$$

where f is the recoilless fraction, which is proportional to the spectral area for thin sample and E_γ is the γ -ray energy. The values of Θ_D obtained from the fits are 410 ± 32 K and 433 ± 25 K for LuFe₂Zn₂₀ and YbFe₂Zn₂₀ respectively.

3.2. GdFe₂Zn₂₀

The ⁵⁷Fe Mössbauer spectra of GdFe₂Zn₂₀, taken at different temperatures, are presented in Fig. 4. Whereas at high temperatures the observed spectra are doublets and are very similar to those measured for LuFe₂Zn₂₀ and YbFe₂Zn₂₀ (Fig. 2), at the base temperature the spectrum splits in six lines indicating the presence of magnetic hyperfine field at the Fe site, in agreement with ferromagnetic ground state of GdFe₂Zn₂₀ proposed based on the bulk magnetic measurements. [2] The low temperature spectra for GdFe₂Zn₂₀ depend on the electric field gradient and magnetic hyperfine field, B_{hf} , at the ⁵⁷Fe nuclei, as well as on the angle between the principal axis of the EFG tensor and the direction of the B_{hf} . The "mixed Q + M static Hamiltonian (powder)" model (the z -axis is parallel to the principal axis of the EFG tensor, the static Hamiltonian includes magnetic and quadrupole interactions with an arbitrary relative orientation) of the MossWinn [14] program was used to analyze these spectra. Due to the relative symmetry of the environment of the Fe atom in the lattice (see Fig. 1), the EFG assymetry parameter, $\eta = (V_{xx} - V_{yy})/V_{zz}$ and the azimuthal angle of the hyperfine field, α were fixed at zero for the fits.

Hyperfine field at the Fe site, B_{hf} as a function of temperature is shown in Fig. 5. At the base temperature the value of B_{hf} is close to 2.4 T. It has been observed [15] that temperature dependence of the hyperfine field follows closely (even if not exactly) the saturation magnetization curve. We can fit $B_{hf}(T)$ data with the function based on the Curie-Bloch equation [16] for saturation magnetization.

$$B_{hf} = B_{hf}(0) \left[1 - \left(\frac{T}{T_C} \right)^\alpha \right]^\beta, \quad (3)$$

where α is an empirical constant and β is a critical exponent. The value $\beta = 1/3$ [16] was used in the fit. As the result of the fit the following values were obtained: $B_{hf}(0) = 2.36 \pm 0.03$ T, $T_C = 84.6 \pm 0.8$ K, $\alpha = 1.4 \pm 0.1$. The bulk measurements of GdFe₂Zn₂₀ [2, 6] yield a $T_C = 86 \pm 2$ K value of the Curie temperature, that is consistent with the results of the fit and with our observation of a magnetic sextet in the 80 K spectrum and a non-magnetic doublet in the 90 K spectrum (Fig. 4).

The values of the V_{zz} component of the EFG tensor and of the polar angle β between the directions of V_{zz} and B_{hf} (Fig. 5) show only insignificant change between ~ 4.5 K and 80 K.

The temperature dependent hyperfine parameters of GdFe₂Zn₂₀ are shown in Fig. 6. The values of $IS(T)$, $QS(T)$ and the general behavior of $SA(T)$ are consistent with those found for LuFe₂Zn₂₀ and YbFe₂Zn₂₀. The overall behavior of the isomer shift can be described by the Debye model with $\Theta_D = 465 \pm 23$ K, similarly, above the magnetic transition, the Debye model with $\Theta_D = 450 \pm 14$ K can be used to describe the temperature dependent relative spectral area. It should be noted that an increase in the spectral area (beyond to what is expected in the Debye model) is apparent on decrease of temperature below the transition. Since significant change in phonon spectra is not anticipated in the ferromagnetic state, and since the $IS(T)$ behavior in GdFe₂Zn₂₀ is very similar to those in LuFe₂Zn₂₀ and YbFe₂Zn₂₀, probably some unknown artifacts contribute to the curious $SA(T)$ behavior shown in Fig. 6. The quadrupole splitting is only weakly temperature dependent both above and below the magnetic transition. $QS(295K) = 0.637 \pm 0.005$ mm/s that is $\sim 10\%$ lower than the room temperature values for LuFe₂Zn₂₀ and YbFe₂Zn₂₀ (0.712 ± 0.002 and 0.708 ± 0.003 , respectively).

3.3. Discussion

The values and the temperature dependence of the isomer shift for all three compounds studied in this work are very similar and are consistent with the results reported for YFe₂Zn₂₀ and DyFe₂Zn₂₀. [11] The Debye temperature, evaluated from the $IS(T)$ data is $\sim 450 - 500$ K for all of them. Similar but slightly lower values of Debye temperature, $\sim 410 - 450$ K, were estimated for RFe₂Zn₂₀ ($R = \text{Lu, Yb, Gd}$) from the temperature dependence of the relative spectral area. This discrepancy may be explained by the fact the area reflects the average mean-square displacements, whereas IS is related to the mean-square velocity of the Mössbauer atom. Both quantities may respond in different ways to the lattice anharmonicities. It has to be mentioned that the Θ_D values inferred from the temperature dependence of the hyperfine parameters above are significantly higher than $\Theta_D \approx 340$ K evaluated for YFe₂Zn₂₀ and LuFe₂Zn₂₀ from the low temperature specific heat capacity. [4] This discrepancy might be related to the fact that the Θ_D estimates from the temperature dependent Mössbauer spectra reflect, to a large extent, the phonon frequencies related to the Mössbauer atom (⁵⁷Fe), whereas specific heat capacity involves more uniform average. Additionally, other, e.g. electronic, contributions might play a role in the temperature evolution of the IS and SA in this family of materials with nearly ferromagnetic Fermi liquid matrix.

The ⁵⁷Fe hyperfine parameters of the Yb-based heavy fermion YbFe₂Zn₂₀ and nearly ferromagnetic Fermi liquid compound LuFe₂Zn₂₀ are virtually indistinguishable. Possibly the effects of additional electron correlations below the Kondo temperature of ~ 33 K in YbFe₂Zn₂₀, [3] as probed on the Fe site, are small in comparison with the effects of strongly correlated Fermi liquid governed by the transition metal density of states.[2]

The quadrupole splitting values are governed by the local environment of the Fe atoms (Fig. 1). To follow the geometrical changes in this environment with the change of

rare earth ionic radius or the lattice parameters is a cumbersome task and does require a detailed knowledge of atomic coordinates for each compound. Experimentally though, it appears that the QS at room temperature decreases as a function of the a -lattice parameter in close to linear fashion (Fig. 7) suggesting the trend of the Fe environment to become more isotropic with increase of a . This observation is consistent with a very simple notion the average electric field on the Fe site becoming "more spherical" with point charges of the local environment moving further apart.

The ferromagnetic transition in GdFe₂Zn₂₀ is seen in ⁵⁷Fe Mössbauer spectra as appearance of a sextet (Fig. 4), additionally, there is an increase in the relative spectral area near Curie temperature (Fig. 6). The value of the hyperfine field for $T \rightarrow 0$ is $B_{hf}(0) = 2.36 \pm 0.03$ T, that is $\sim 25\%$ higher than ~ 1.84 T measured in DyFe₂Zn₂₀. [11] It is possible that the crystal electric field effect present in DyFe₂Zn₂₀ and absent in GdFe₂Zn₂₀ should be considered in the analysis of this difference.

The temperature dependence of the ⁵⁷Fe site hyperfine field in GdFe₂Zn₂₀ allows us to address the question of interpretation of the strong deviation of the temperature dependence of magnetic susceptibility from the Curie-Weiss law between T_C and ~ 250 K, discussed in Ref. [6]. Absence of hyperfine field on the ⁵⁷Fe site at and above 90 K, argues against the picture of formation of magnetic droplets consisting of the Gd⁺³ local moments and oppositely polarized electron cloud, leaving the hypothesis of the temperature dependent coupling between Gd³⁺ local moments as more viable.

4. Summary

To summarize, we performed measurements of ⁵⁷Fe Mössbauer spectra for RFe₂Zn₂₀ ($R = \text{Lu, Yb, Gd}$) from ~ 4.5 K to room temperature. The isomer shift values are very similar for all three compounds, and give an estimate of Θ_D in the range of 480-500 K. The values of quadrupole splitting at room temperature change with the lattice constant in a linear fashion. Ferromagnetic order is seen as a sextet in the spectra. The $B_{hf}(T)$ data can be fitted with the function based on the Curie-Bloch equation for saturation magnetization that yields the value of B_{hf} at $T \rightarrow 0$ of ~ 2.4 T. No features in the Mössbauer spectra were associated with the Yb-based heavy fermion behavior in the YbFe₂Zn₂₀ at low temperatures.

This study suggests that further work on understanding of the observed trend in QS within this family as well as on elucidation of the relation of the measured hyperfine field and bulk magnetic properties of the compounds in the RFe₂Zn₂₀ family is desired.

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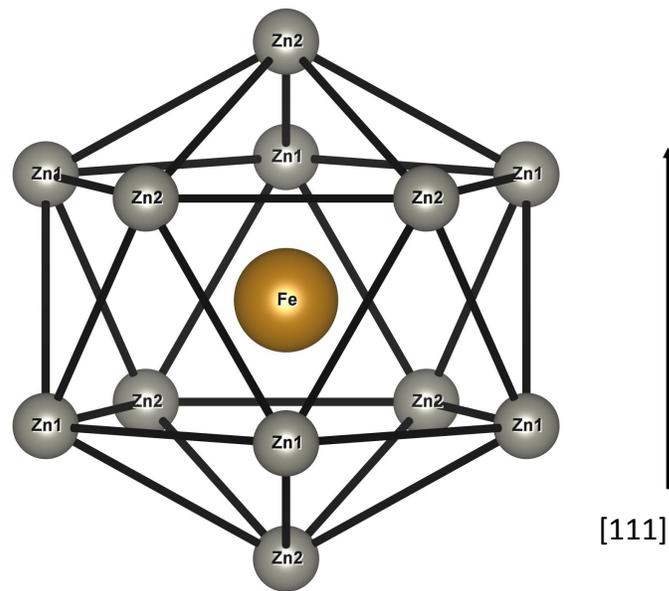
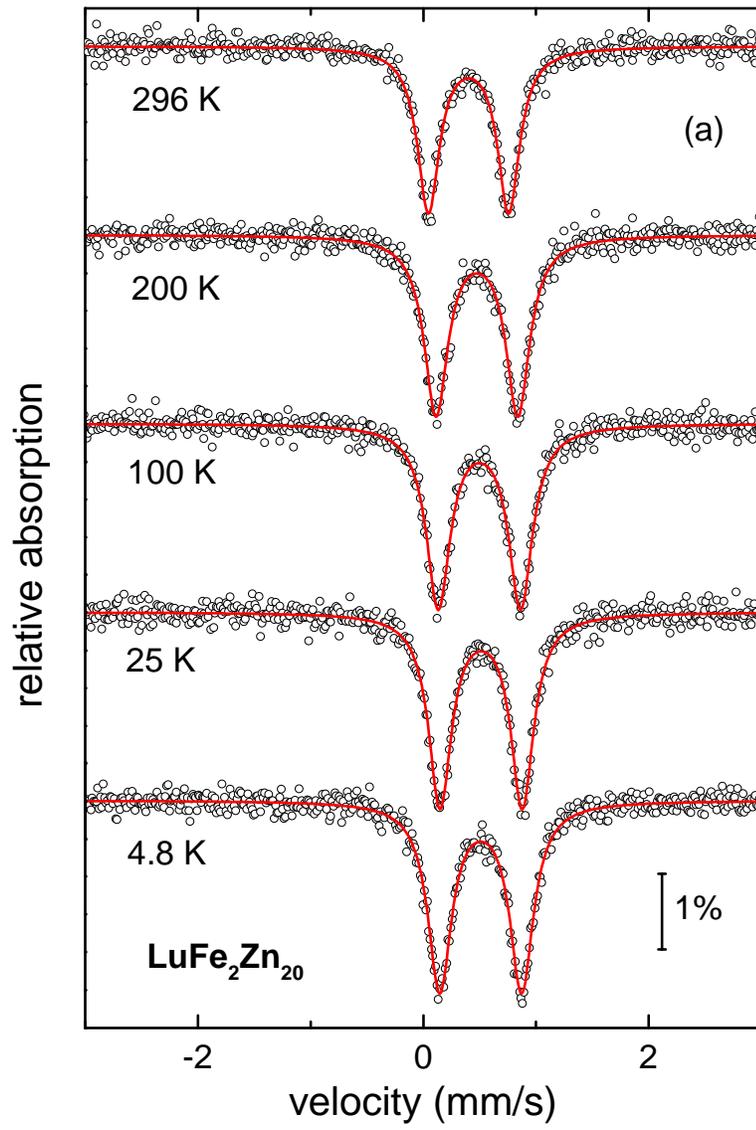


Figure 1. (Color online) Local environment of the Fe atoms in $R\text{Fe}_2\text{Zn}_{20}$. (After Ref. [1]).



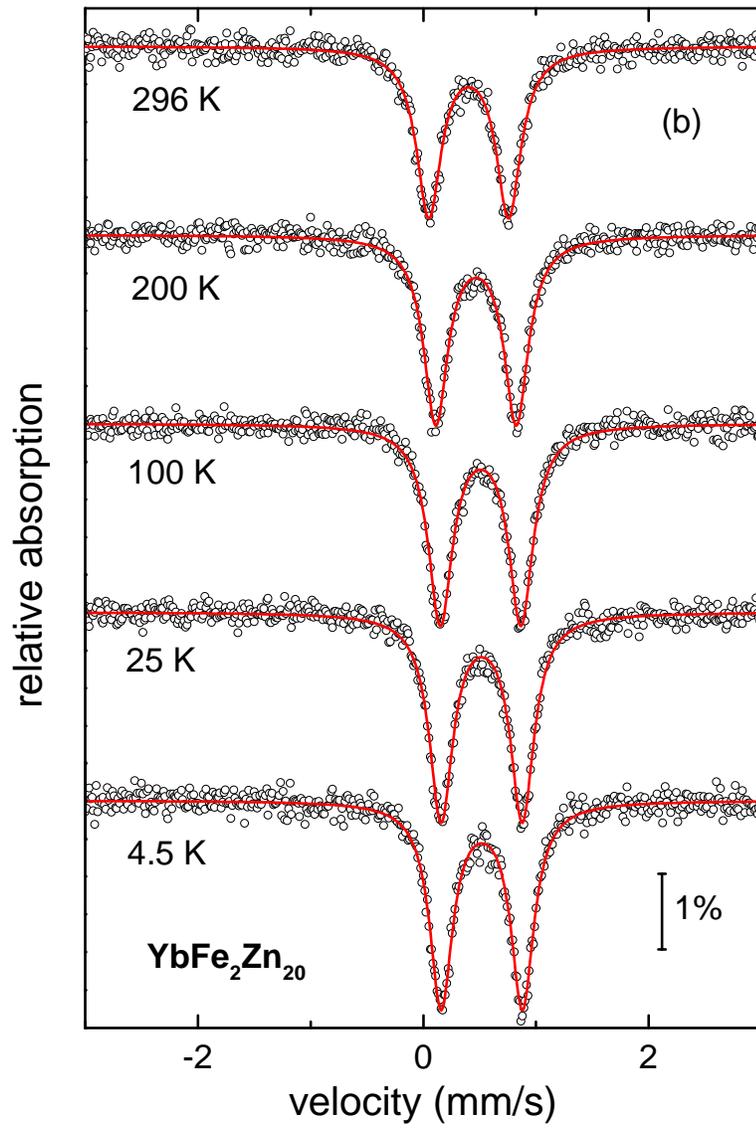


Figure 2. (Color online) ^{57}Fe Mössbauer spectra of (a) $\text{LuFe}_2\text{Zn}_{20}$ and (b) $\text{YbFe}_2\text{Zn}_{20}$ at selected temperatures. Circles - experimental data, lines - fits.

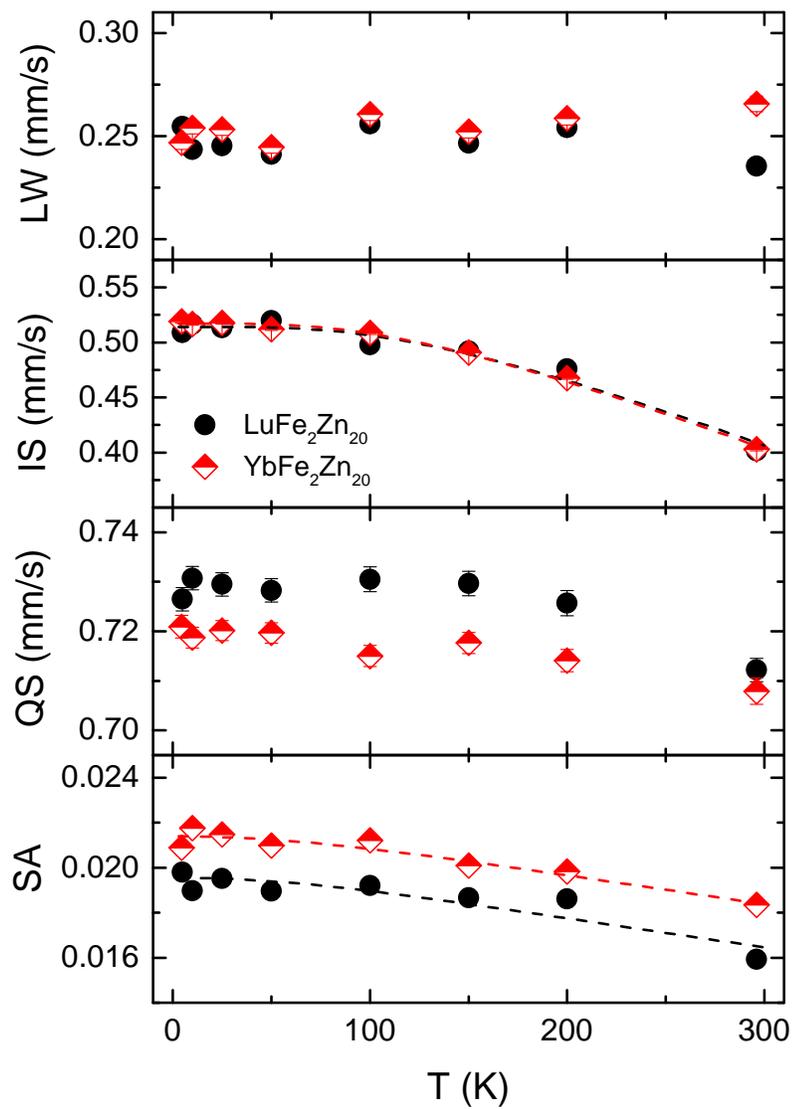


Figure 3. (Color online) Temperature dependence of the hyperfine parameters obtained from ^{57}Fe Mössbauer spectra of $\text{LuFe}_2\text{Zn}_{20}$ (filled circles) and $\text{YbFe}_2\text{Zn}_{20}$ (half filled diamonds). LW - linewidth, IS - isomer shift, QS - quadrupole splitting, SA - normalized spectral area. Dashed lines - Debye fits as described in the text.

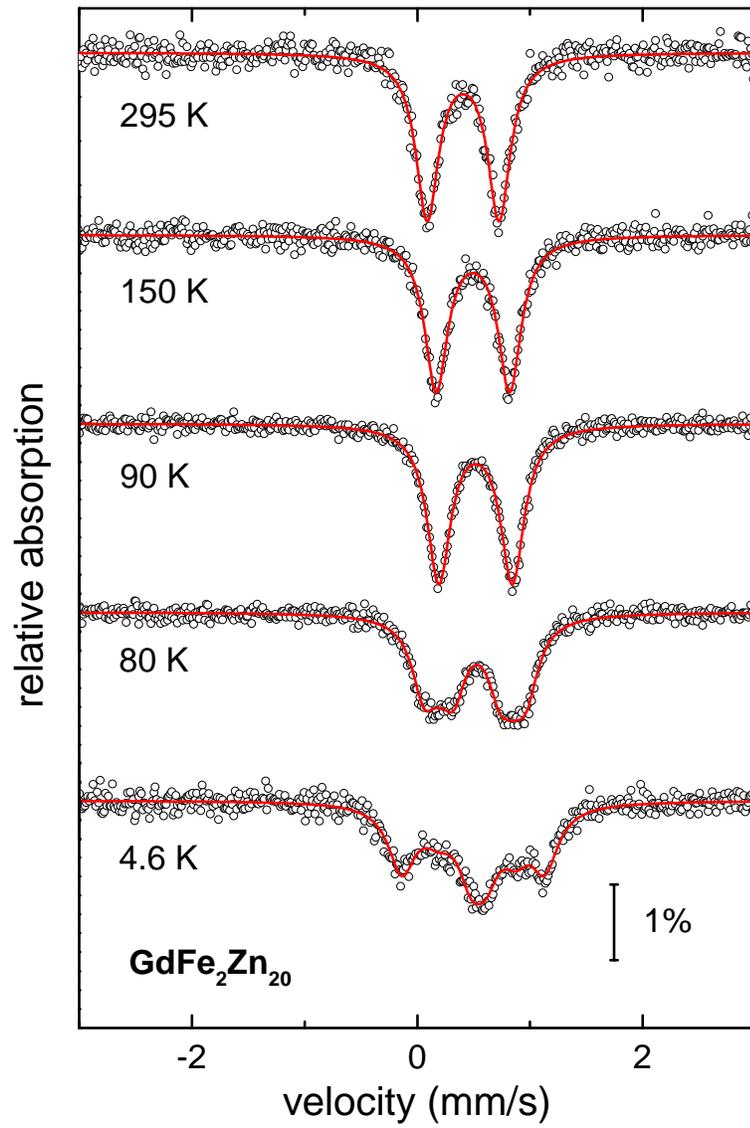


Figure 4. (Color online) ^{57}Fe Mössbauer spectra of $\text{GdFe}_2\text{Zn}_{20}$ at selected temperatures. Circles - experimental data, lines - fits.

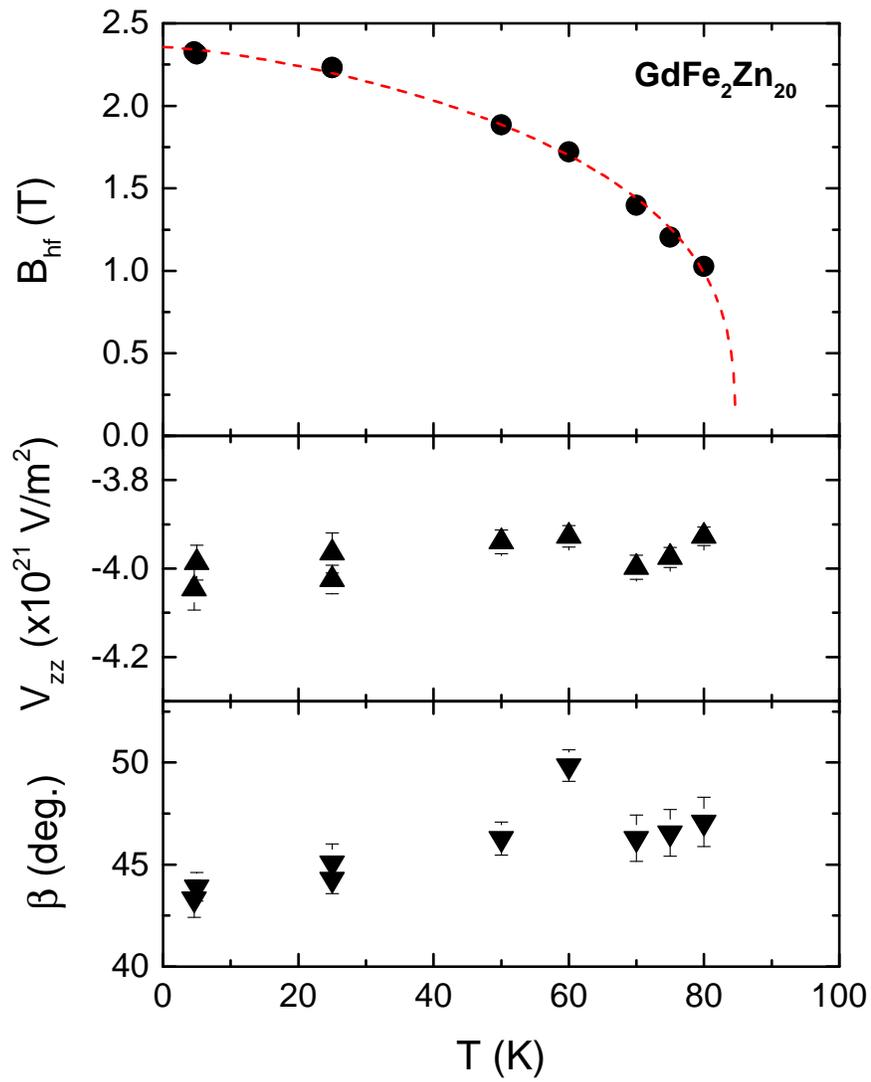


Figure 5. (Color online) Top panel: magnetic hyperfine field, $B_{hf}(T)$, for $\text{GdFe}_2\text{Zn}_{20}$. Symbols - data, dashed line - fit to Curie - Bloch equation - see the text. Middle panel: V_{zz} component of the electric field gradient tensor as a function of temperature. Bottom panel: temperature dependence of the angle between V_{zz} and the direction of the B_{hf} .

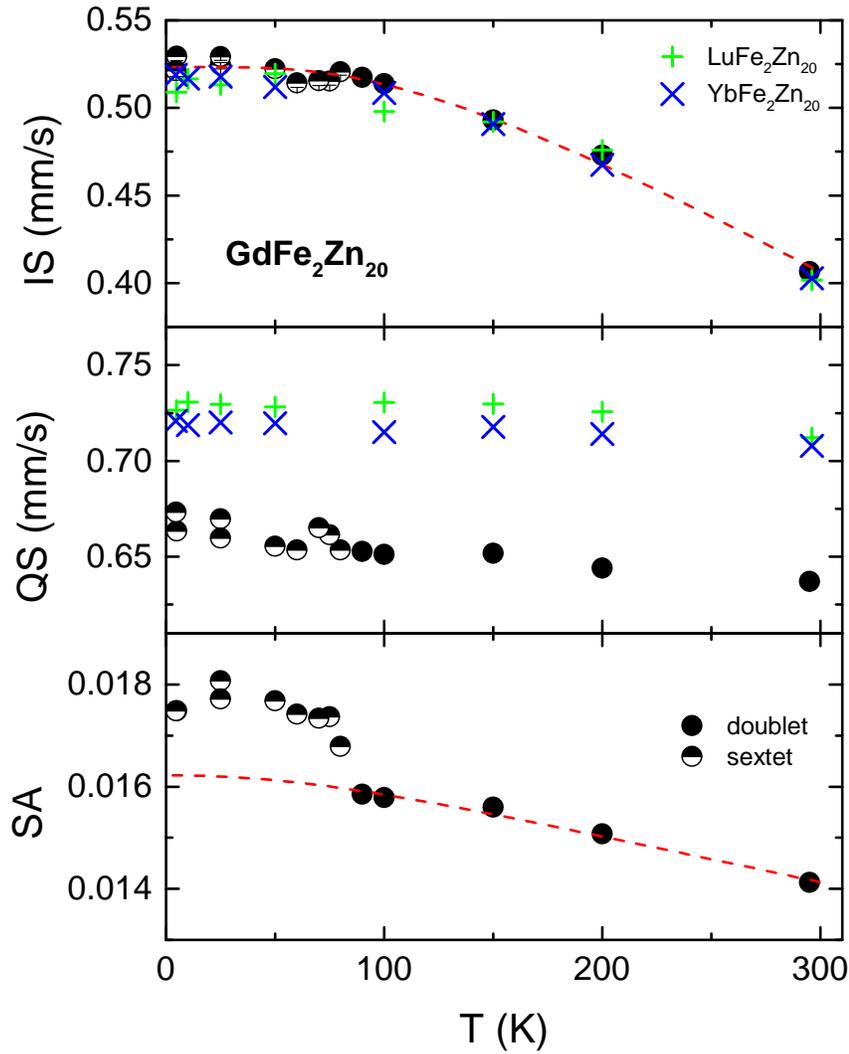


Figure 6. (Color online) Temperature dependence of the hyperfine parameters obtained from ^{57}Fe Mössbauer spectra of $\text{GdFe}_2\text{Zn}_{20}$. IS - isomer shift, QS - quadrupole splitting, SA - normalized spectral area. Dashed lines - Debye fits as described in the text. For $IS(T)$ and $QS(T)$ data for $\text{LuFe}_2\text{Zn}_{20}$ (+) and $\text{YbFe}_2\text{Zn}_{20}$ (\times) are plotted for comparison.

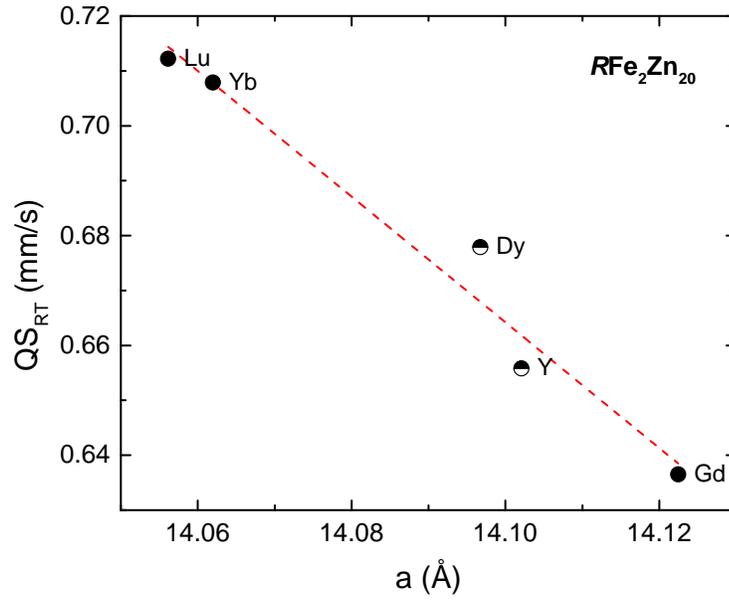


Figure 7. (Color online) Quadrupole splitting at room temperature as a function of the a lattice parameter. Filled circles - this work, half-filled circles - from Ref. [11]. Dashed is a guide to the eye. Values of the lattice parameter a were taken from Ref. [4]